

# Methoden moderner Röntgenphysik II

## Streuung und Abbildung

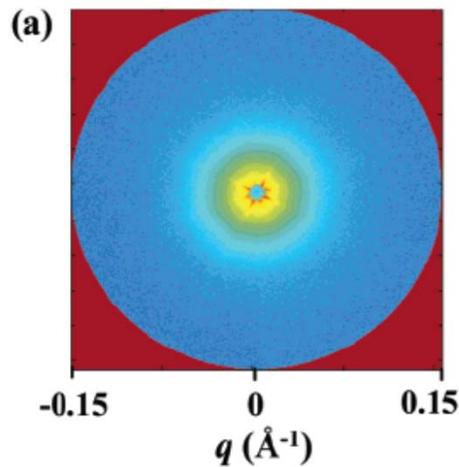
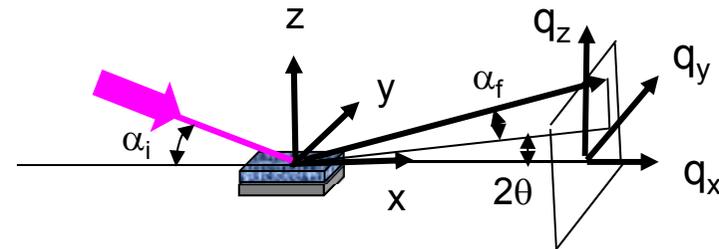
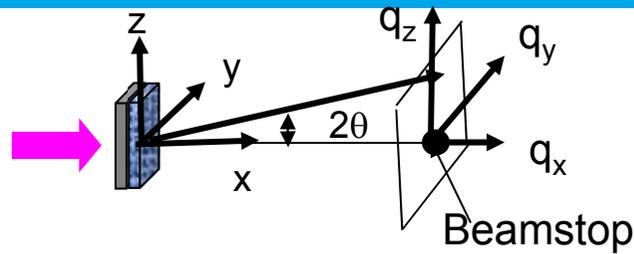
Stephan V. Roth  
DESY  
12.05.2015

# Outline

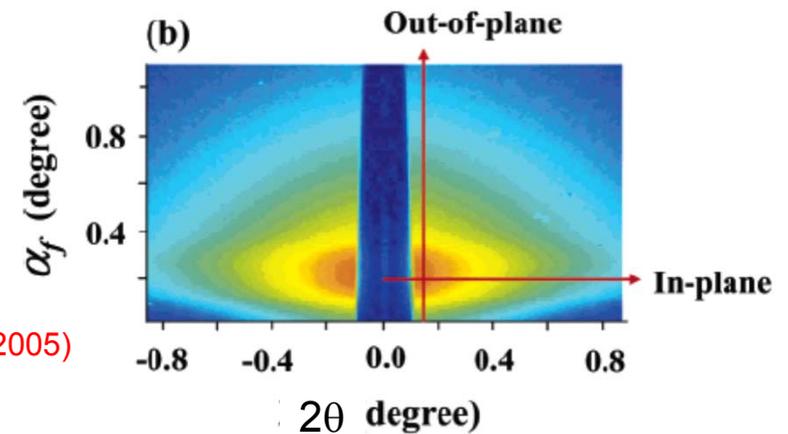
- > 12.05. : Small-Angle X-ray Scattering (SAXS)
- > 19.05. : Applications &  
A short excursion into Polymeric materials
- > 21.05. : Grazing incidence SAXS (GISAXS)
- > 02.06. : The polymer-metal interface – application of GISAXS  
On the route to organic electronics
- > 04.06.: In-situ studies of metal layer growth



# T-SAXS vs. GISAXS



Lee et al., *Macromolecules*, 38, 8991 (2005)



- Easy measurement
- Easy analysis
- In-plane information ( $q_y, q_z$ )
- Any possible scattering from substrate
- Transparency of substrate
- High energy

- Strong intensity
- Easy preparation of samples
- Full information ( $q_x, q_y, q_z$ )
- Scattering from surface / internal structure
- Scattering from reflected AND transmitted beam
- Refraction effects (DWBA)
- Special setup

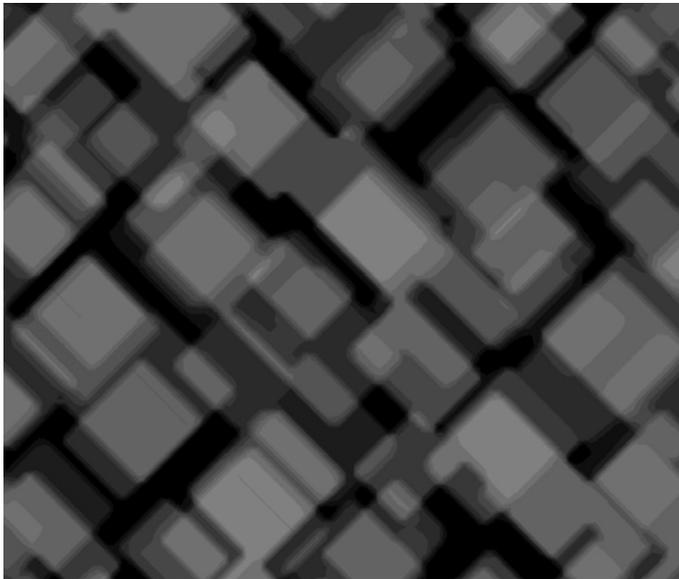


# Aim

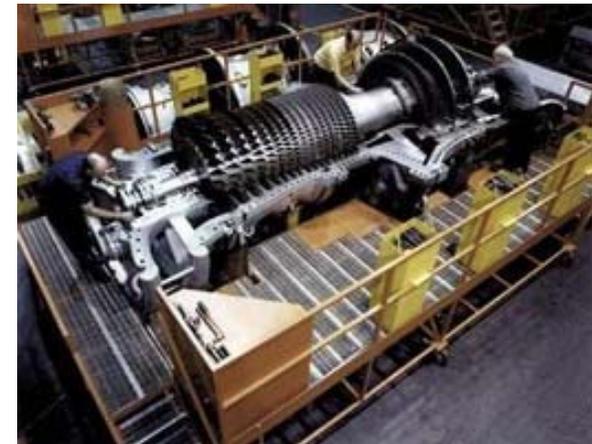
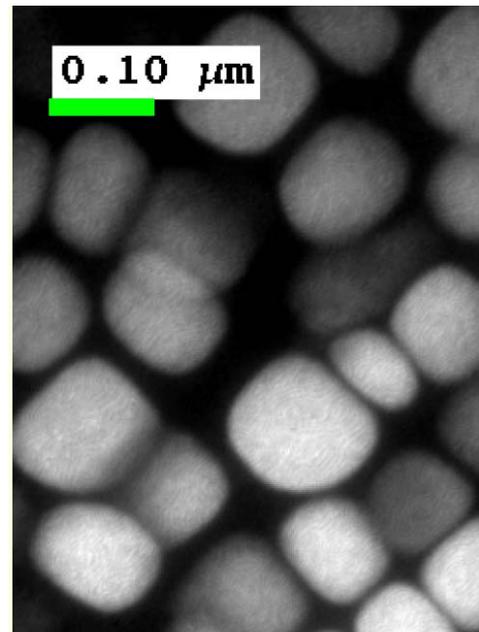
> To understand the **structure – property relation** of materials on **multiple length scales**

- **q-resolution**
- Maximum q-value
- **Beam size**

- Real pieces & materials
- Model systems
- Nanotechnology



Courtesy: R. Gilles (TUM)



