

Methoden moderner Röntgenphysik: Streuung und Abbildung

Lecture 15	Vorlesung zum Haupt- oder Masterstudiengang Physik, SoSe 2018 G. Grübel, A. Philippi-Kobs, <u>F. Lehmkuhler</u> , O. Seeck, L. Frenzel, M. Martins, W. Wurth		
Location	Lecture hall AP, Physics, Jungiusstraße		
Date	Tuesday	13:00 - 14:30	(starting 3.4.)
	Thursday	8:30 - 10:00	(until 12.7.)



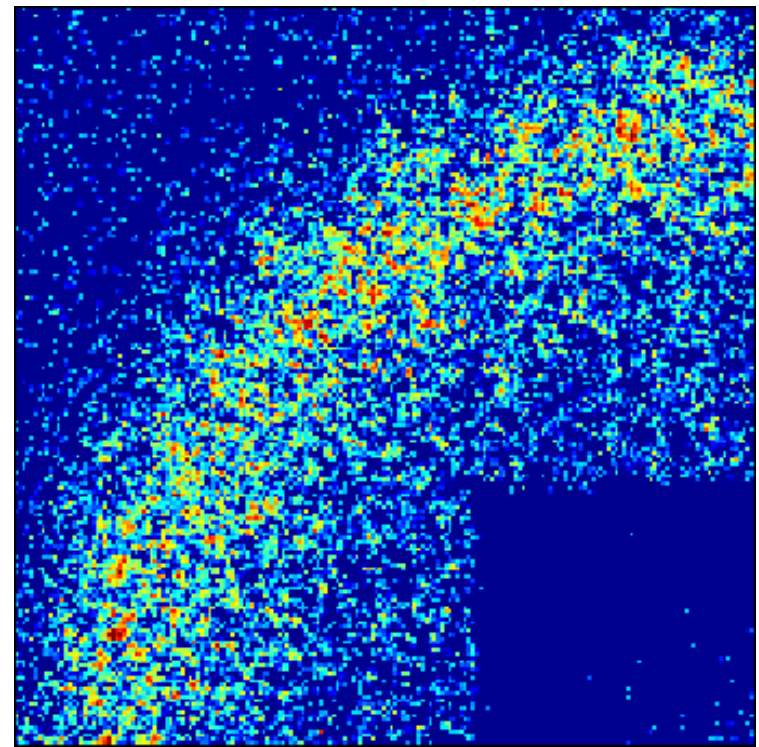
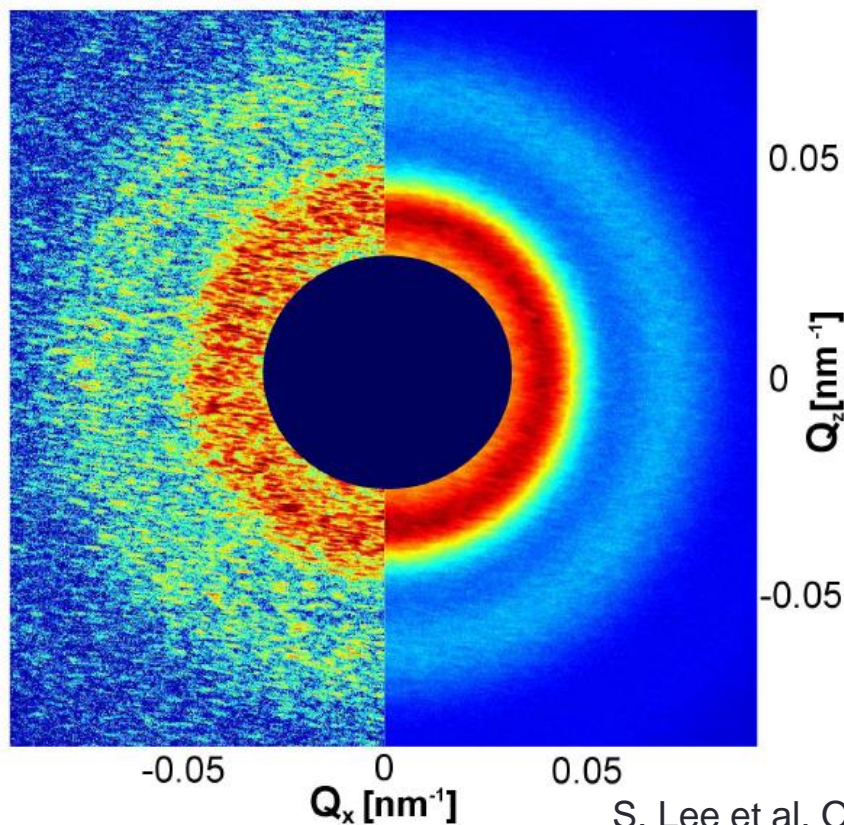
Soft Matter – Timeline

- Di 29.05.2018 Soft Matter studies I: Methods & experiments
Definitions, complex liquids, colloids, storage ring and FEL experiments, setups, liquid jets, ...
- Do 31.05.2018 Soft Matter studies II: Structure
SAXS & WAXS applications, X-ray cross correlations, ...
- Di 05.06.2018 **Soft Matter studies III: Dynamics**
XPCS applications, diffusion, dynamical heterogeneities, ...
- Do 07.06.2018 no lecture!
- Di 12.06.2018 Case study I: Glass transition **!!! CFEL Sem. Rm. 1 !!!**
Supercooled liquids, glasses vs. crystals, glass transition concepts, structure-dynamics relations, ...
- Do 14.06.2018 Case study II: Water
Phase diagram, anomalies, crystalline and glassy forms, FEL studies, ...



Probing dynamics with coherent X-rays: X-ray photon correlation spectroscopy (XPCS)

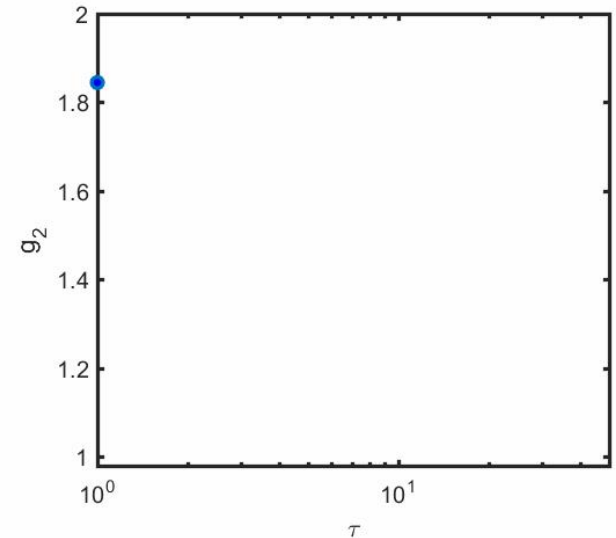
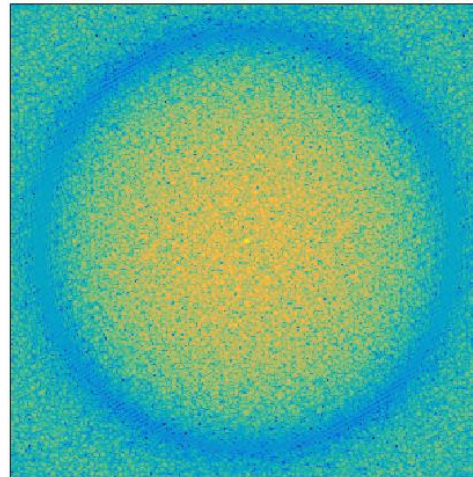
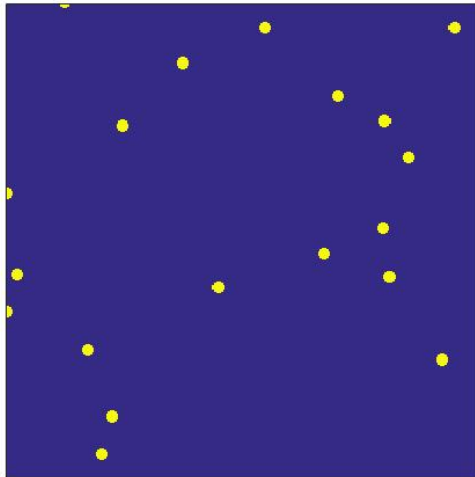
X-ray scattering from disordered samples: speckles
→ structure decoded



S. Lee et al. Optics
Express 21, 24647 (2013)

XPCS

- **Time domain:** changing sample structure → change of speckle pattern
- Correlation function $g_2(q, \tau) = \frac{\langle I(q,t)I(q,t+\tau) \rangle_t}{\langle I(q,t) \rangle_t^2} = 1 + \beta^2 |f(q, \tau)|^2$, speckle contrast $\beta = \text{std}(I)/\langle I \rangle$
- Intermediate scattering function $f(q, \tau) = S(q, \tau)/S(q, 0)$



Diffusing particles

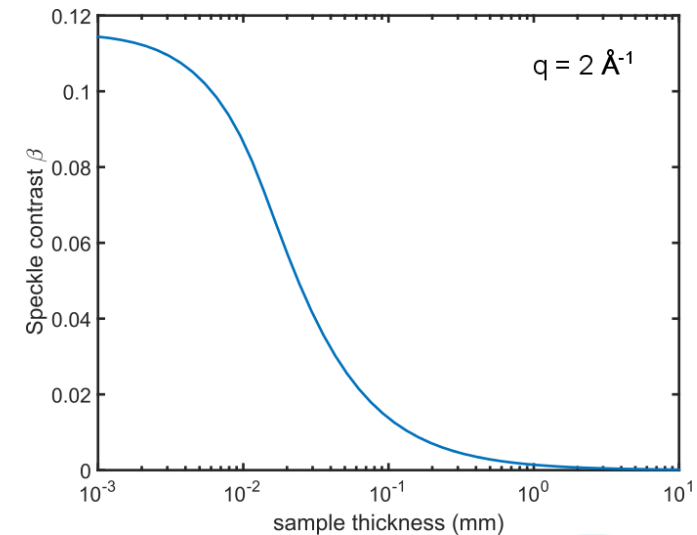
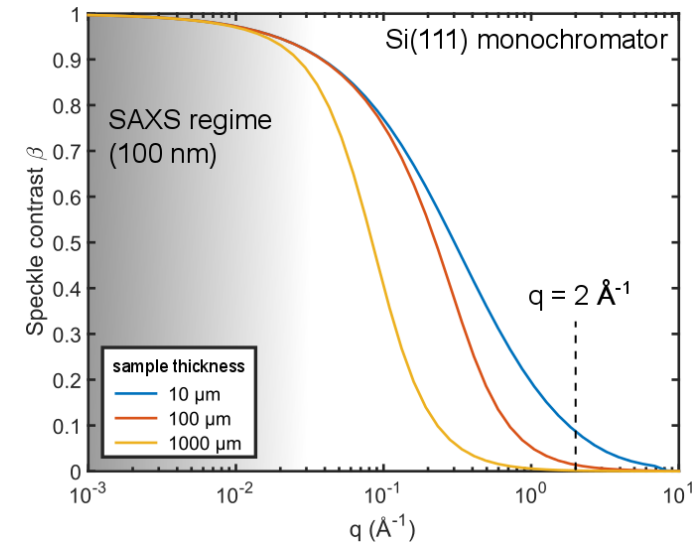
Speckle pattern

