

# Methoden moderner Röntgenphysik II: Streuung und Abbildung

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Lecture 8	Vorlesung zum Haupt- oder Masterstudiengang Physik, SoSe 2015 G. Grübel, M. Martins, E. Weckert	
Location	Lecture hall AP, Physics, Jungiusstraße	
Date	Tuesday	12:45 - 14:15
	Thursday	8:30 - 10:00



# Methoden moderner Röntgenphysik II: Streuung und Abbildung

## Part I:

### Basics of X-ray Physics

by Gerhard Grübel (GG)

#### Introduction

Overview, Introduction to X-ray Scattering

#### X-ray Scattering Primer

Elements of X-ray Scattering

#### Sources of X-rays, Synchrotron Radiation

Laboratory Sources, Accelerator Bases Sources

#### Reflection and Refraction from Interfaces

Snell's Law, Fresnel Equations

#### Kinematical Diffraction (I)

Diffraction from an Atom, a Molecule, from Liquids, Glasses, ...

#### Kinematical Diffraction (II)

Diffraction from a Crystal, Reciprocal Lattice, Structure Factor, ...



# Methoden moderner Röntgenphysik II: Streuung und Abbildung

## Small Angle Scattering, and Soft Matter

Introduction, Form Factor, Structure Factor, Applications, ...

## **Anomalous Diffraction**

**Introduction into Anomalous Scattering, ...**

## Introduction into Coherence

Concept, First Order Coherence, ...

## Coherent Scattering

Spatial Coherence, Second Order Coherence, ...

## Applications of Coherent Scattering

Imaging and Correlation Spectroscopy, ...



# Resonant Scattering (phasing, magnetism,...)

Scattering length of an atom:  $-r_0 f^0(\mathbf{Q})$

$f^0(\mathbf{Q})$  atomic form factor (Fourier transform of charge distribution)

$r_0$  Thomson scattering length of single electron

In order to include absorption effects ( $f''$ ) atoms, a more elaborate model than the free electron gas is needed.

→ Electrons are bound to atoms

→ Forced oscillator model with resonant frequency  $\omega_s$  and damping constant  $\Gamma$

Include dispersion corrections ( $f'$ ,  $f''$ ):

[ Note:  $f'' = \left(\frac{k}{4\pi r_0}\right) \sigma_a$  ]

$$f(\mathbf{Q}, \omega) = f^0(\mathbf{Q}) + f'(\omega) + i f''(\omega) \quad \text{[in units of } r_0\text{]}$$



# Resonant Scattering

Classical model of an electron bound in an atom in E field

$$\mathbf{E}(\mathbf{r}, t) = \hat{x} E_0 e^{-i\omega t}$$



equation of motion of the electron

$$\ddot{x} + \Gamma \dot{x} + \omega_s^2 x = - \left( \frac{e E_0}{m} \right) e^{-i\omega t}$$

$\Gamma$  = damping  
 $\omega_s$  resonant frequency

Solution:  $x(t) = x_0 e^{-i\omega t}$  →  $x_0 = - \left( \frac{e E_0}{m} \right) \frac{1}{(\omega_s^2 - \omega^2 - i\omega\Gamma)}$  (A)

Radiated field strength at distance R and time t

$$E_{\text{rad}}(R, t) = \left( \frac{e}{4 \epsilon_0 R c^2} \right) \ddot{x} \left( t - \frac{R}{c} \right)$$

(B)

↑  
acceleration at “earlier” time  $(t - \frac{R}{c})$



# Resonant Scattering

Inserting  $\ddot{x} \left( t - \frac{R}{c} \right) = \omega^2 x_0 e^{-i\omega t} e^{i\left(\frac{\omega}{c}\right)R}$  using (A) into (B):

$$E_{\text{rad}}(R, t) = \left( \frac{\omega^2}{(\omega_s^2 - \omega^2 + i\omega\Gamma)} \right) \left( \frac{e^2}{4 \epsilon_0 R c^2} \right) E_0 e^{-i\omega t} \left( \frac{e^{ikR}}{R} \right)$$

or

$$\frac{E_{\text{rad}}(R, t)}{E_{\text{in}}} = -r_0 \underbrace{\frac{\omega^2}{(\omega_s^2 - \omega^2 + i\omega\Gamma)}}_{\text{atomic scattering length } f_s} \left( \frac{e^{ikR}}{R} \right)$$

atomic scattering length  $f_s$  (in units of  $-r_0$ ) for bound electron (C)  
 note:  $f_s \rightarrow 1$  ( $\omega \gg \omega_s$ )