<http://news.thomasnet.com/companystory/GE-Gas-Turbine-Technology-Selected-for-Pearl-GTL-Project-in-Qatar-495497>

# Outline II - today

## SAXS – Introduction

Neutrons, X-rays and Light: Scattering Methods Applied to Soft Condensed Matter.  
Eds: P. Lindner, Th. Zemb. North Holland Delta Series, Elsevier, Amsterdam (2002)  
ISBN: 0-444-51122-9

### > Instrumentation

- P03/MiNaXS @ PETRA III

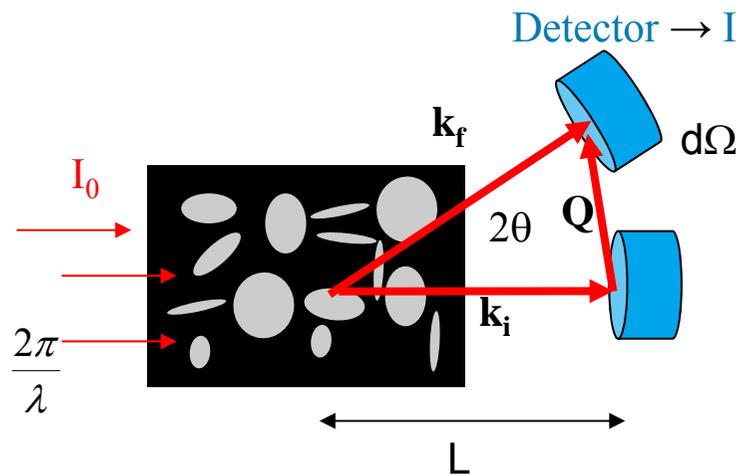
### > Bulk materials → Transmission U/SAXS:

- Porous materials
- Ni-base superalloys
- Droplet drying



# Cross-section

## > Differential cross section



$$d\sigma = \frac{I}{I_0} (L^2 d\Omega)$$

$$\frac{d\sigma}{d\Omega} = \frac{I}{I_0} (L^2) \Rightarrow \frac{d\Sigma}{d\Omega} = \frac{1}{V} \frac{d\sigma}{d\Omega}$$

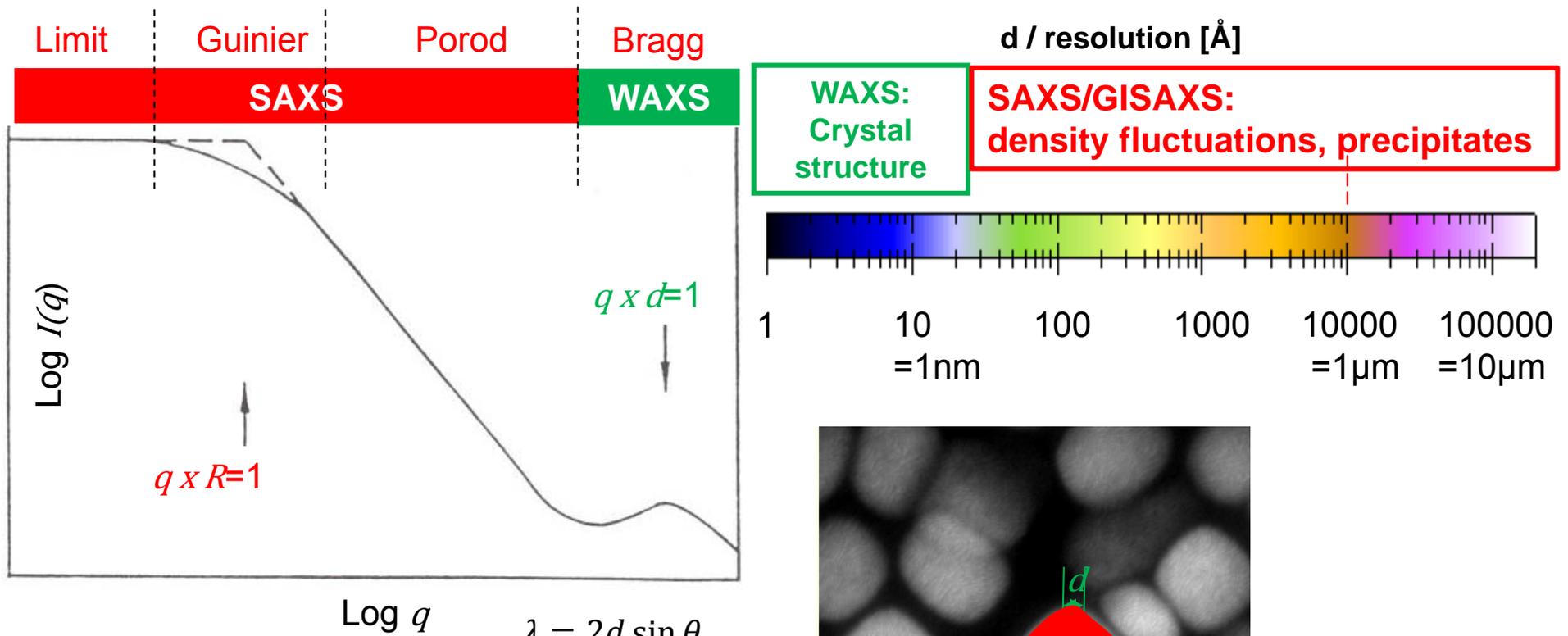
$V$  = Sample volume

$$\vec{q} = \vec{k}_f - \vec{k}_i \quad |\vec{k}_f| = |\vec{k}_i| = \frac{2\pi}{\lambda} \quad |\vec{q}| = 2 \frac{2\pi}{\lambda} \sin(\theta)$$

## > Scattering occurs due to density differences

# WAXS, SAXS, GISAXS...

Source: Streumethoden zur Untersuchung kondensierter Materie  
1996; ISBN 978-3-89336-180-9

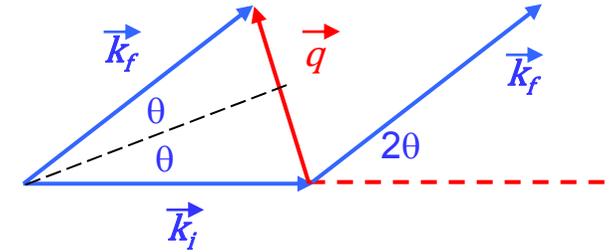
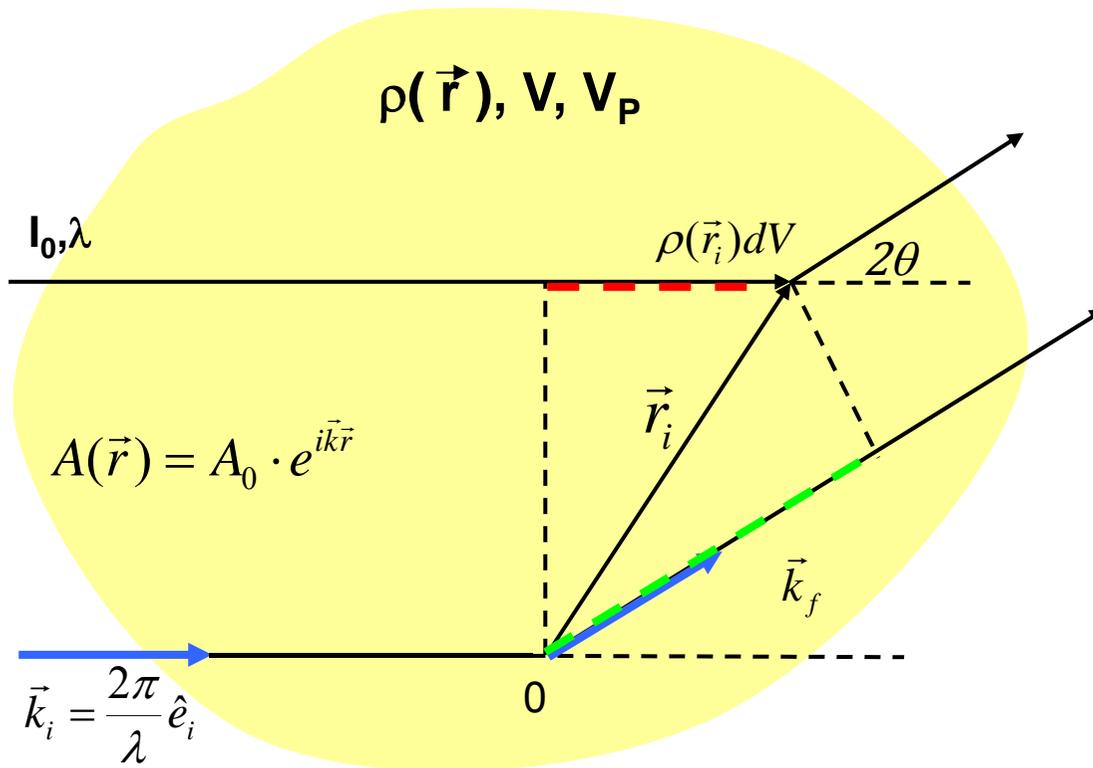


