

# Methoden moderner Röntgenphysik: Streuung und Abbildung

Lecture 14	Vorlesung zum Haupt- oder Masterstudiengang Physik, SoSe 2018 G. Grübel, A. Philippi-Kobs, <u>F. Lehmkühler</u> , O. Seeck, L. Frenzel, M. Martins, W. Wurth		
Location	Lecture hall AP, Physics, Jungiusstraße		
Date	Tuesday Thursday	13:00 - 14:30 8:30 - 10:00	(starting 3.4.) (until 12.7.)





# **Soft Matter – Timeline**

- Di 29.05.2018 Soft Matter studies I: Methods & experiments Definitions, complex liquids, colloids, storage ring and FEL experiments, setups, liquid jets, ...
- Do 30.05.2018 Soft Matter studies II: Structure
   SAXS & WAXS applications, X-ray cross correlations, ...
- Di 05.06.2018 Soft Matter studies III: Dynamics
   XPCS applications, diffusion, dynamical heterogeneities, ...
- Do 07.06.2018 cancelled!
- Di 12.06.2018 Case study I: Glass transition at DESY campus! Supercooled liquids, glasses vs. crystals, glass transition concepts, structure-dynamics relations, ... + DESY photon science site visit
- Do 14.06.2018 Case study II: Water
   Phase diagram, anomalies, crystalline and glassy forms, FEL
   studies, ...

DESY.

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## **Small-angle X-ray scattering**



Web of knowledge topic search: "Small angle X-ray scattering"





# SAXS – Analysis methods: Formfactor

Lecture 7: 
$$I_{SAXS}(Q) = (\rho_{sI,p} - \rho_{sI,0})^2 \left| \int_{V_p} e^{iQr} dV_p \right|^2$$
 for particle (p) in solvent (0)

**Diluted case: Formfactors** 

- Spheres:  $F(q) = 3 \frac{\sin(qR) qR\cos(qR)}{(qR)^3}$
- In general difficult to calculate  $\rightarrow$  numerical approaches
- Soft Matter: Polydispersity & (solvent) background



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# **SAXS – Analysis methods: Formfactor**

Ab initio methods (use "dummy" bead models) → BioSAXS







# **SAXS – Analysis methods: Structure factors**

From lecture 5 (Kinematical Diffraction I): Structure factor of a liquid (or glass)

$$S(q) = 1 + \rho_0 \int_0^\infty \frac{4\pi r}{q} [g(r) - 1] \sin(qr) dr$$

With the radial pair correlation function g(r). This relates to the potential of mean force between two particles  $U_{MF}(r)$ 

$$g(r) = \exp\left(-\frac{U_{MF}(r)}{k_B T}\right)$$

For very dilute systems  $U_{MF}(r)$  equals the interaction potential U(r).

Relation of S(q) or g(r) and  $U(r) \rightarrow$  **Ornstein-Zernike equation** relating total correlations  $h(r) \equiv g(r) - 1$  to direct two-particle correlations c(r) and indirect correlations  $c(|\mathbf{r} - \mathbf{r}'|)$ 

$$h(r) = c(r) + \rho_0 \int c(|\mathbf{r} - \mathbf{r}'|)h(|\mathbf{r}'|)d\mathbf{r}'$$





# **SAXS – Analysis methods: Structure factors**

$$h(r) = c(r) + \rho_0 \int c(|\mathbf{r} - \mathbf{r}'|)h(|\mathbf{r}'|)d\mathbf{r}'$$

c(r) short range part

Can be solved using so-called "closure relations". Percus-Yevick closure:

$$c(r) = g(r) \left[ 1 - \exp\left(\frac{U(r)}{k_B T}\right) \right]$$

- → solves the hard-sphere potential  $U_{HS}(r) = \begin{cases} \infty, r \leq 2R \\ 0, r > 2R \end{cases}$  analytically.
- → Mean-spherical approximation closure relation  $c(r) = -\frac{U_{ES}(r)}{k_B T}$  solves electrostatic interactions (DLVO) [→ Lecture 13]