g_2 function

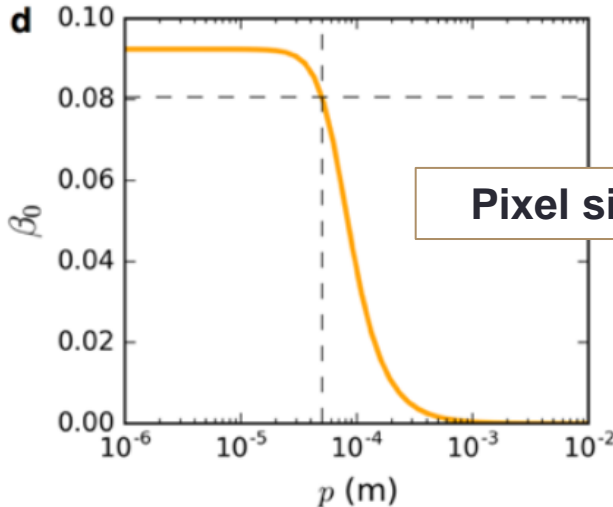
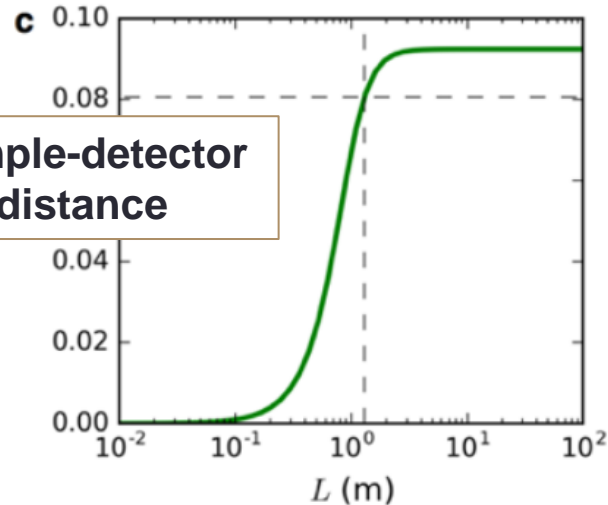
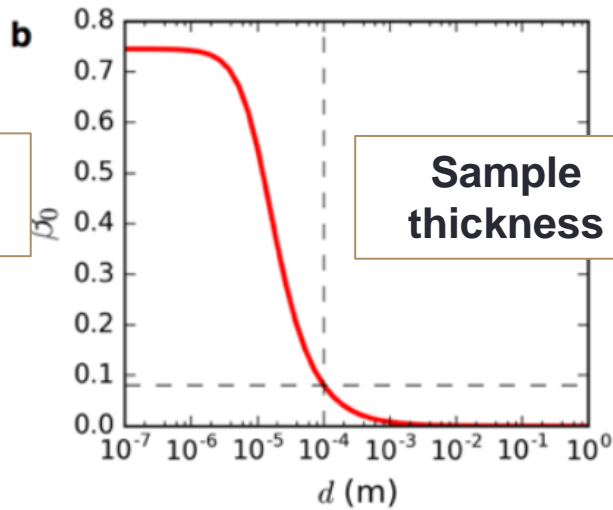
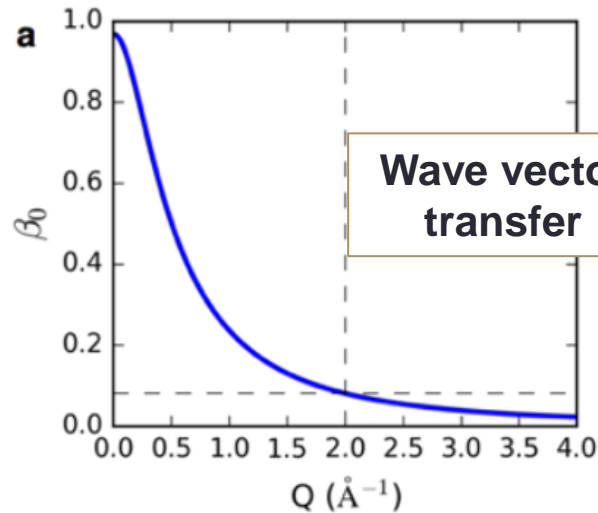


XPCS experiments – requirements

- Degree of coherence → contrast
- Need to resolve speckles
 - speckle size $s \approx \frac{\lambda D}{b}$
 - using hard X-rays ($\lambda \sim 10^{-10}$ m)
 - $\frac{bs}{D} = 10^{-10}$ m
 - $bs \sim 10^{-10}$ m² for $D \sim 1$ m
- Statistics and q-dependence: 2D detectors (e.g. CCD)
 - Typical pixel sizes of $\sim 10 - 100$ μm
 - Consequently beam sizes in the μm regime
- Limit of time scales by detector read-out
 - CCD: \sim seconds
 - Photon counting: $>$ kHz



XPCS experiments – requirements



Speckle contrast as a function of various parameters

Nature Comm. 9, 1917 (2018).



Diffusion in Soft Matter

- Brownian motion: random movement of particles (pollen collision with water molecules (Einstein 1905))
- Omnipresent in soft matter systems
- Derivation (after Langevin, here only one direction x):

$$m \frac{d^2x}{dt^2} = F - f \frac{dx}{dt}$$

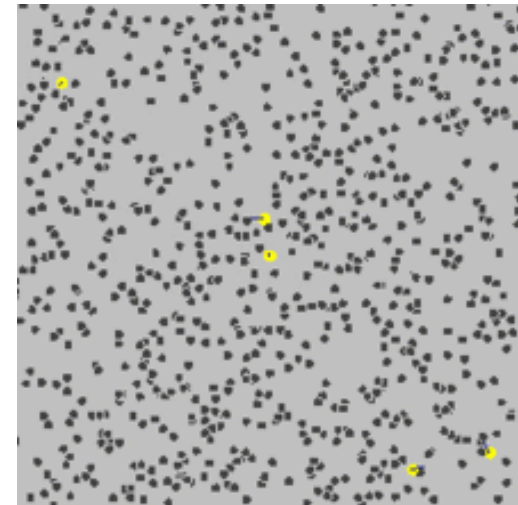
(with force F and viscous friction $F_R = -f \frac{dx}{dt}$)

$$\Leftrightarrow m \frac{d}{dt} \langle v \rangle = \langle F \rangle - f \langle v \rangle \quad (\text{Averaging})$$

$\langle F \rangle = 0$ for random particle collisions

$$\frac{d}{dt} \langle v \rangle = -\frac{f}{m} \langle v \rangle$$

$$\Rightarrow \langle v(t) \rangle = v(0) \exp\left(-\frac{m}{f} t\right)$$



Diffusion of particles

Diffusion in Soft Matter

→ Mean drift velocity $\langle v \rangle$ decays with time. Back to $m \frac{d}{dt} v = F - f v$. Multiply by instantaneous position r of an particle and average yields:

$$\frac{d^2}{dt^2} \langle r^2 \rangle + \frac{f}{m} \frac{d}{dt} \langle r^2 \rangle = 2 \langle v^2 \rangle$$

Following the equipartition theorem ($\langle v^2 \rangle = \frac{3k_B T}{m}$) the equation can be solved with the result

$$\langle r^2 \rangle = \frac{6k_B T m}{f^2} \left(\frac{f}{m} t - \left[1 - \exp \left(-\frac{f}{m} t \right) \right] \right)$$

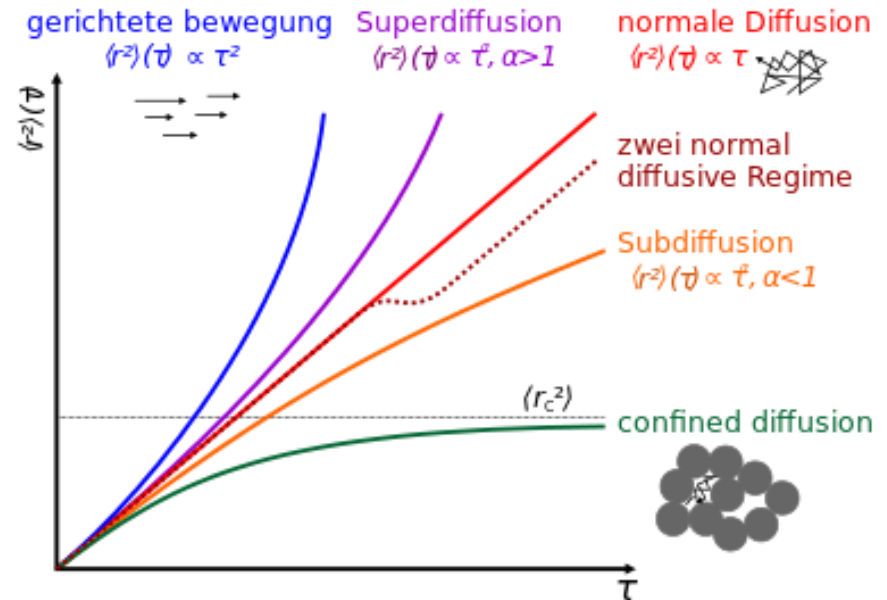
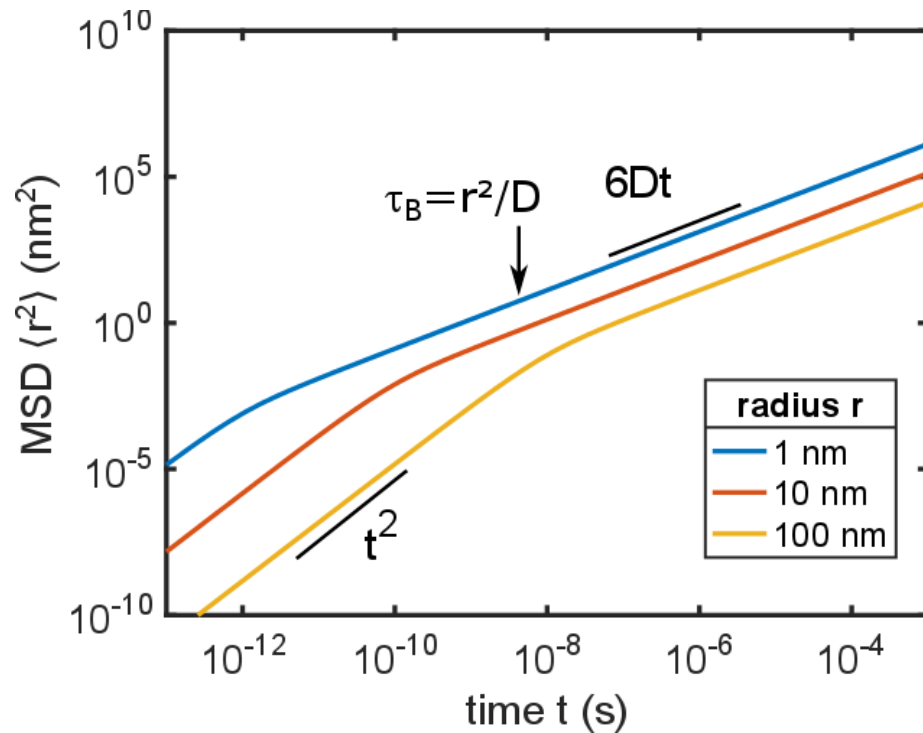
For $t \gg \frac{m}{f}$ we obtain with Stoke's law (friction of spheres, $f = 6\pi r \eta$)

$$\langle r^2 \rangle = \left(\frac{k_b T}{\pi R \eta} \right) t = 6Dt \text{ with diffusion coefficient } D = \frac{k_b T}{6\pi \eta R}$$