- >  $R$ ~particles "radius"
- >  $d$ ~interatomic distance
- > SAXS:  $\theta < 5^\circ$



# Scattering amplitude

- > Interference in far field



$$|\vec{q}| = \frac{4\pi}{\lambda} \sin \theta$$

- > Phase difference:  $\Delta\varphi_i = (\vec{k}_f - \vec{k}_i) \cdot \vec{r}_i = \vec{q} \cdot \vec{r}_i$
- > Scattering amplitude:  $A(\vec{q}) = \int \rho(\vec{r}) e^{-i\vec{q}\vec{r}} dV = \int \rho(\vec{r}) e^{-i\vec{q}\vec{r}} d^3\vec{r}$
- > Intensity:  $I(\vec{q}) = \frac{1}{V} |A(\vec{q})|^2$

# Form factor and structure factor: Fourier transform

Single particle: Fourier transformation

$$A(\vec{q}) = \int \rho_P(\vec{r}) e^{-i\vec{q}\vec{r}} dV = \int \rho_P(\vec{r}) e^{-i\vec{q}\vec{r}} d^3\vec{r}$$

Particle distribution function  $G(\mathbf{r}) \rightarrow$  Electron density distribution

$$\rho(\vec{r}) = \sum_i \rho_P(\vec{r}_i) = \int \rho_P(\vec{r}') G(\vec{r} - \vec{r}') d^3r' = \rho_P(\vec{r}) * G(\vec{r})$$

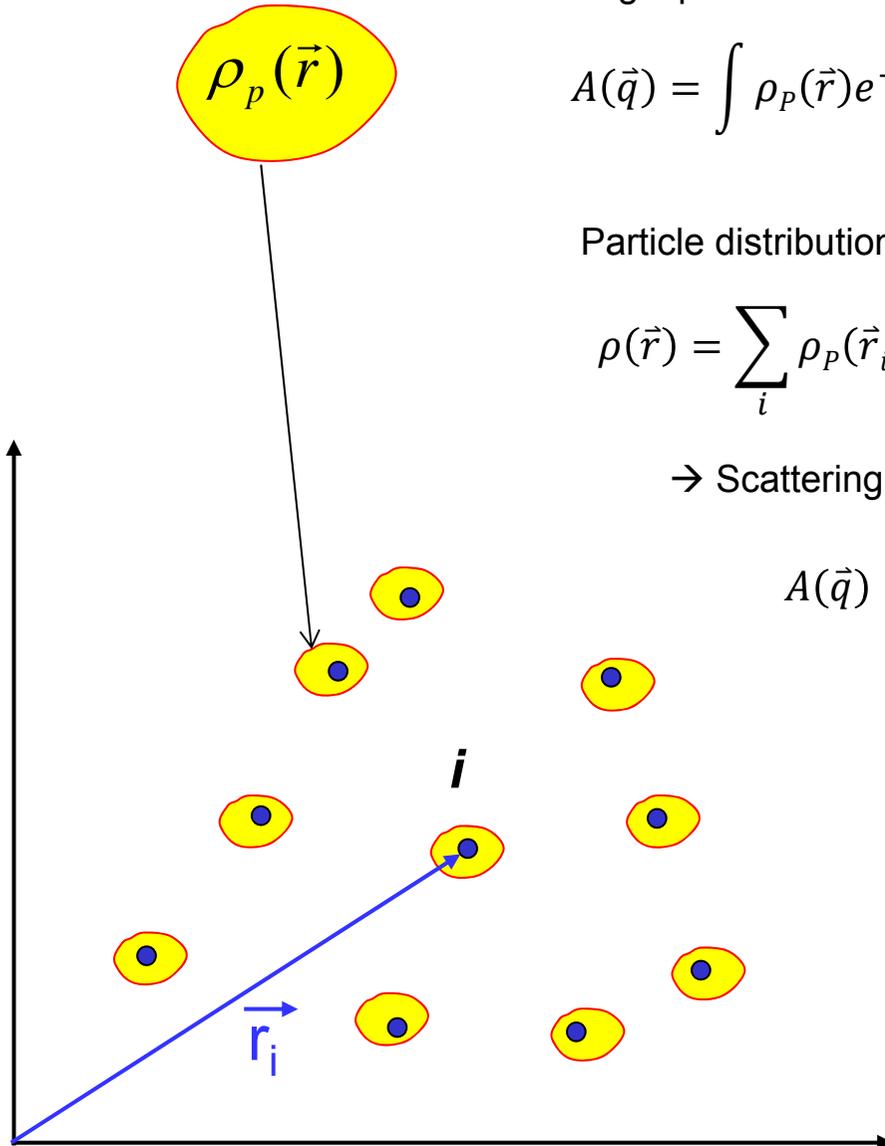
$\rightarrow$  Scattering amplitudes of the whole arrangement

$$\begin{aligned} A(\vec{q}) &= \int \rho(\vec{r}) e^{-i\vec{q}\vec{r}} dV = \int [\rho_P(\vec{r}) * G(\vec{r})] e^{-i\vec{q}\vec{r}} d^3\vec{r} \\ &= \int \rho_P(\vec{r}) e^{-i\vec{q}\vec{r}} dV \cdot \int G(\vec{r}) e^{-i\vec{q}\vec{r}} dV \end{aligned}$$