Total cross-section:  $\sigma_T = \left( \frac{8\pi}{3} \right) r_0^2$  (free electron)

$$\sigma_T = \left( \frac{8\pi}{3} \right) \frac{\omega^4}{(\omega^2 - \omega_s^2)^2 + (\omega\Gamma)^2} r_0^2$$

for  $\Gamma = 0$  and  $\omega \ll \omega_s$ :  $\sigma_T = \left( \frac{8\pi}{3} \right) r_0^2 \left( \frac{\omega}{\omega_s} \right)^4$  : “Rayleigh Scattering”



# Resonant Scattering

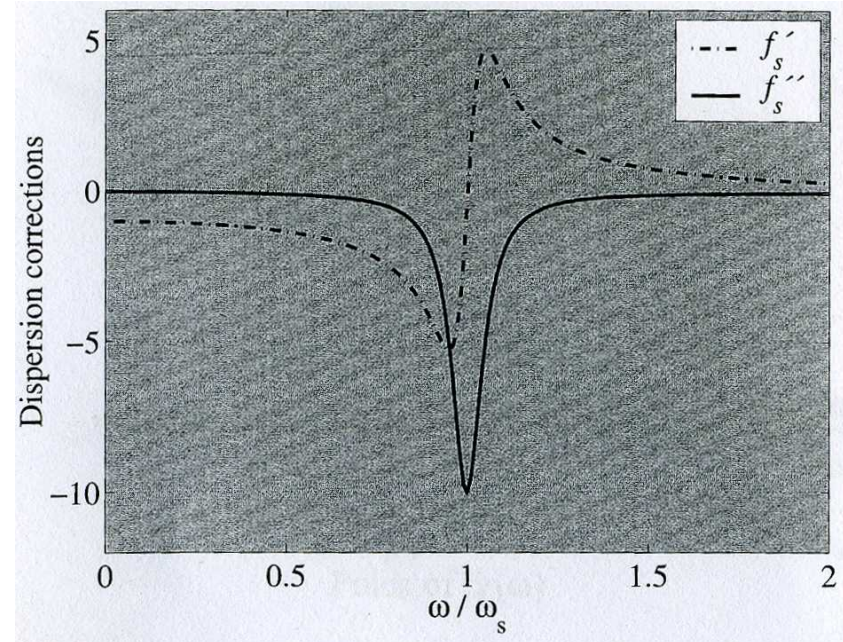
$$f_s = \frac{\omega^2 - \omega_s^2 + i\omega\Gamma + \omega_s^2 - i\omega\Gamma}{(\omega^2 - \omega_s^2 + i\omega\Gamma)}$$

$$= 1 + \frac{\omega_s^2 - i\omega\Gamma}{(\omega^2 - \omega_s^2 + i\omega\Gamma)}$$

$$\approx 1 + \underbrace{\frac{\omega_s^2}{(\omega^2 - \omega_s^2 + i\omega\Gamma)}}_{\text{dispersion correction } \chi(\omega)}$$

dispersion correction  $\chi(\omega)$

$$x(\omega) = f'_s + i f''_s = \frac{\omega_s^2}{(\omega^2 - \omega_s^2 + i\omega\Gamma)}$$



With:

$$f'_s = \frac{\omega_s^2(\omega^2 - \omega_s^2)}{(\omega^2 - \omega_s^2)^2 + (\omega\Gamma)^2}$$

$$f''_s = \frac{\omega_s^2\omega\Gamma}{(\omega^2 - \omega_s^2)^2 + (\omega\Gamma)^2}$$

# Resonant Scattering

Note: since  $f'' = -\left(\frac{k}{4\pi}\right) \sigma_a(E)$  (see J. A-N. & D. McM. p. 70) it follows that the absorption cross-section for a single oscillator model is:

$$\sigma_{a,s}(\omega) = 4 \pi r_0 c \frac{\omega_s^2 \Gamma}{(\omega^2 - \omega_s^2)^2 + (\omega \Gamma)^2}$$

This function has:

- sharp peak at  $\omega = \omega_s$
- $\Delta\omega_{\text{FWHM}} \approx \Gamma$