Diffusion in Soft Matter

Mean squared displacement – particles in water



Characteristic time $\tau_b = \frac{R^2}{D}$ to move by one radius (here $4.5 \cdot 10^{-9}$ s for $R = 1$ nm)

Diffusion in Soft Matter – XPCS

Intermediate scattering function $f(q, \tau) = S(q, \tau)/S(q, 0)$ with

$$S(q, \tau) = \frac{1}{N} \left\langle \sum_{i=1}^N \sum_{j=1}^N \exp(i\mathbf{q} \cdot [\mathbf{r}_i(0) - \mathbf{r}_j(\tau)]) \right\rangle$$

For diffusion, only single particle properties are probed \rightarrow cross terms $i \neq j$ average out and $S(q) = 1 \rightarrow$ we obtain

$$f(q, \tau) = \frac{1}{N} \left\langle \sum_{i=1}^N \exp(i\mathbf{q} \cdot [\mathbf{r}_i(0) - \mathbf{r}_i(\tau)]) \right\rangle$$

And finally (cf. Physica 32, 415 (1966)) the result for diffusion

$$f(q, \tau) = \exp(-Dq^2\tau)$$



Diffusion by XPCS – Notes

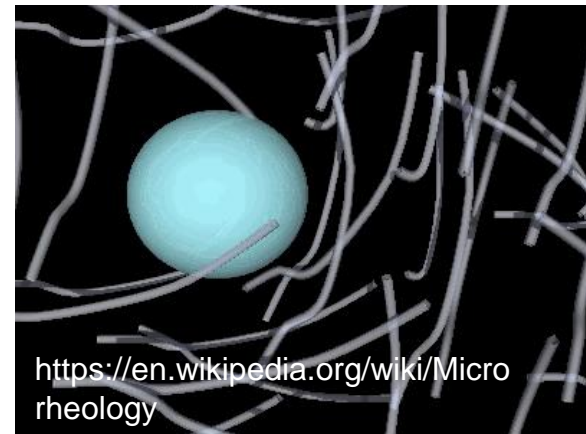
- In XPCS, correlation function for diffusion:

$$g_2(q, \tau) = 1 + \beta^2 |f(q, \tau)|^2 = 1 + \beta^2 \exp(-2Dq^2\tau)$$

- Relaxation time $\tau_0 = \frac{1}{\Gamma} = \frac{1}{Dq^2} \rightarrow$ characteristic $\tau \propto q^{-2}$

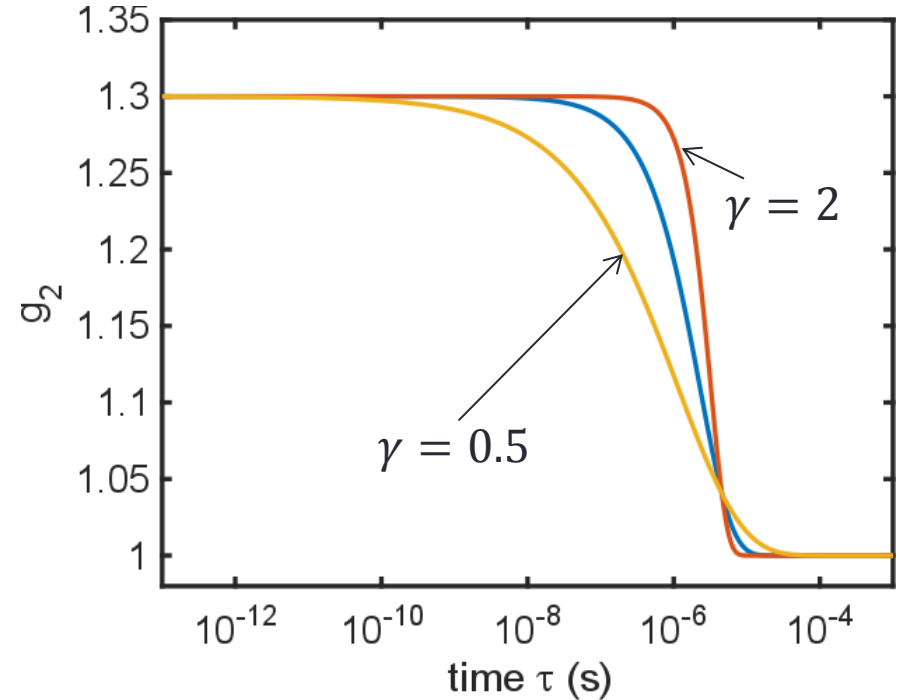
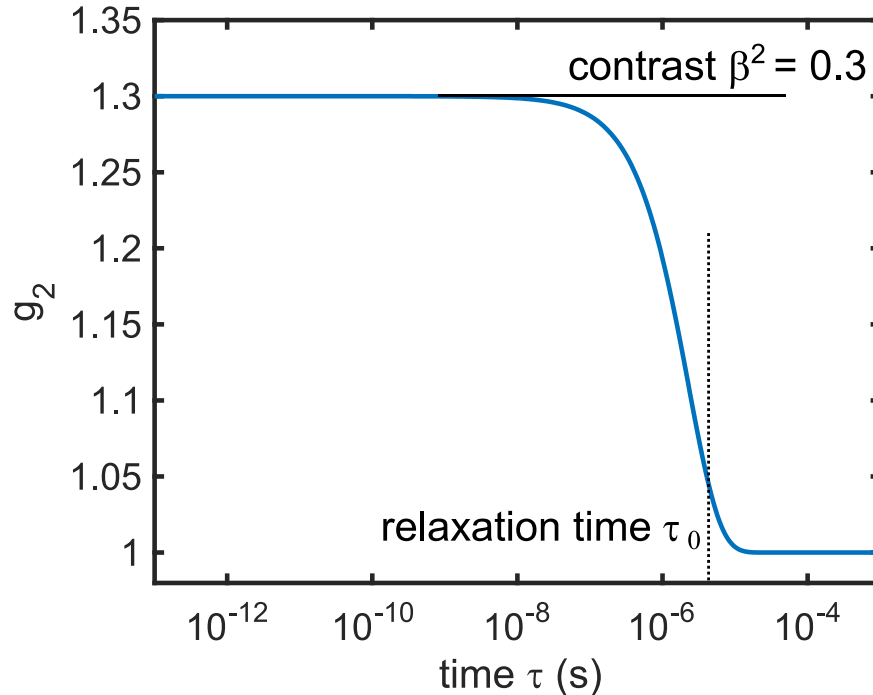
- Measuring g_2 allows to obtain particle size $r = \frac{k_B T \tau_0 q^2}{6\pi\eta}$ when solvent properties are known \rightarrow Dynamical light scattering

- On the other hand, known particles can be used to probe solvent properties, in particular viscosity $\eta \rightarrow$ microrheology



XPCS – correlation functions

Spherical particles with $R = 10$ nm in water



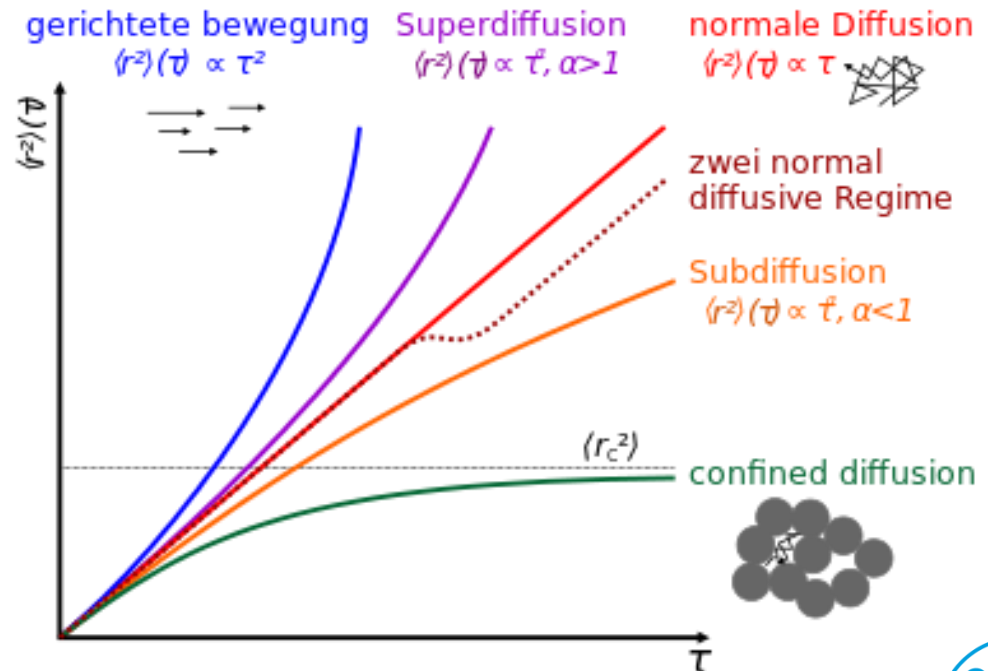
- Stretched and compressed correlation functions: Kohlrausch-Williams-Watts function $f(q, \tau) = \exp(-(\Gamma\tau)^\gamma)$
- Measure of width of distribution of (local) relaxation times