$\rightarrow$  Scattered Intensity

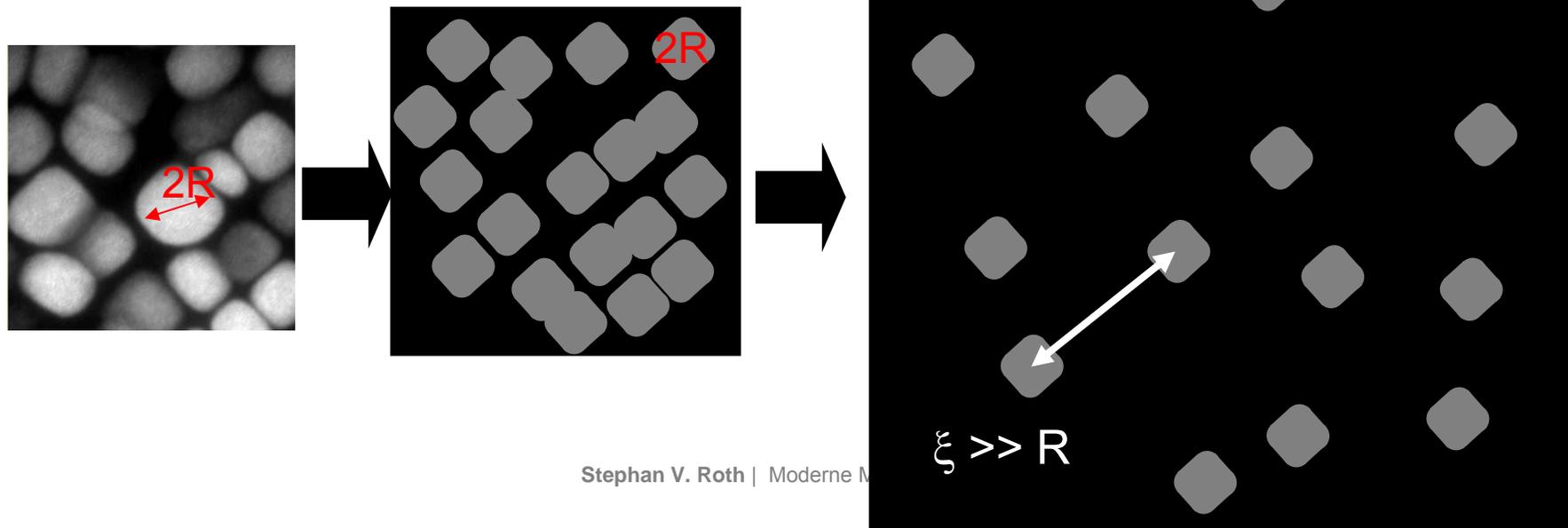
$$I(\vec{q}) = \frac{1}{V} |A(\vec{q})|^2 = P(\vec{q}) S(\vec{q})$$

Form factor    Structure factor



# Two-phase model: Dilute systems

- > Only form of particle relevant
- > Matrix  $M$ , volume fraction  $\Phi$   
Particles  $P$ , volume fraction  $(1-\Phi)$   
Electron density:  $\rho_{M,P} = n_{M,P} * f_{M,P}$
- $f_{M,P}$ : atomic form factor („extension of the electron cloud“, resonances)
- $n_{M,P}$ : number density of atoms
- > Consider  $\rho_{M,P}$  as constant resp.



# Two phase Model

- > Scattering amplitude:

$$A(\vec{q}) = \int \rho(\vec{r}) e^{-i\vec{q}\vec{r}} d^3\vec{r} = \int_{\Phi V} \rho_M(\vec{r}) e^{-i\vec{q}\vec{r}} d^3\vec{r} + \int_{(1-\Phi)V} \rho_P(\vec{r}) e^{-i\vec{q}\vec{r}} d^3\vec{r}$$

$$A(\vec{q}) = (\rho_M - \rho_P) \int_{\Phi V} e^{-i\vec{q}\vec{r}} d^3\vec{r}$$

$$A(\vec{q}) = \Delta\rho \int_{\Phi V} e^{-i\vec{q}\vec{r}} d^3\vec{r}$$

- >  $I(\vec{q}) = \frac{1}{V} |A(\vec{q})|^2 \sim \Delta\rho^2$

- > Porod Invariant Q (Porod, 1982):

$$Q = \int I(\vec{q}) d^3\vec{q} = 4\pi\Phi(1-\Phi)\Delta\rho^2$$

Ableiten!

Mittlung <..> erklären S.25, S.51

- > Only dependent on density contrast  $\Delta\rho$



# Herleitung Porod-Invariante

- > Siehe Handzettel und Übung
- > Q-Berechnung Übung



# Two phase Model – single particle approximation

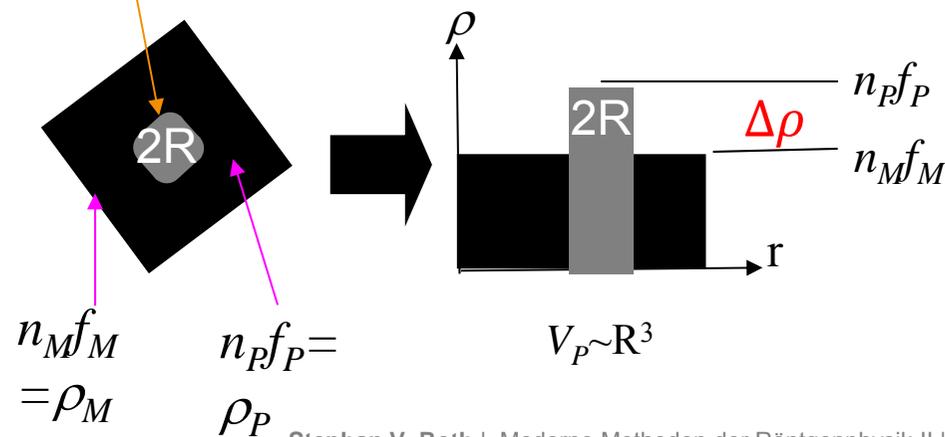
> Amplitude:  $A(\vec{q}) = \Delta\rho \int_{\Phi_V} e^{-i\vec{q}\vec{r}} d^3\vec{r}$

> Intensity:  $I(\vec{q}) = \frac{1}{V} |A(\vec{q})|^2$

> Closer look at  $I(q)$  for dilute systems:  $N_P$  independent scatterers

> Incoherent sum of intensities:

$$I(\vec{q}) \sim NP V_P^2 \Delta\rho^2 \left| \frac{1}{V_P} \int_{V_P} e^{-i\vec{q}\vec{r}} d^3r \right|^2$$



# Two phase Model – single particle approximation

> Amplitude:  $A(\vec{q}) = \Delta\rho \int_{\Phi_V} e^{-i\vec{q}\vec{r}} d^3\vec{r}$

> Intensity:  $I(\vec{q}) = \frac{1}{V} |A(\vec{q})|^2$

> Closer look at  $I(q)$  for dilute systems:  $N_P$  independent scatterers

> Incoherent sum of intensities:

$$I_m(\vec{q}) \sim NP V_P^2 \Delta\rho^2 \left| \frac{1}{V_P} \int_{V_P} e^{-i\vec{q}\vec{r}} d^3r \right|^2$$

$$P(q) = \left| 3 \frac{\sin(qR) - qR \cos(qR)}{(qR)^3} \right|^2$$

- Form factor of a **sphere of radius  $R$**
- Isotropic scattering



# Colloid: homogeneous sphere of radius R

A simple, but important calculation:

$$F(\vec{q}) = \int_{V=\text{particleVolume}} \rho(\vec{r}) \cdot e^{-i\vec{q}\vec{r}} \cdot d^3r = \int_0^R \int_0^{2\pi} \int_0^\pi \rho_0 \cdot e^{-i\vec{q}\vec{r}} r^2 \sin(\theta) d\theta d\varphi dr$$