# Hard spheres

- Volume fraction as only parameter
- Does not include crystallisation/glass transition!
- I.e. typically breaks down close to  $\Phi \approx 0.5$

# **Sticky hard spheres**

$$\frac{U_{SHS}(r)}{k_B T} = \begin{cases} 0, \\ \ln\left(\frac{12\tau\Delta}{\sigma+\Delta}\right) \\ 0, \end{cases}$$

 $r < \sigma$ 

$$\sigma \le r \le \sigma + \Delta$$





#### **Structure factors – RMSA**

Charge stabilized systems  $\rightarrow$  rescaled mean spherical approximation (RMSA)

Structure factor as function of  $\Phi$ , charge, screening

High screening  $\rightarrow$  hard spheres







#### **Example 1: Structure and Formfactors from charge stabilized colloids**





PMMA spheres in water

Westermeier et al. JCP 137, 114504 (2012)





# **Example 2: High pressure studies**

- Structure at high pressures → solid sample chambers (diamond windows)
- X-rays to penetrate diamond windows
- Functionalized core-shell particles at pressures up to 4 kbar: transition from repulsion to attraction (sticky hard spheres!)



a)

b)

1 bar

500 bar 1000 bar

0



# **Example 3: nucleation and growth of quantum dots**







Combination of SAXS & XRD

→ cristallinity of nanoparticle





# **Example 4: Phase transitions in liquid crystals**











Isotropic

Nematic

Smectic

Goethite [ $\alpha$ -FeO(OH)] particles in water may form

- Isotropic •
- Nematic
- Smectic

Phases  $\rightarrow$  SAXS







#### **Example 4: Phase transitions in liquid crystals**

n



- **Disc-systems**
- (a) Discotic nematic phase
- (b) Hexagonal columnar phase
- (c) Rectangular columnar phase



Combined SAXS/WAXS from columnar phase

- SAXS: hexagonal intercolumnar order
- WAXS: disorder inside column

de Jeu: "Basic X-ray scattering for Soft Matter", 2016



# **Example 4: Phase transitions in liquid crystals**



Liquid crystal phase of the system monoglyceride-water

de Jeu: "Basic X-ray scattering for Soft Matter", 2016





# **Further methods and applications**

- Anomalous SAXS → ASAXS
- Scanning SAXS
- Phase transitions and self-assembly
- Time resolved techniques
- SAXS tomography
- BioSAXS
- ...





SAXS: 1D information (typically)

 $\rightarrow$  How to make use of the 2D information obtained from a 2D scattering pattern?

 $\rightarrow$  Angular correlations



1D information (standard SAXS)

•  $I(\mathbf{q}) = \langle I(q, \varphi) \rangle_{\varphi} = I(q)$ 

2D information: Angular correlations

- $C(q, \Delta) = \frac{\langle I(q, \phi)I(q, \phi + \Delta) \rangle_{\phi} \langle I(q, \phi) \rangle_{\phi}^2}{\langle I(q, \phi) \rangle_{\phi}^2}$ , i.e. correlations of fluctuations
- Coherent X-rays
- Two possibilities:
  - Solve structures in solution
  - Hidden symmetries





# **Reminder: coherent X-rays**

- Correlations of speckles  $\rightarrow$  coherent X-rays
- Reminder: degree of coherence of partial coherent source  $\rightarrow$  speckle contrast  $\beta = \frac{\sigma^2}{\langle I \rangle^2} = \frac{\operatorname{var}(I)}{\langle I \rangle^2} \le 1$
- Intensity follows Gamma distribution (Lecture 10)
- Low intensities  $\rightarrow$  Poisson noise  $\rightarrow$  Negative binomial probability function

$$P_{nb}(i) = \frac{\Gamma(i+M)}{\Gamma(M)\Gamma(i+1)} \left(1 + \frac{M}{\langle i \rangle}\right)^{-i} \left(1 + \frac{\langle i \rangle}{M}\right)^{-M}, \text{ with number of modes } M = \frac{1}{\beta}$$



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# **Reminder: coherent X-rays**