XPCS – correlation functions

Q-dependencies

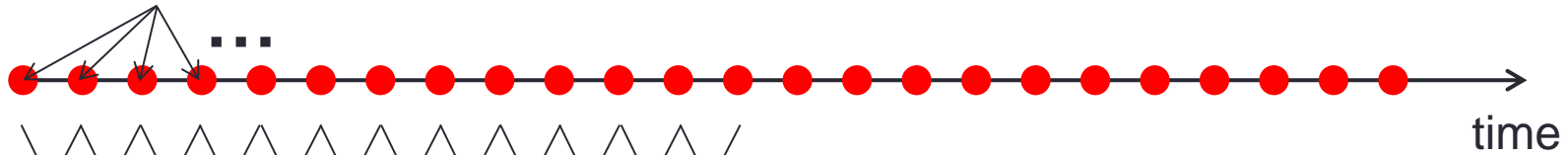
- Diffusion: $\tau_0 \propto q^{-2}$, $\langle r^2 \rangle \propto \tau$
- More general ($\tau_0 \propto q^{-\delta}$): $\langle r^2 \rangle \propto \tau^\alpha \rightsquigarrow r \propto \tau^{\frac{\alpha}{2}} \rightsquigarrow \tau \propto q^{\frac{2}{\alpha}}$ $\Rightarrow \alpha\delta = -2$
- Diffusion: $\delta = 2$
- Subdiffusion: $\delta > 2$
- Superdiffusion: $\delta < 2$
- Ballistic motion: $\delta = 1$

→ The analysis of both exponents γ and δ provides information on the type of dynamics



XPCS – instantaneous correlation function

Measured speckle patterns

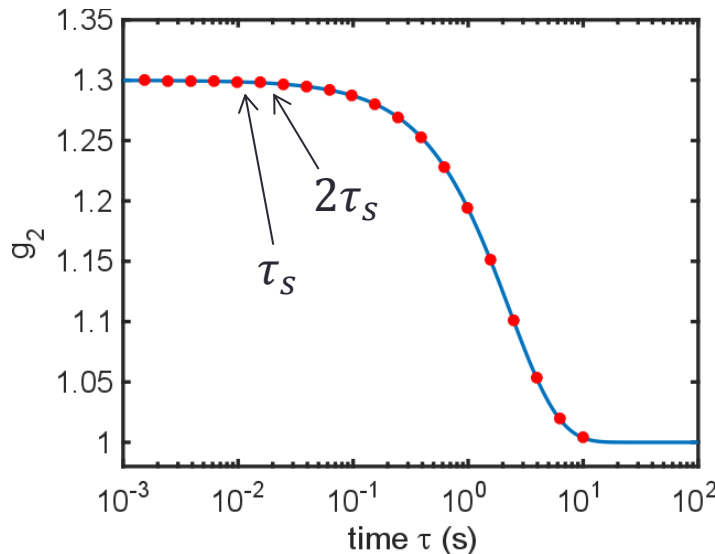


Correlate & average → Shortest lag time τ_S



Correlate & average → 2nd shortest lag time $2\tau_S$

⋮



→ $g_2(q, \tau)$ averaged over pairs of same lag time τ (...and pixels!) taken during the experimental run

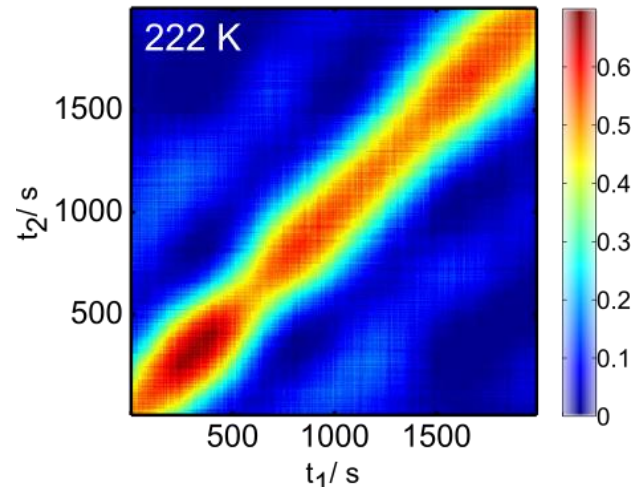
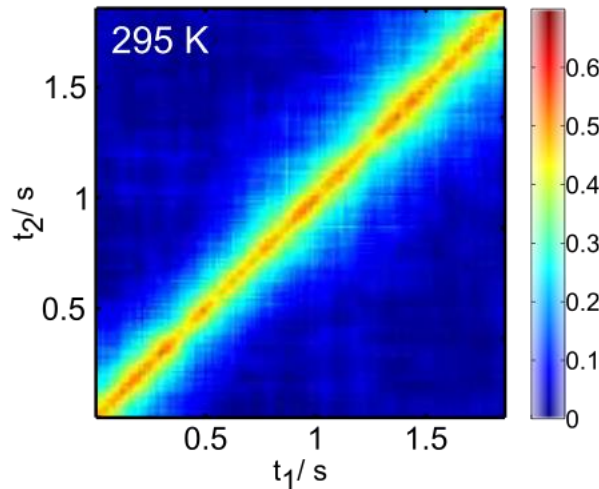
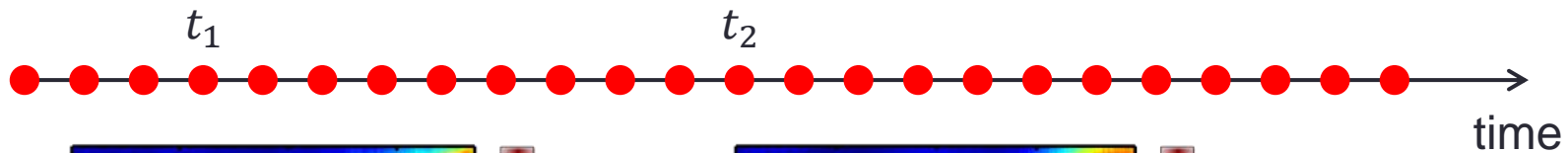


XPCS – instantaneous correlation function

But: sample may change during the measurement

→ Two-time correlation function $C_I(q, t_1, t_2) = \frac{\langle I(q, t_1) I(q, t_2) \rangle}{\langle I(q, t_1) \rangle \langle I(q, t_2) \rangle}$

- t_1, t_2 are points in experiment time, e.g.:



→ Dynamical heterogeneities!

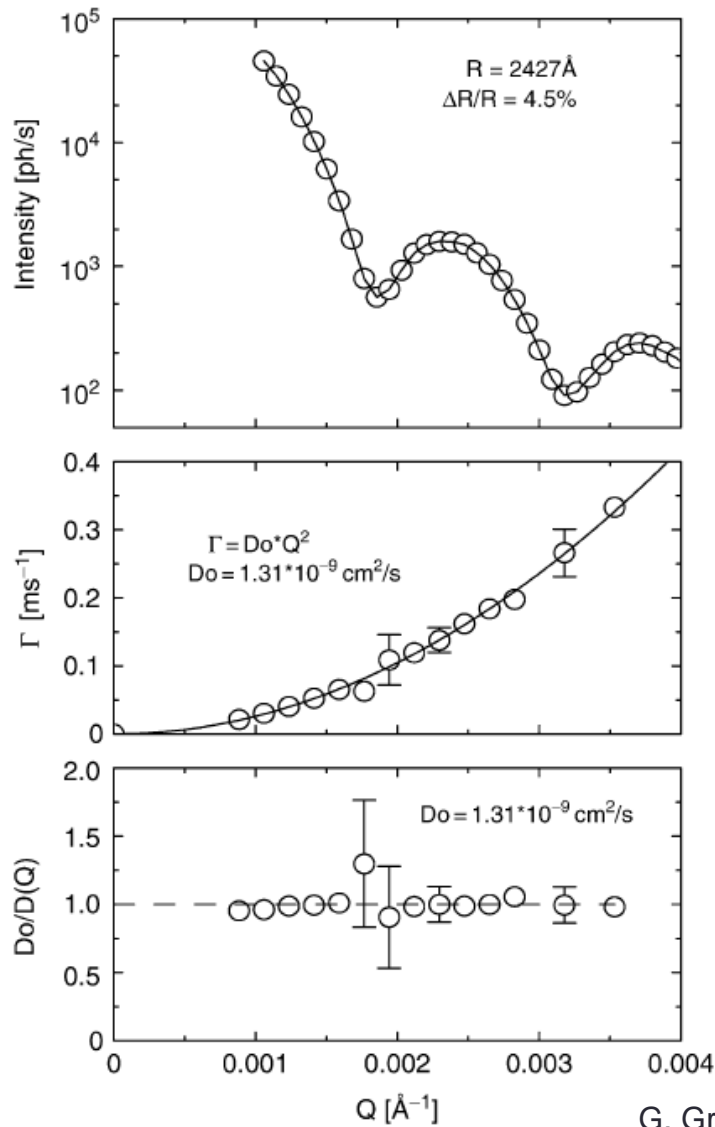
- Sample ages along $t_1 = t_2$ diagonal $t_{age} = \frac{t_2 - t_1}{2}$
- Lag time $\tau = |t_2 - t_1|$



XPCS example 1 – dynamics in colloidal systems

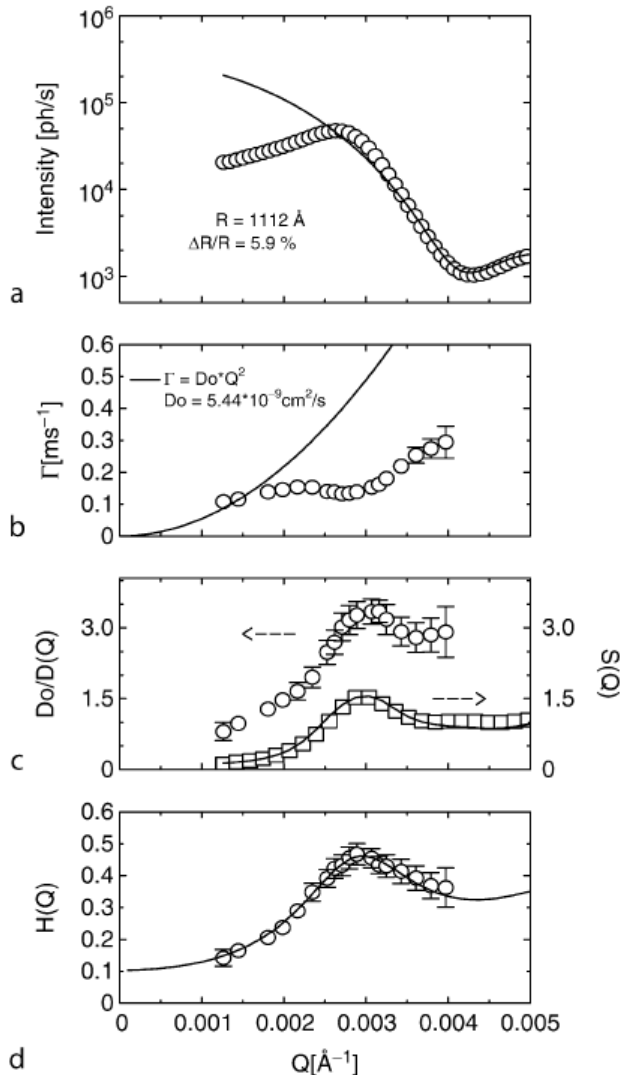
SiO₂ colloids in glycerol/water

- Low concentration: volume fraction 1%
- SAXS: Formfactor
- XPCS: diffusion with $\Gamma \propto q^2$



G. Grübel et al. In "Soft Matter Characterization", Springer (2008)

XPCS example 1 – dynamics in colloidal systems



PMMA particles in decalin

- High concentration: volume fraction 37%
- SAXS: Structure factor
- XPCS results deviate from $\Gamma \propto q^2$
- Effective diffusion constant $D(q) = D_0 H(q)/S(q)$ for short times, hydrodynamic function $H(q)$
- $H(q) = 1 \Rightarrow D(q) = D_0/S(q)$: de Gennes narrowing, i.e. slowing down around next-neighbour distances.