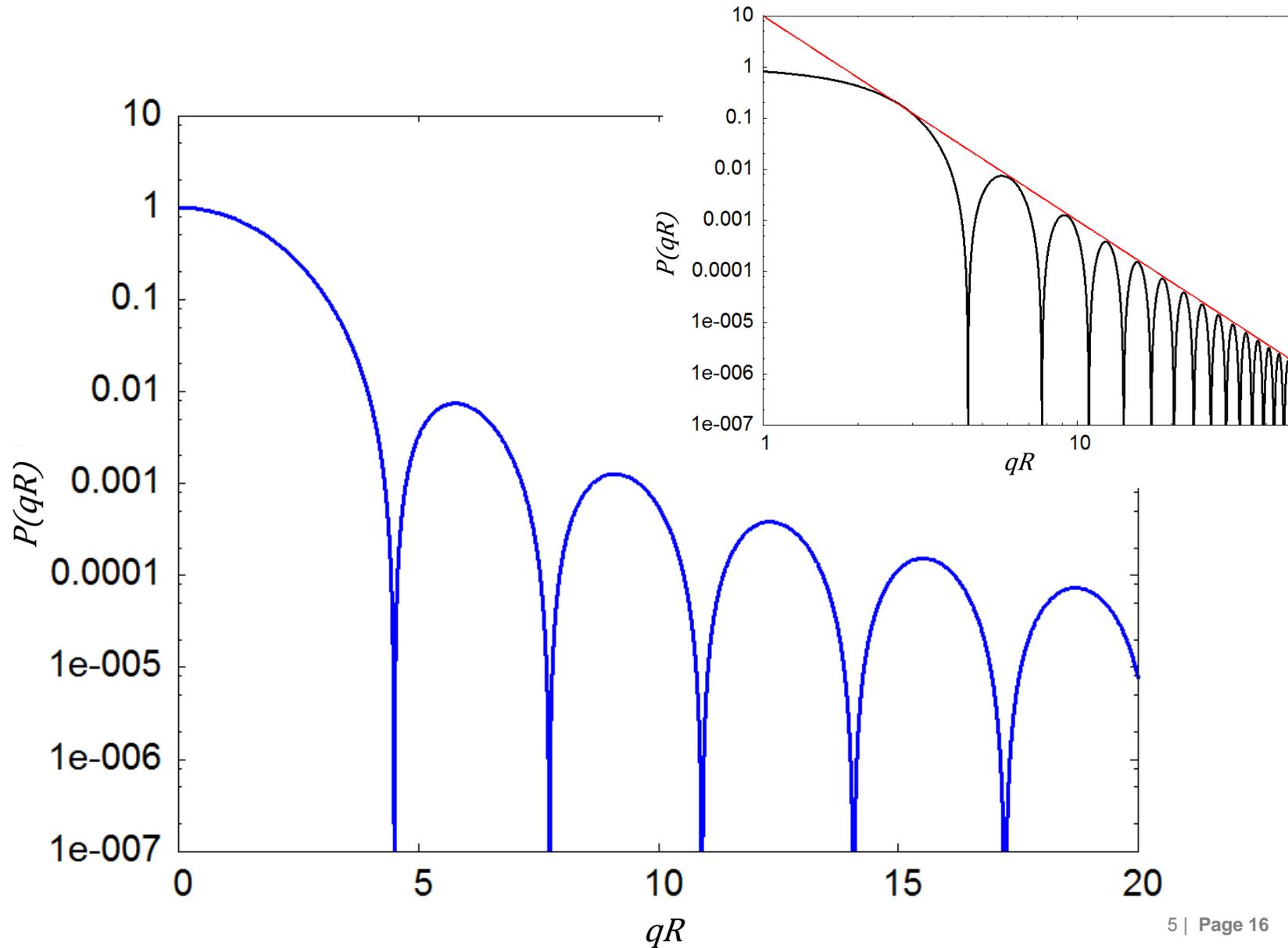
$$F(\vec{q}) = \rho_0 2\pi \int_0^R \int_0^\pi e^{-iqr \cos(\theta)} r^2 \sin(\theta) d\theta d\varphi dr = \rho_0 2\pi \int_0^R \frac{e^{iqr} - e^{-iqr}}{qr} r^2 \sin(\theta) dr$$

$$F(\vec{q}) = \rho_0 2\pi \cdot \frac{2}{q} \int_0^R \sin(qr) r dr = \frac{4\pi\rho_0}{q} \left[ -\frac{r \cos(qr)}{q} \Big|_0^R + \int_0^R \frac{\cos(qr)}{q} dr \right]$$

$$F(\vec{q}) = \frac{4\pi\rho_0}{q} \left[ -\frac{R \cos(qR)}{q} + \frac{\sin(qR)}{q^2} \right] = 4\pi R^3 \rho_0 \frac{(\sin(qR) - qR \cos(qR))}{(qR)^3}$$



# Colloid: homogeneous sphere of radius R



# Guinier radius

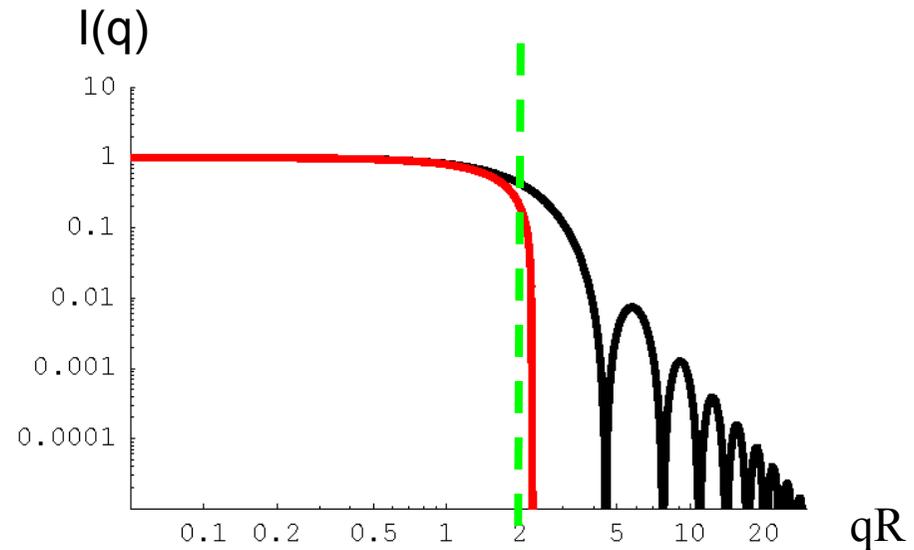
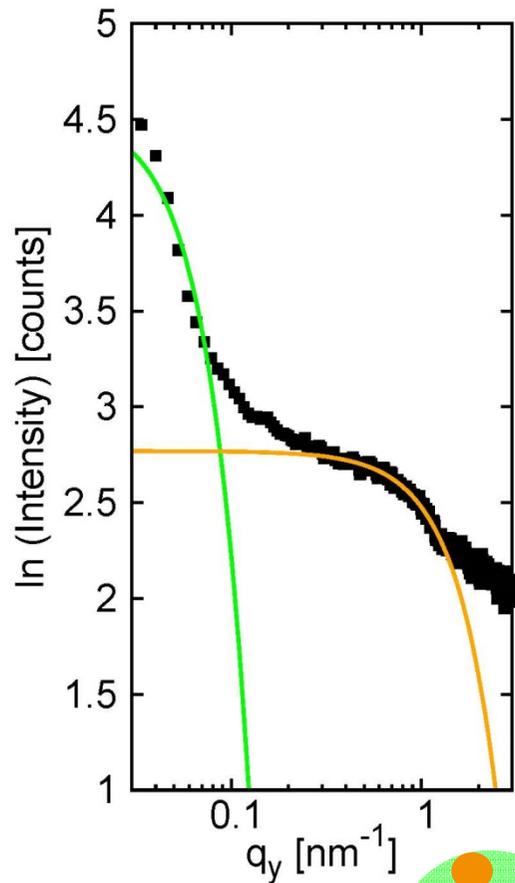
- >  $Q \rightarrow 0$
- > Homogenous sphere of radius  $R$

$$P(q) = \left| 3 \frac{\sin(qR) - qR \cos(qR)}{(qR)^3} \right|^2 \sim 1 - \frac{1}{5} q^2 R^2 \sim \exp\left(-\frac{1}{5} q^2 R^2\right) \quad \text{Ableiten}$$