Thus  $\sigma_a(E)$  may be written with help of a delta function:

$$\sigma_{a,s}(\omega) = 4 \pi r_0 c \frac{\pi}{2} \delta(\omega - \omega_s)$$

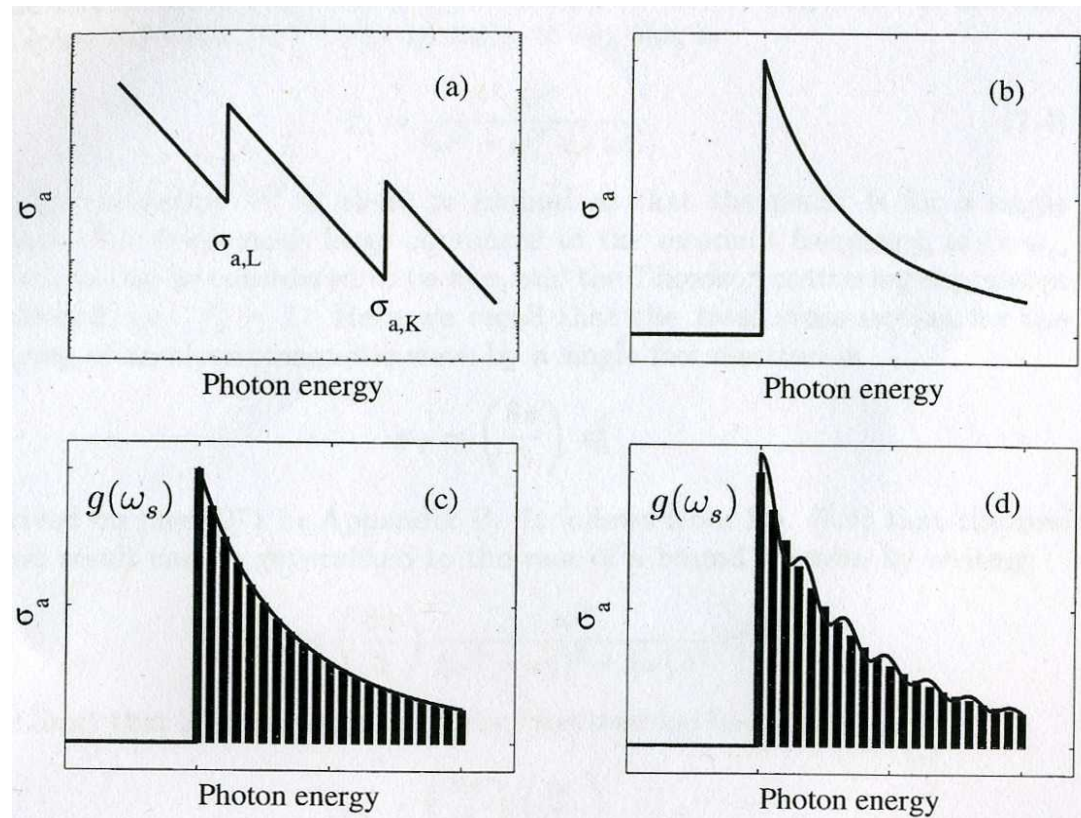
(D)



# Resonant Scattering

The experimentally observed absorption cross-section is NOT a single line spectrum as suggested by (D).

There is a continuum of free states above an absorption edge that the electron can be excited into. This implies a series of different  $\omega_s$ :



# Resonant Scattering

Absorption cross section for multiple harmonic oscillators:

$$\sigma_a(\omega) = 2 \pi^2 r_0 c \sum_s g(\omega_s) \delta(\omega - \omega_s)$$

where  $g(\omega_s)$  is the relative weight of each transition

Geben Sie hier eine Formel ein.

The real part of the dispersion becomes:

$$f''(\omega) = \sum_s g(\omega_s) f'_s(\omega, \omega_s) \quad (F)$$

(F) does not describe e.g. “white lines” or “EXAFS” oscillations (see figure) in the absorption cross-section arising from the particular environment of the resonantly scattering atom.



# Resonant Scattering

Measure absorption cross-section and use (E) to obtain  $f''$ :

$$f''(\omega) = - \left( \frac{\omega}{4 \pi r_0 c} \right) \sigma_a(\omega)$$

Use **Kramers-Kronig relations** to obtain  $f'$ :

$$f''(\omega) = \frac{1}{\pi} P \int_{-\infty}^{+\infty} \frac{f''(\omega')}{(\omega' - \omega)} d\omega' = \frac{2}{\pi} P \int_0^{+\infty} \frac{\omega' f''(\omega')}{(\omega'^2 - \omega^2)} d\omega'$$

$$f'(\omega) = - \frac{1}{\pi} P \int_{-\infty}^{+\infty} \frac{f'(\omega')}{(\omega' - \omega)} d\omega' = - \frac{2\omega}{\pi} P \int_0^{+\infty} \frac{f'(\omega')}{(\omega'^2 - \omega^2)} d\omega'$$

$P$  stands for “principal value” (see also comments J. A-N & D. McM p. 242)

# Resonant Scattering

Friedel's Law and Bijvoet Pairs

The Phase Problem in Crystallography

The MAD Method

(Resonant) Magnetic Scattering