G. Grübel et al. In "Soft Matter Characterization", Springer (2008)

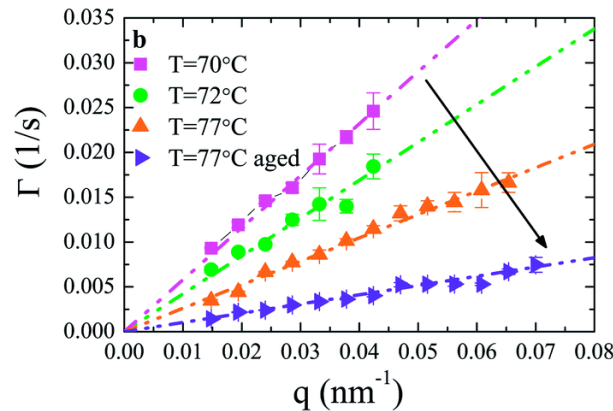
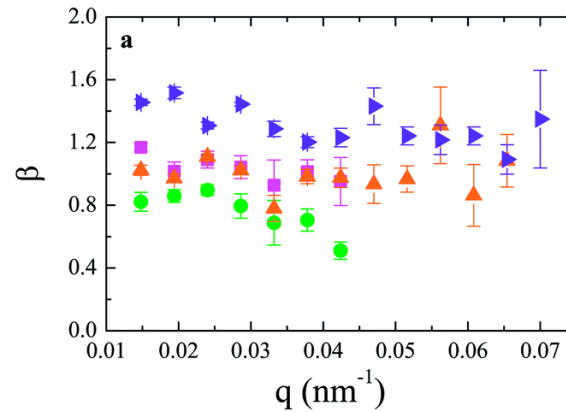
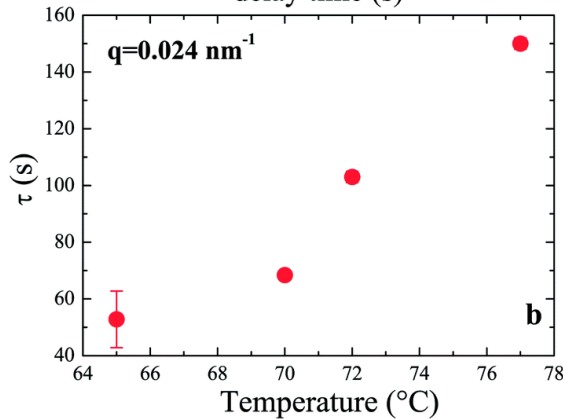
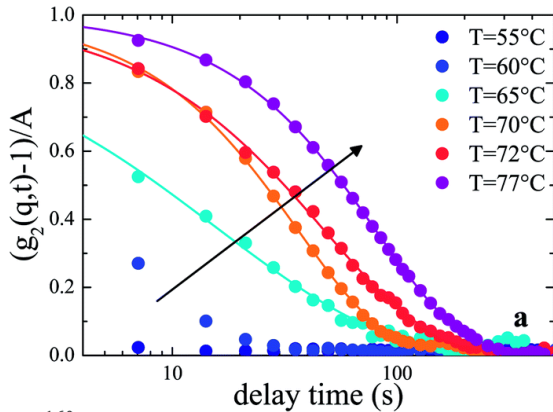
XPCS example 2 – microrheology

Tracer particles to measure solvent properties

- Weak scattering signal from solvent
- Large q -region has to be probed (low speckle contrast)
- Slower dynamics in SAXS regime
- Indirect access to solvent properties only
- Length scale of several 10 nm given by the tracer particle size
- Low tracer particle concentrations, so that $S(q) = 1$ for the particles \rightarrow avoid any particle-solvent interactions



XPCS example 2 – microrheology



B. Ruta et al. *Soft Matter* 10, 4547 (2014)

Glass transition studies: → Lecture 16

- SiO₂ particles as tracer for gelation of methylcellulose in water
- Gel-gel-transition: Turbid gel for $T \geq 60$ °C
- Stretched ($\beta < 1$) to compressed ($\beta > 1$) transition (KWW exponent!)
- Hyper-diffusive & compressed at high temperatures → stress-dominated

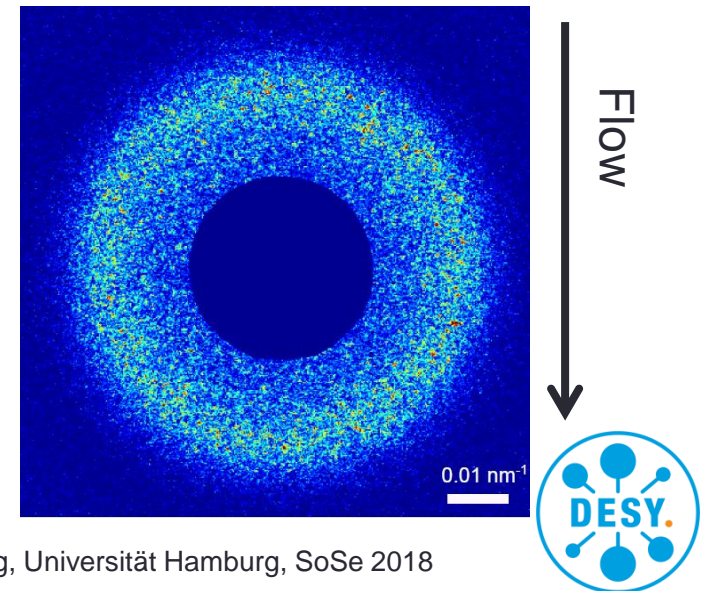
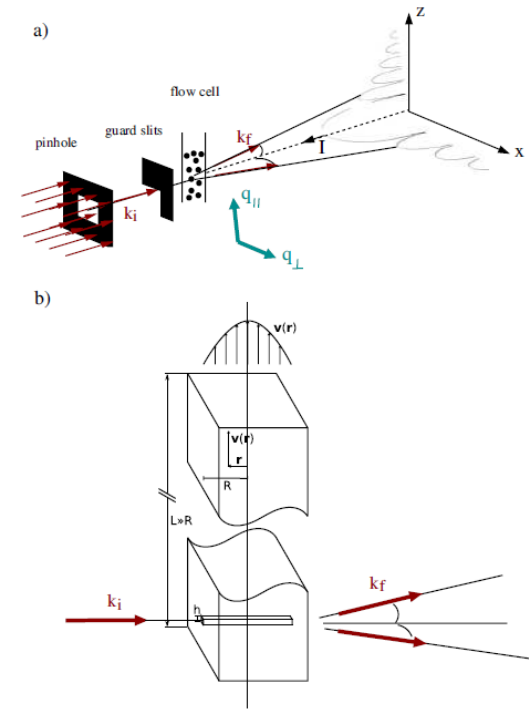
XPCS example 3 – directed motion

Sample undergoes (shear) flow

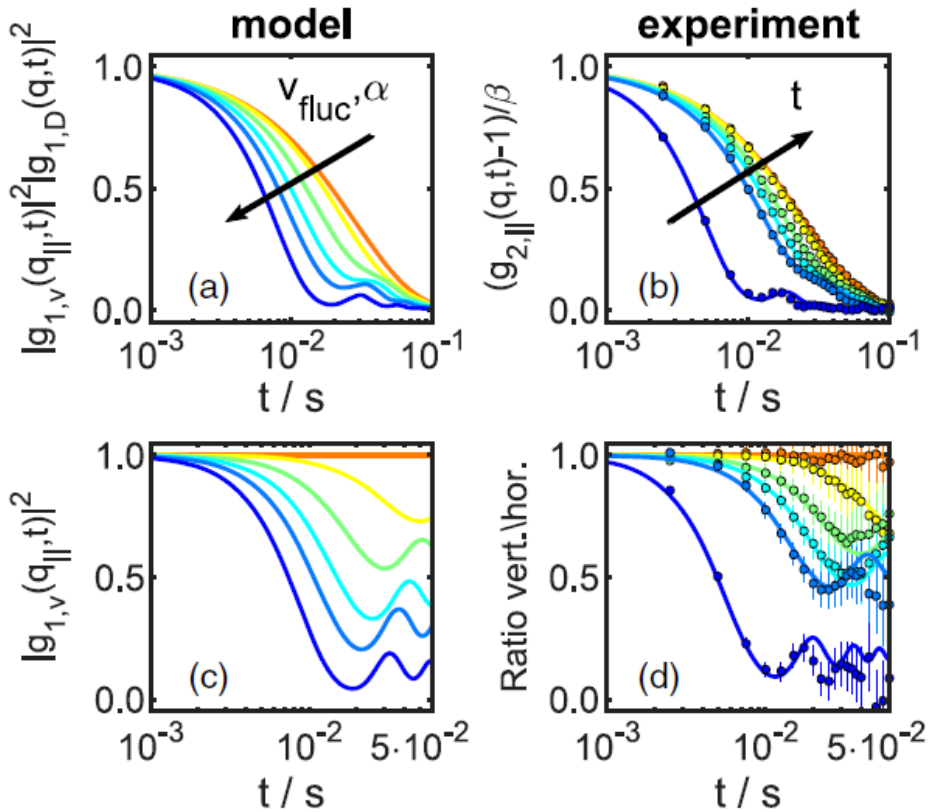
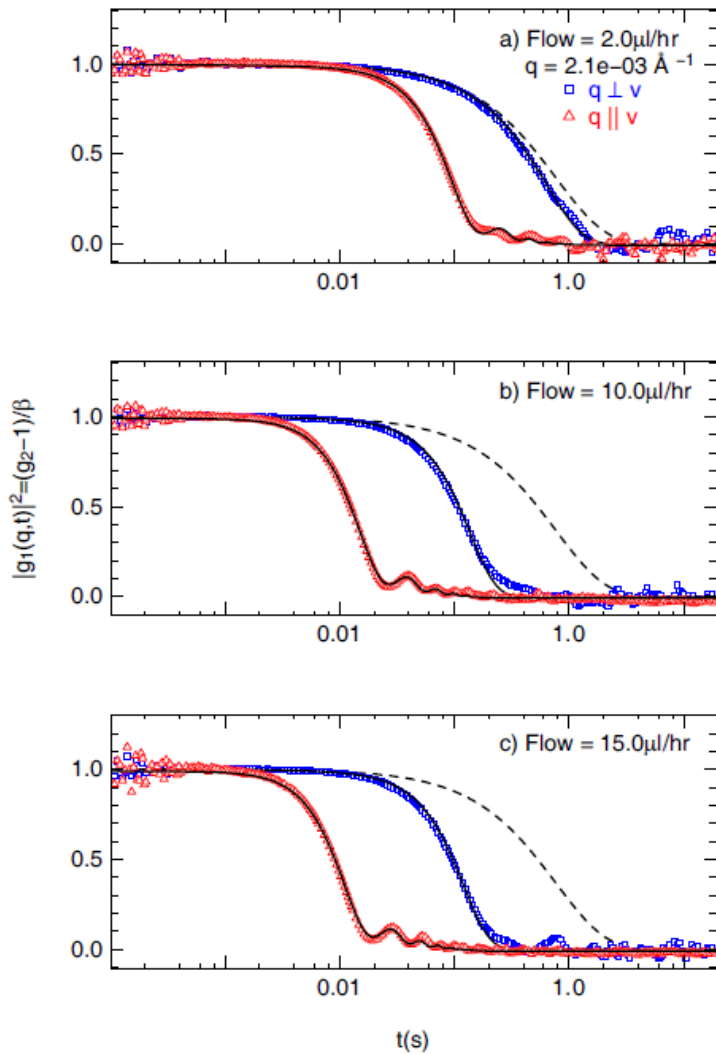
- Flowing sample to avoid radiation damage
- Sedimentation of particles
- $f(q, \tau)$: product of diffusive and advective contributions
 $f(q, \tau) = f_d(q, \tau) \cdot f_a(q, \tau) = \exp(-\Gamma t) \cdot f_a(q, \tau)$
- $f_a(q, \tau)$ can become complicated