- > Radius of gyration:  
replace homogenous sphere by shell of same moment of inertia:  $R_g$
- >  $R_g = \sqrt{3/5} R$
- >  $P(q) \sim \exp\left(-\frac{1}{3} q^2 R_g^2\right)$  general form of Guinier law [Guinier (1955)]
- > Independent of particle form



# Guinier Approximation



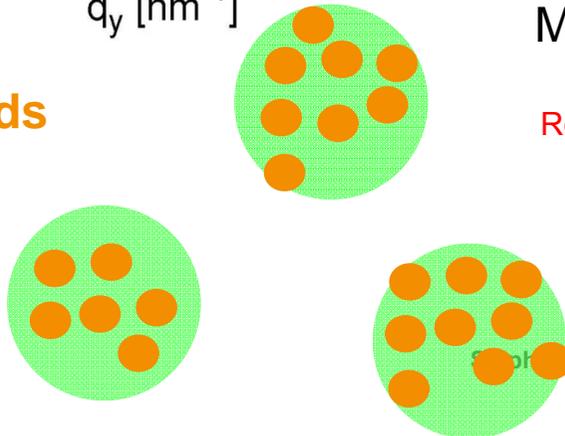
$$\lim_{q \rightarrow 0} I(q) = \Delta\rho^2 \cdot V^2 \cdot \exp\left(-q^2 \cdot \frac{R_g^2}{3}\right)$$

Radius of Gyration  $R_g$

Monodisperse spheres of radius  $R$ :  $R_g = \sqrt{3/5} \cdot R$

Roth et al., Appl. Phys. Lett. **91**, 091915 (2007)

2nm Colloids  
domains

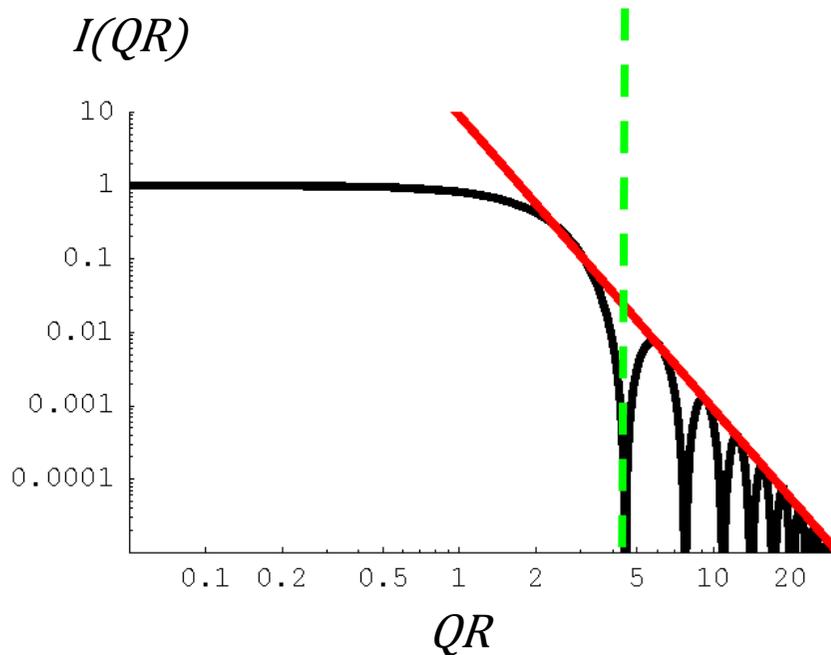


# Porod's law: large $q$

Scattered intensity:  $\sim \left| 4\pi R^3 \rho_0 \frac{(\sin(qR) - qR \cos(qR))}{(qR)^3} \right|^2$

Look at maxima of form factor

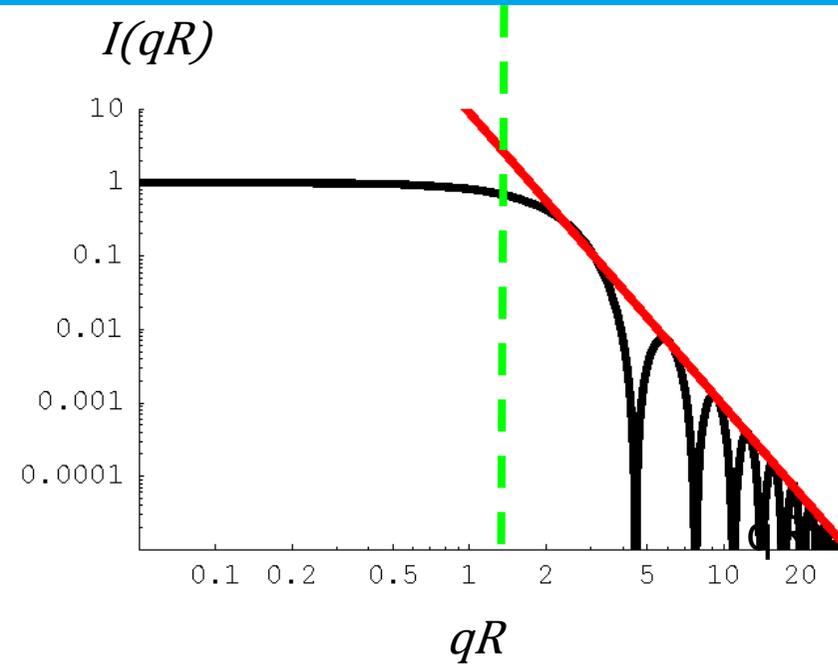
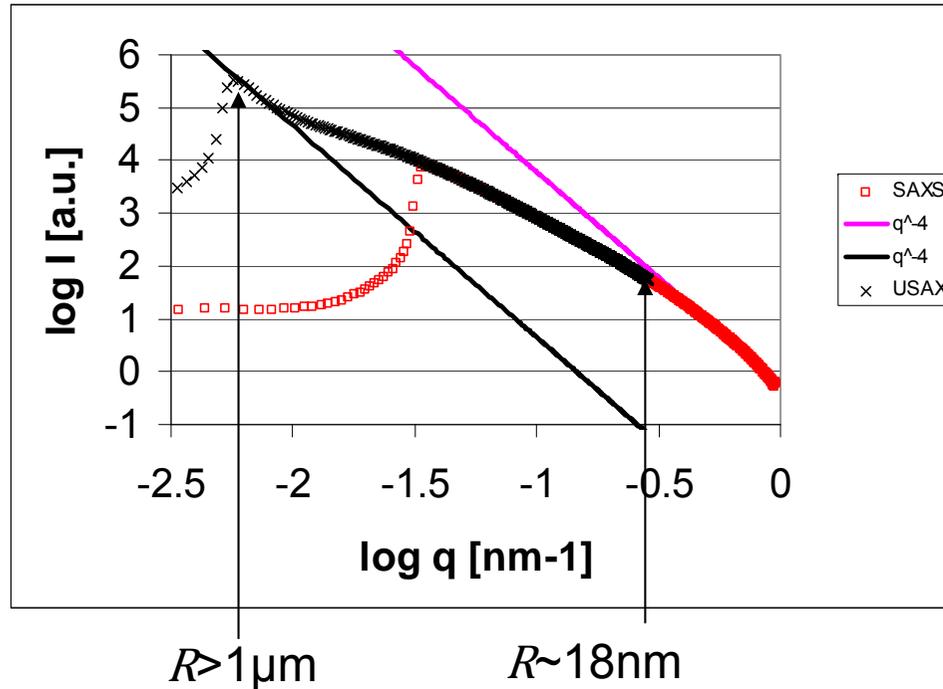
$$\begin{aligned} &\sim \left| 4\pi \rho_0 \frac{(\sin(qR) - qR \cos(qR))}{(qR)^3} \right|^2 \\ &\leq \left( 4\pi \rho_0 \frac{|\sin(qR)| + qR |\cos(qR)|}{(qR)^3} \right)^2 \\ &\sim \left( 4\pi \rho_0 \frac{1 + qR}{(qR)^3} \right)^2 \sim \left( 4\pi \rho_0 \frac{qR}{(qR)^3} \right)^2 \\ &\sim \frac{1}{(q)^4} \frac{R^2}{R^6} \sim \frac{S}{V_P^2} q^{-4} \end{aligned}$$