Excercise: Contrast calculation at low count rates





Consider coherent X-ray scattering experiment in transmission geometry (e.g. SAXS) with 2D detector on disordered sample of N identical particles

$$A_{j}(\mathbf{q}) = \int \rho_{j}(\mathbf{r})e^{i\mathbf{q}\mathbf{r}}d\mathbf{r} \to I(\mathbf{q}) = \sum_{j_{1},j_{2}=1}^{N} e^{i\mathbf{q}\mathbf{R}(j_{1},j_{2})}A_{j_{1}}^{*}(\mathbf{q})A_{j_{2}}(\mathbf{q})$$
$$= \sum_{j_{1},j_{2}=1}^{N} \int \int \rho_{j_{1}}^{*}(\mathbf{r}_{1})\rho_{j_{2}}(\mathbf{r}_{2})e^{i\mathbf{q}(\mathbf{R}(j_{1},j_{2})+\mathbf{r}_{21})}d\mathbf{r}_{1}d\mathbf{r}_{2}$$

Partially coherent illumination and dilute system (particles distance > coherence length)  $\rightarrow$  interparticle correlations can be neglected:

$$I(\mathbf{q}) = \sum_{j=1}^{N} I_j(\mathbf{q}) = \sum_{j=1}^{N} |A_j(\mathbf{q})|^2$$

Angular information: Fourier decomposition

$$I(\mathbf{q}) = I(q,\phi) = \sum_{l=-\infty}^{\infty} \hat{I}_{\ell}(q) e^{il\phi}; \quad \hat{I}_{\ell}(q) = \frac{1}{2\pi} \int_{0}^{2\pi} I(q,\phi) e^{-i\ell\phi} \, \mathrm{d}\phi$$





Now consider 2D disordered system in the dilute limit, e.g. pentagonal arrangement of particles (polar coordinates,  $R_0$  radius of pentagon,  $\theta_j = \frac{2\pi j}{r}$ )

$$\rho(r,\theta) = \frac{\delta(r-R_0)}{R_0} \sum_{j=1}^5 \delta(\theta - \theta_j)$$

Expansion of scattering amplitude in Fourier series yields

$$A(q,\phi) = \sum_{\ell=-\infty} \hat{a}_{\ell}(q) e^{il\phi}$$
(1)

with Fourier coefficients

$$\hat{a}_{\ell}(q) = i^{-\ell} J_{\ell}(qR_0) \sum_{j=1}^{5} e_j^{il\theta_j}$$
(2)

- Pentagonal symmetry: only contribution if  $\ell = 0 \mod 5$  in (2).
- Odd terms cancel out pairwise (e.g.  $\ell = 5$  and  $\ell = -5$ ) in (1)  $\rightarrow$  Friedel's law!
- Only contributions with  $\ell = 0 \mod 10$
- $F_l(q) \propto J_\ell(qR_0) \rightarrow$  higher-order terms at large q





- Corresponding correlation function  $C(q, \Delta) = \frac{\langle I(q,\phi)I(q,\phi+\Delta)\rangle_{\phi} \langle I(q,\phi)\rangle_{\phi}^2}{\langle I(q,\phi)\rangle_{\phi}^2}$  with Fourier coefficients  $\hat{c}_{\ell}(q) = |\hat{I}_{\ell}(q)|^2$  (Wiener–Khinchin theorem)
- Correlations between different q possible



Adv. Chem. Phys. 161, 1 (2016)

• 3D systems: curvature of Ewald sphere  $\rightarrow$  odd symmetries











Hard-sphere glass





→ Hidden symmetries
 → Structural information beyond SAXS

PNAS 109, 11511 (2009)





# **XCCA & microscopy**



# Thin colloidal films

# Orientational order with 500 nm resolution



IUCrJ 5, 354 (2018)





# Liquid crystals

High number of symmetries  $\rightarrow$  strongly developed hexatic order

Measure of correlation length

XCCA to provide measure of degree of order and as order parameter for phase transitions



Adv. Chem. Phys. 161, 1 (2016)



#### **XCCA – sample reconstruction**