Perpendicular to flow it simplifies to
 $g_{2,\perp}(q, \tau) = 1 + \beta^2 \exp(-2\Gamma\tau) \exp(-(\nu_{tr}t)^2)$,
 transit-induced frequency ν_{tr}

In flow direction: oscillating behaviour



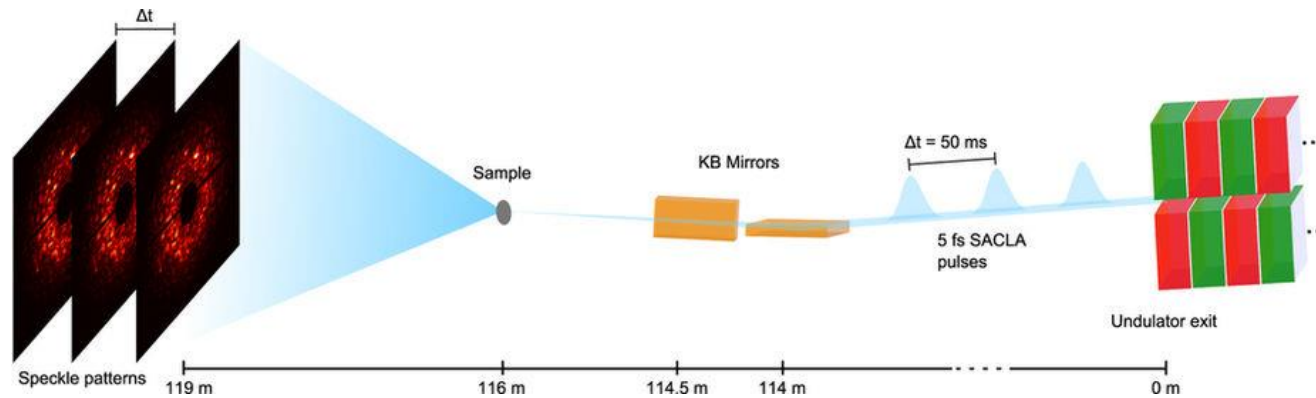
XPCS example 3 – directed motion



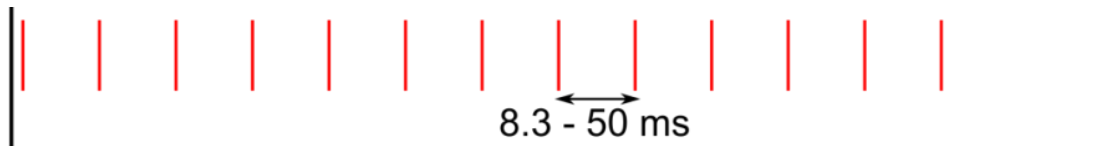
Particle sedimentation
 J. Möller et al. PRL 118, 198001 (2017)

S. Busch et al. Eur. Phys. J. E 26, 55 (2008)

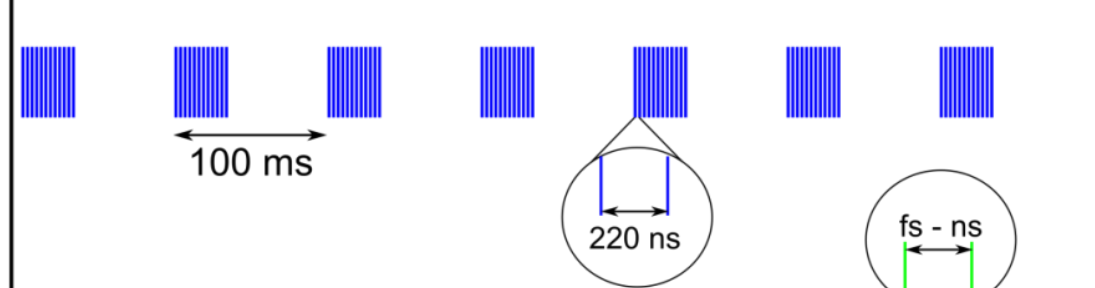
XPCS example 4 – XPCS at FEL



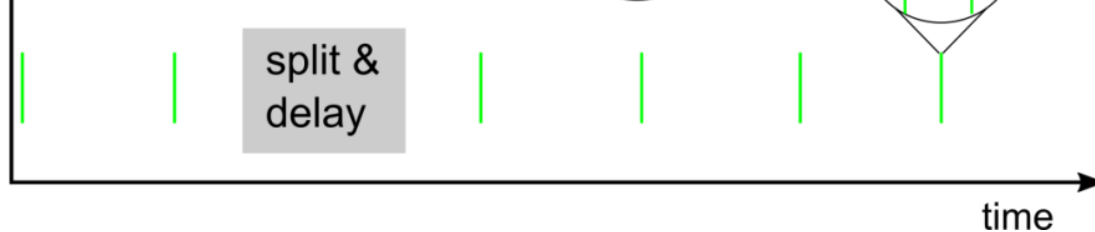
SACLA / LCLS



European XFEL



Split-pulse XPCS



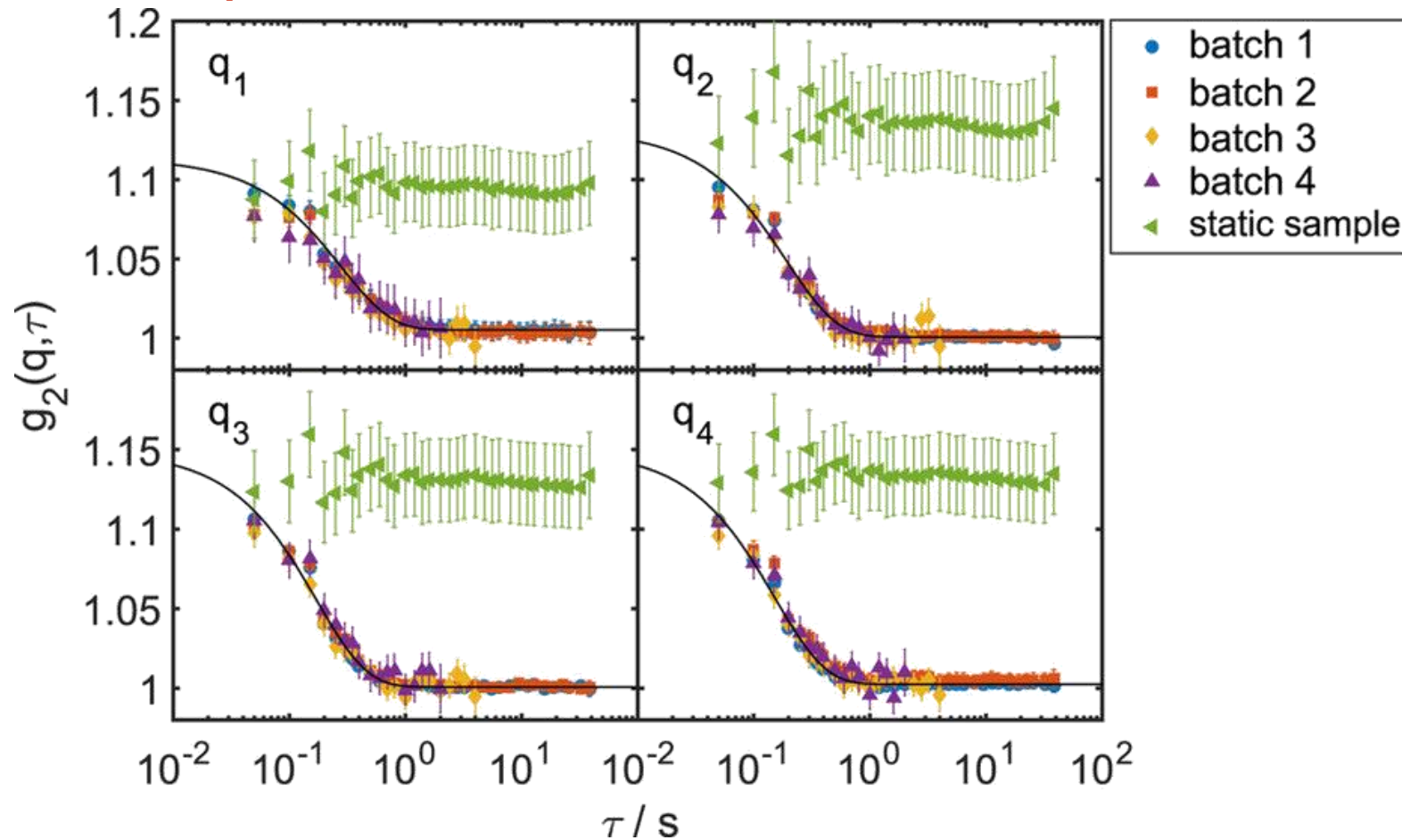
XPCS example 4 – XPCS at FEL

Parameter	Storage ring	FEL
Time Structure	Continuous	Pulses
Coherence	Partial	Full
Intensity	Stable	Fluctuations
Position / pointing	Stable	Fluctuations
Energy spectrum	Stable	Fluctuations
Time lag	Detector- and flux-limited ($\geq 10^{-6}$ s)	Repetition rate