Surface of sphere



# Porod's Law



$$P(qR > 4.5) = 2\pi \left( \frac{S}{V_P^2} \right) q^{-4}$$

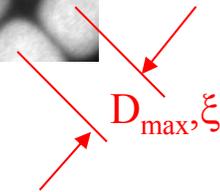
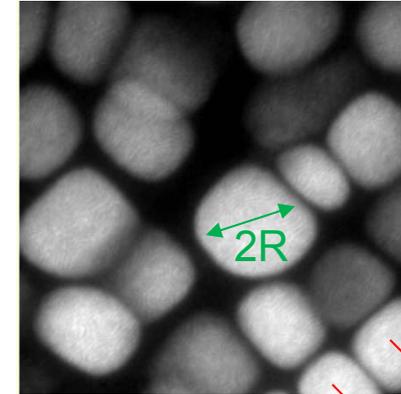
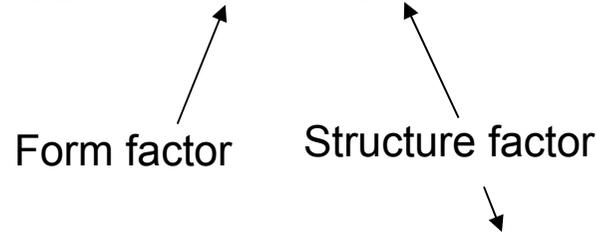
- > Depends only on Surface and particle Volume
- > No shape dependance



# The structure factor – many particles, close distance

- > Real systems: not dilute, many particles...
- > Generalisation of Bragg's Law in crystallography:

$$I(q) = c P(q) S(q)$$



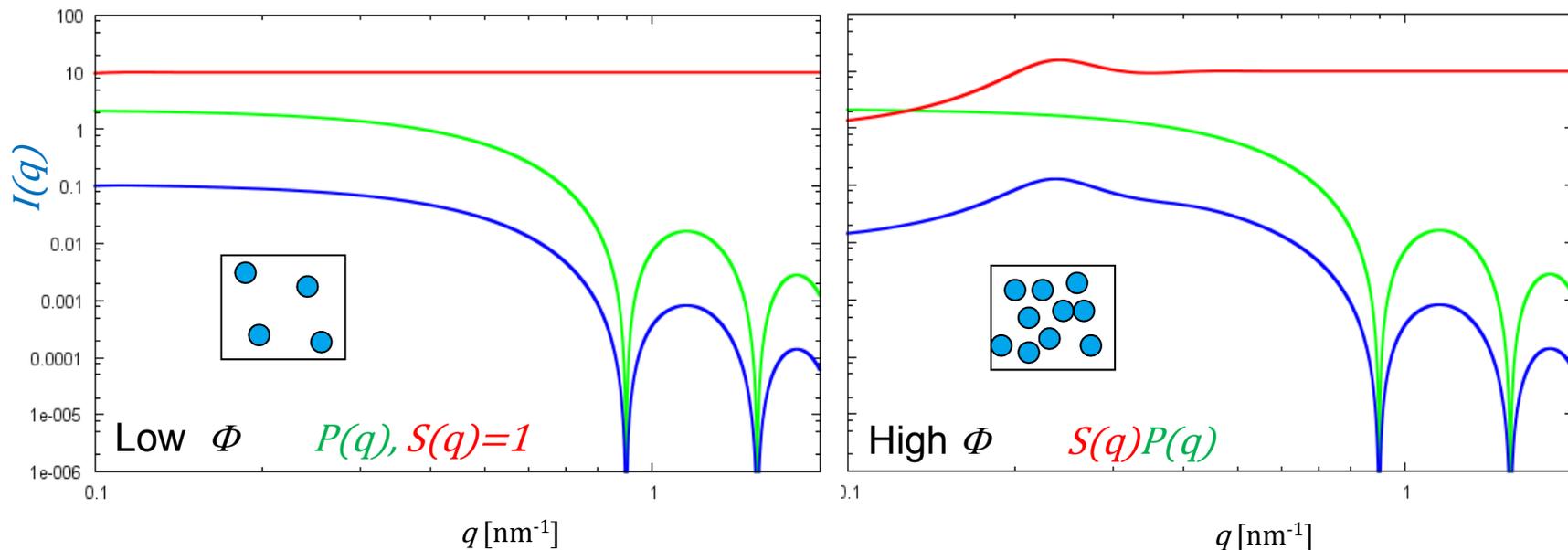
- > Periodic ordering with periodicity  $d, \xi$  in the electron density :
- >  $I(q)$  shows a corresponding maximum at  $q = 2\pi / (D_{max}, \xi)$

$$S(q) \propto \frac{1 - \exp(-2\sigma_D^2 q^2)}{1 - 2 \exp(-\sigma_D^2 q^2) \cos(qD_{max}) + \exp(-2\sigma_D^2 q^2)}$$

Smearing      Distance of particles

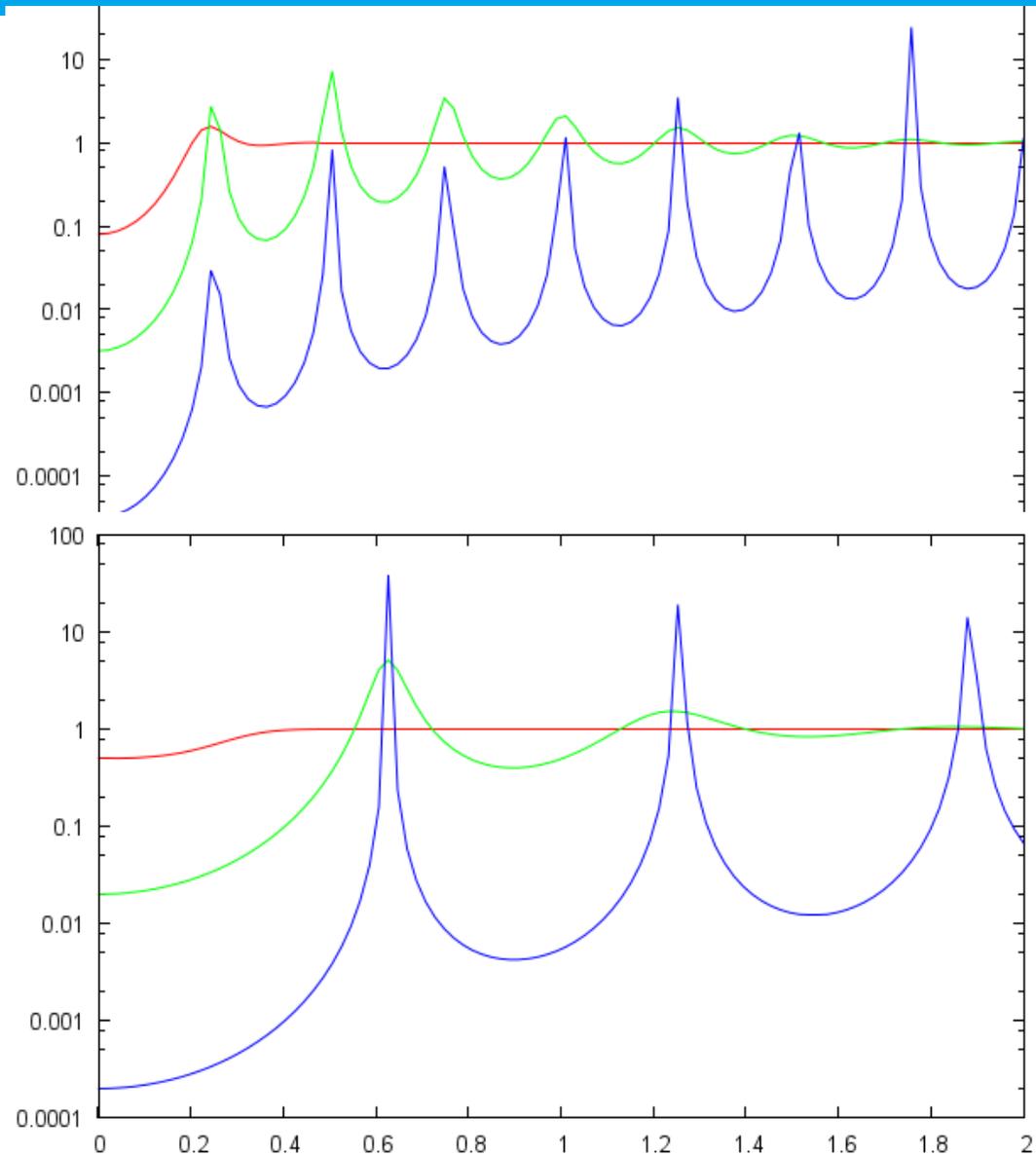
# The structure factor – many particles, close distance

- > Real systems: not dilute, many particles...
- > Generalisation of Bragg's Law in crystallography:  
 $I(q) = c P(q) S(q)$
- > Examples:  $R=5\text{nm}$ ,  $D_{\max}=100\text{nm}$ ,  $25\text{nm}$ ,  $\sigma D/D_{\max}=25\%$



# Structure factor and form factor

- >  $D_{max} = 25\text{nm}$   
 $D_{max} = 10\text{nm}$
- >  $\sigma_D = 5\text{nm}, 1\text{nm}, 0.1\text{nm}$
- >  $S(q) \rightarrow 1 \quad q \rightarrow \infty$   
well separated particles

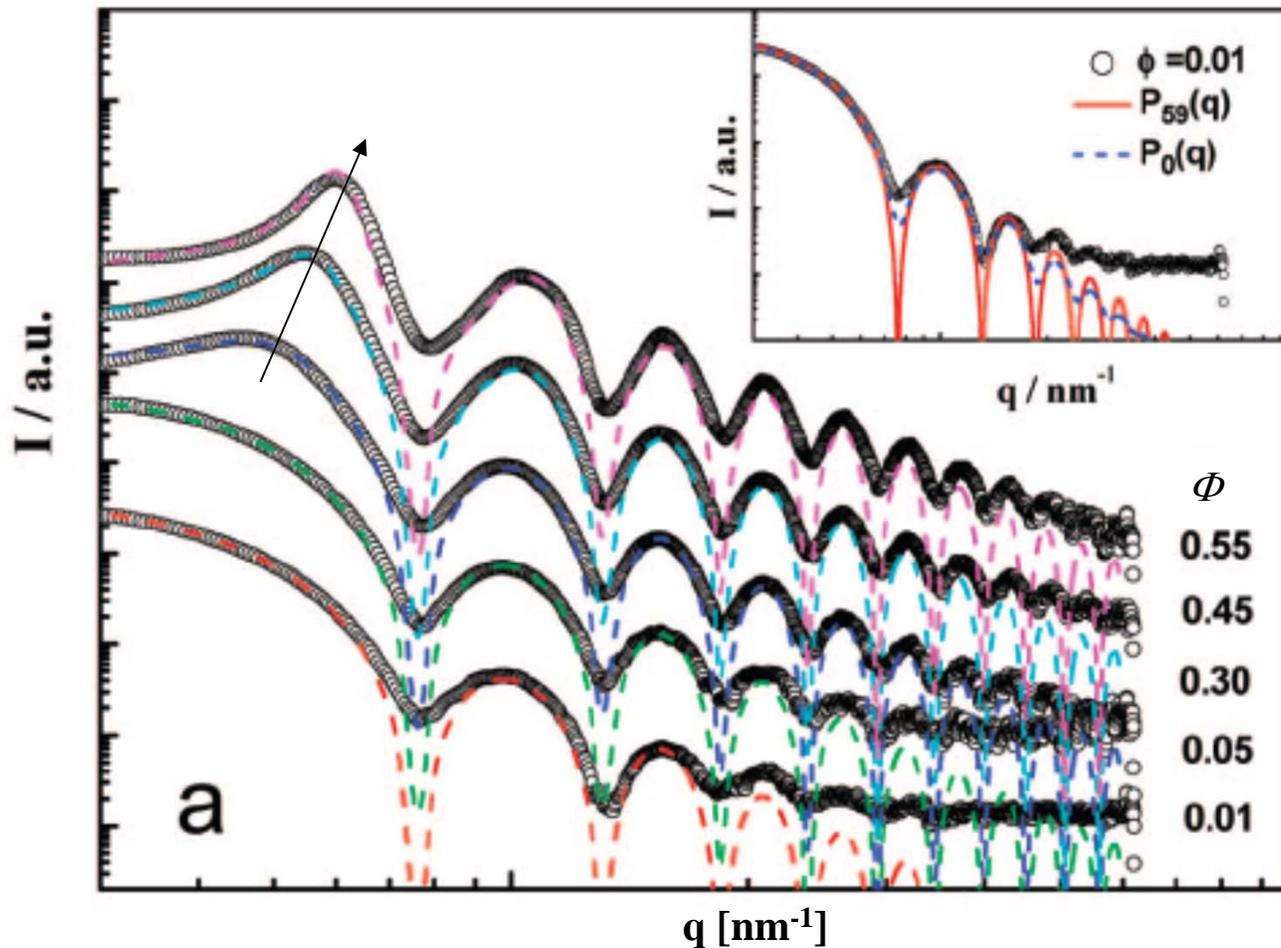
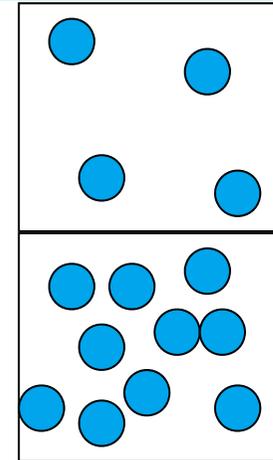


# Colloidal systems

> Latex spheres in water

$$I(q) = c P(q) S(q)$$

Low  $\Phi$   $P(q), S(q) = 1$   
 High  $\Phi$   $S(q)P(q)$



- > Gaussian distribution of particle sizes
- > Shift in maximum: Decreasing distance