State-of-the-art detectors at storage rings

Maxipix (ESRF) ~ 300 Hz
 Lambda (PETRA III, APS) ~ >2 kHz
 Eiger (PETRA III, NSLS-II, ESRF) ~ kHz

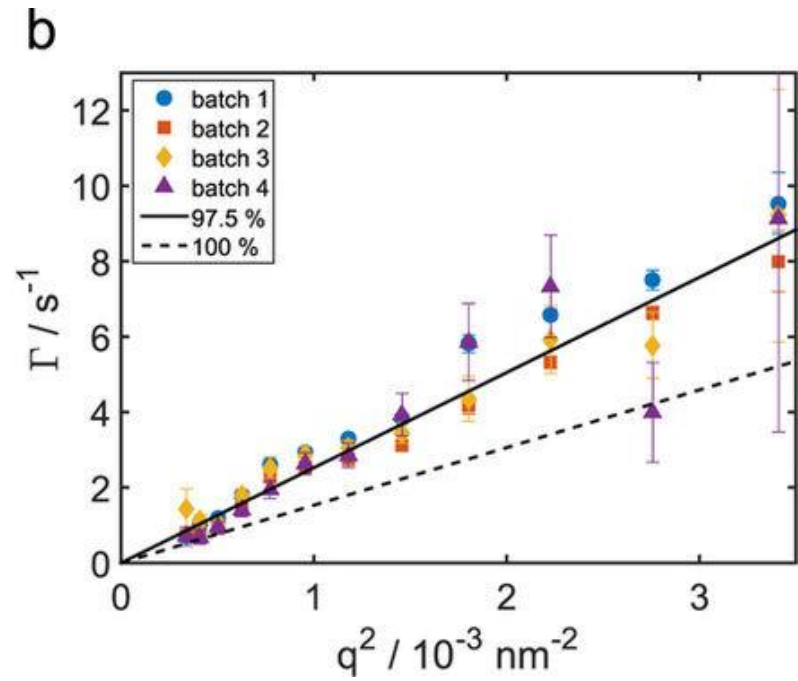
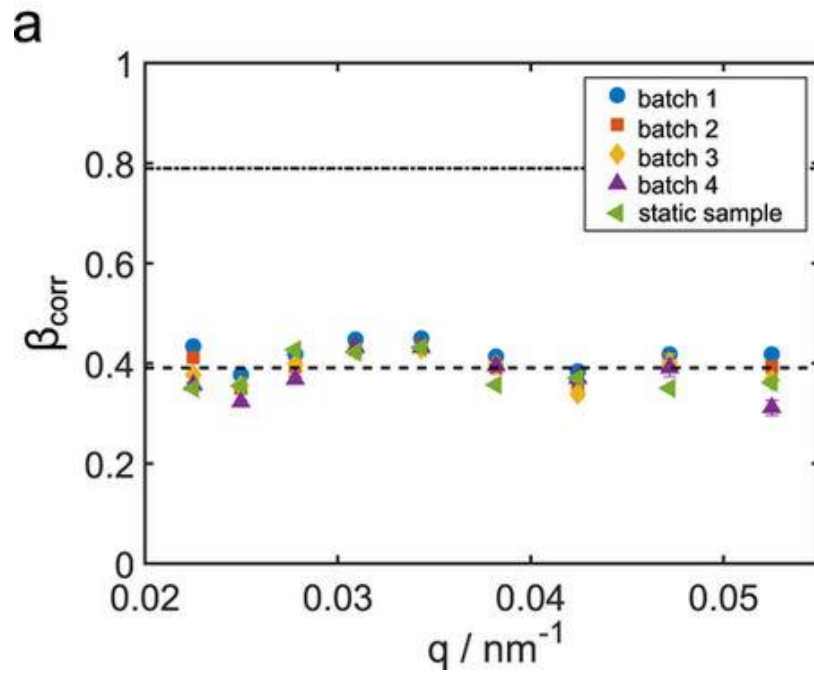
XPCS example 4 – XPCS at FEL



Diffusion of nanoparticles in glycerol and static sample with 20 Hz rep. Rate

FL et al. Sci. Rep. 5, 17193 (2015)

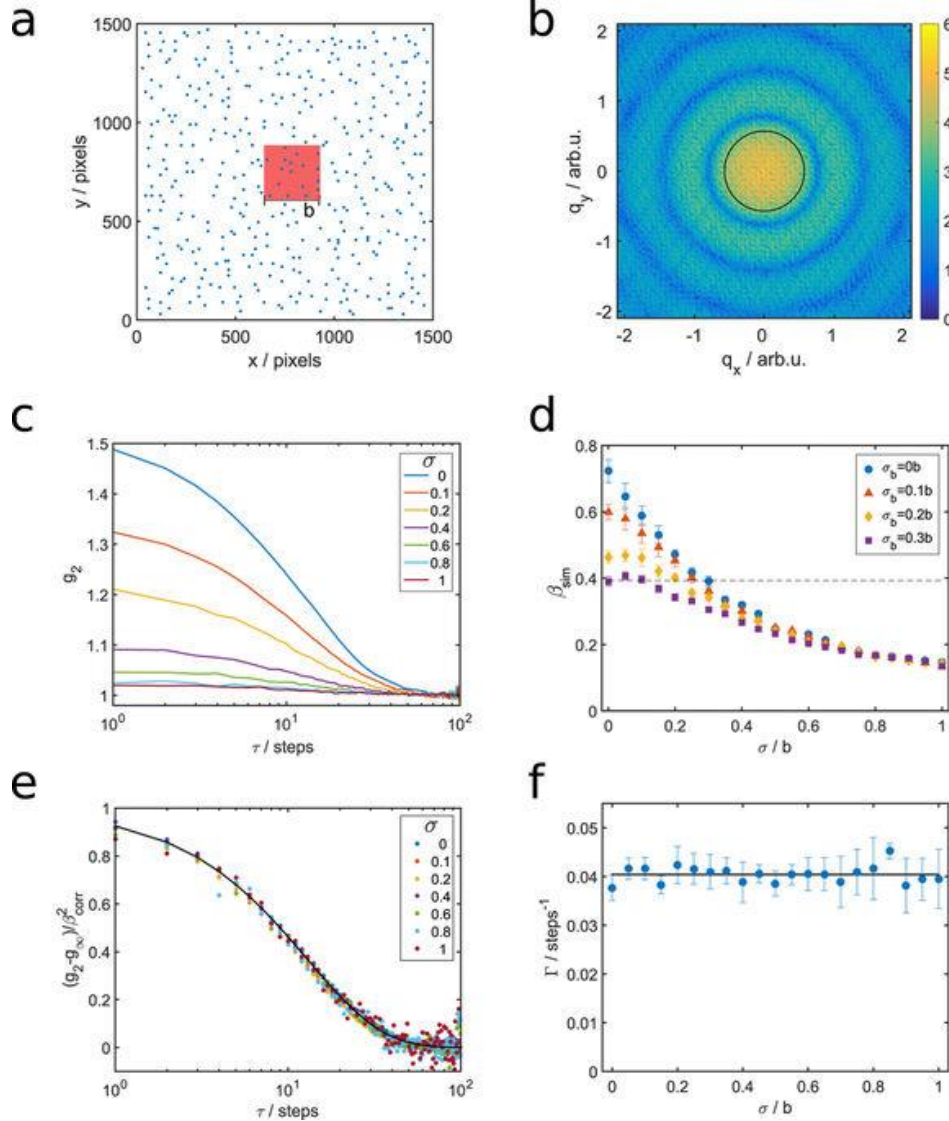
XPCS example 4 – XPCS at FEL



- Contrast well below the expected value
- Extracted relaxation rates well-described by diffusion

FL et al. Sci. Rep. 5, 17193 (2015)

XPCS example 4 – XPCS at FEL

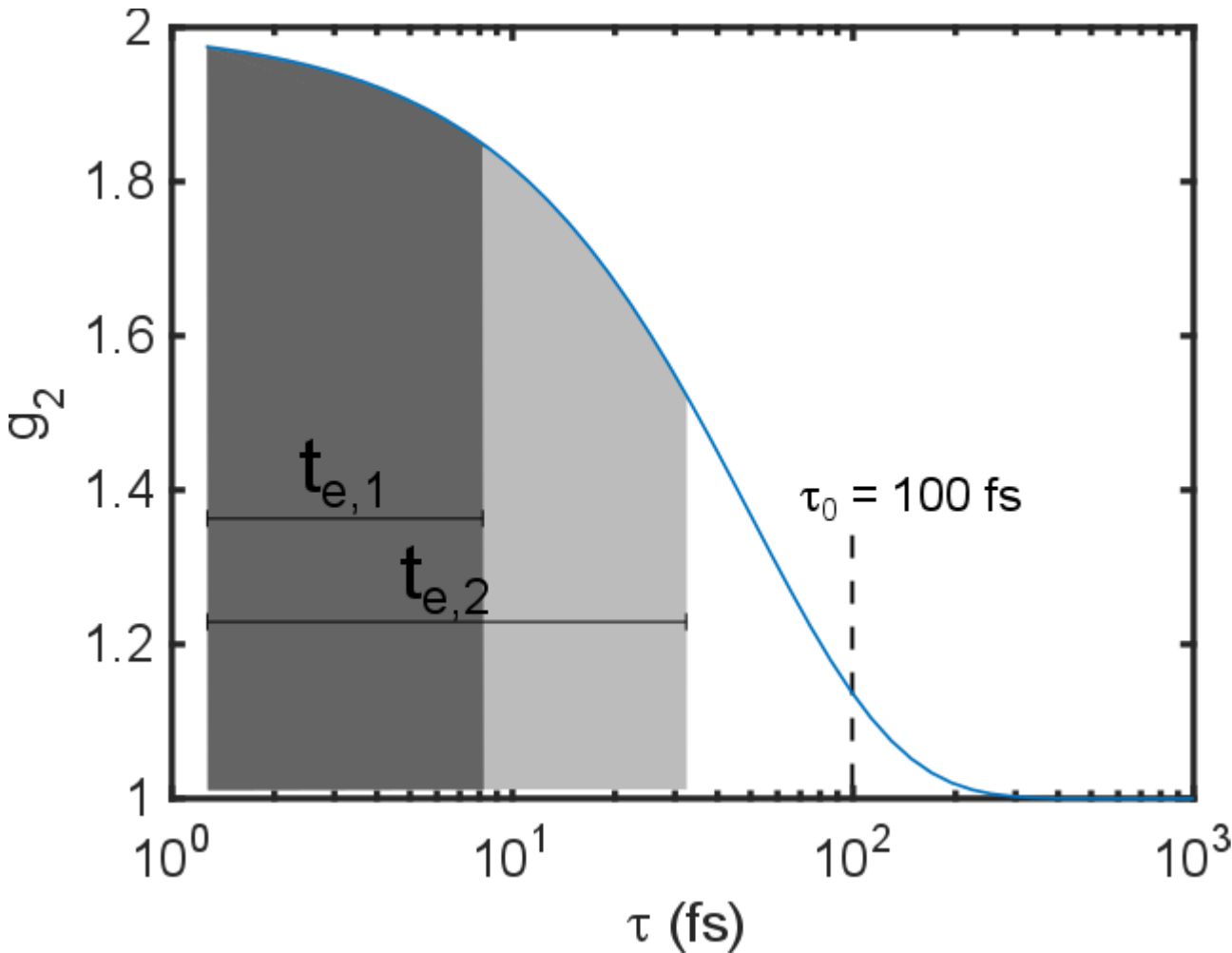


- XPCS simulations
- Moving beam & beam size modifications on shot-to-shot basis
- Extracted relaxation rates well-described by diffusion
- Drop of effective contrast

FL et al. Sci. Rep. 5, 17193 (2015)



XPCS example 5 – ultrafast XPCS: XSVS



- So far: short exposure time
 $t_e \ll \tau_0$
- Dynamics change during exposure \rightarrow reduction of contrast



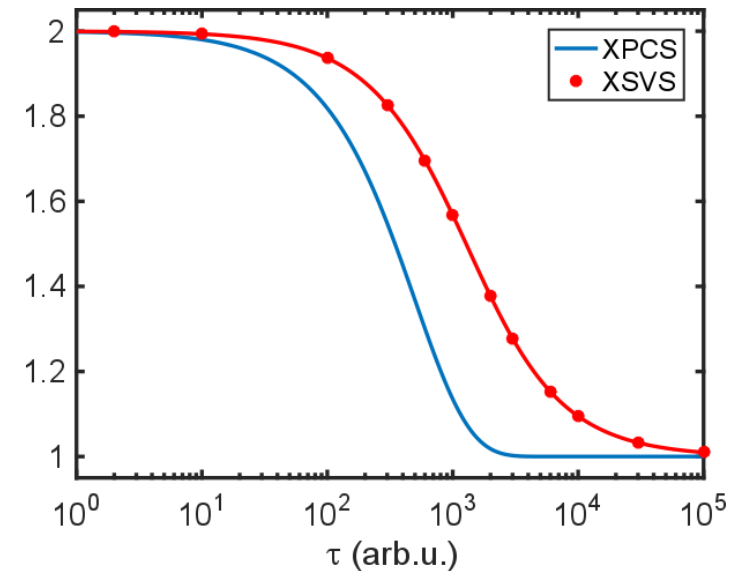
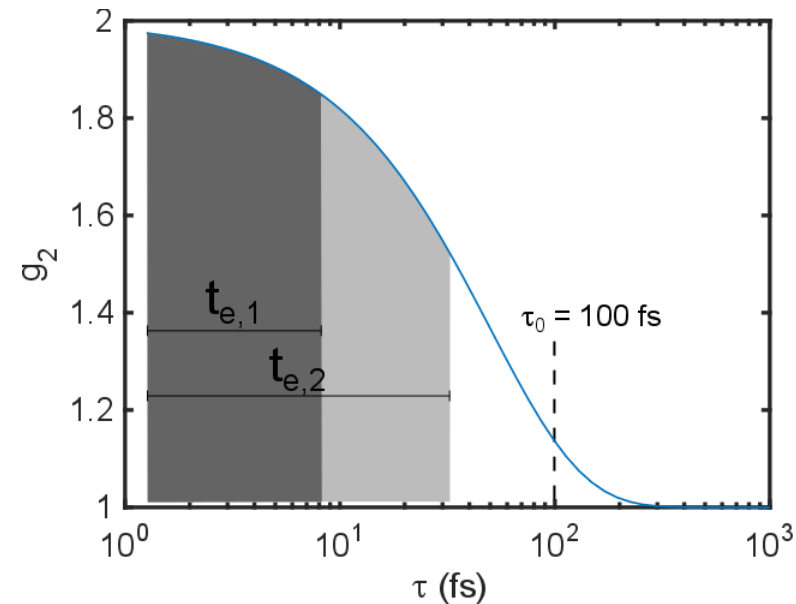
XPCS example 5 – ultrafast XPCS: XSVS

Finite pulse lengths

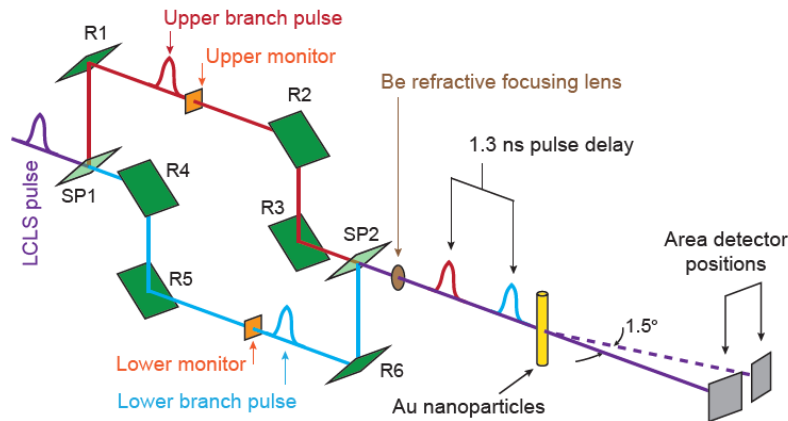
- Contrast as function of exposure ($\beta^2 = \frac{\langle I^2 \rangle}{\langle I \rangle^2} - 1$)

$$\beta^2(q, t_e) = \frac{2\beta_0^2}{t_e} \int_0^{t_e} \left(1 - \frac{\tau}{t_e}\right) |f(q, \tau)|^2 d\tau$$

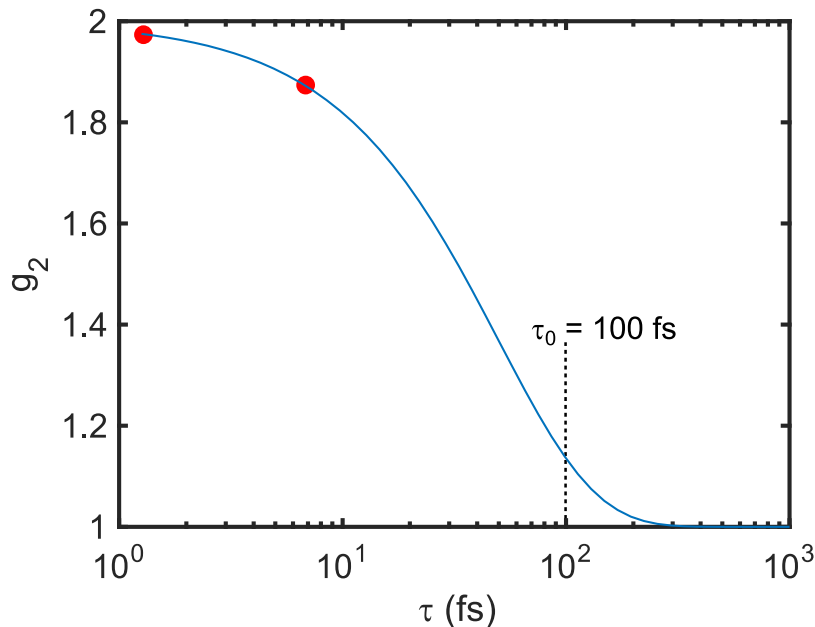
- For diffusion, this can be solved analytically (→ Exercise!)
- Limited by accessible exposure times
 - Pulse lengths (FEL)
 - Detector read out & flux (storage rings)
- FEL: pulse lengths variations & split-pulse applications



XPCS example 5 – ultrafast XPCS: double shot



- Split FEL pulse in two and delay one of them
- Speckle pattern: sum of two patterns
- XSVS-type of analysis
- Typically low count rates → obtain contrast from distribution functions of intensity



W. Roseker et al. Nature Comm. 9, 1704 (2018)



XPCS: further examples

- Glass dynamics and glass transition (→ lecture 16)
- In general, large q experiments
 - Water (→ lecture 17)
 - Domain wall dynamics
 - Network glasses (SiO_2 , ...)
 - ...
- Liquid surfaces: capillary waves → oscillating ISF
- ...
- Relation to neutron spin echo and dynamic light scattering



XPCS – signal to noise

- Reasonable definition $SNR = \frac{g_2 - 1}{\sqrt{\text{var}(g_2)}}$
 - $\text{var}(g_2) = \frac{g_2}{N_p \langle n_c \rangle} = \frac{g_2}{n_x n_y T t_a I^2}$
 - N_p number of correlated pairs averaged for g_2
 - $\langle n_c \rangle$ mean number of counts per exposure time
 - Using count rate I per pixel, accumulation time t_a , number of pixels P , total experimental duration T .
 - $SNR = \frac{g_2 - 1}{\sqrt{\text{var}(g_2)}} = \sqrt{PTt_a/g_2} \cdot I(g_2 - 1)$
 - Substitute g_2 with the limit of $\tau \rightarrow 0$: $g_2 = \beta^2 + 1$
 - Low contrast limit: $\sqrt{g_2} \approx 1$
- $SNR = \beta^2 I \sqrt{PTt_a}$
- Consequence: Increase of I by 10 → accessible time scale 100x smaller!

see P. Falus et al. J. Synchr. Rad. 13, 253 (2006)



Accessible timescales for XPCS



Sequential XPCS at storage ring and FEL sources: $\tau \gtrsim 0.1$ ms



XSVS at FEL sources: $\tau \sim$ pulse lengths ~ 0 -100 fs



Split-pulse XPCS: $\tau \leq 1$ ns



Sequential XPCS at European XFEL: $600 \mu\text{s} \geq \tau \geq 220$ ns



XPCS at DLSR: $\tau \leq$ sub- μ s

