

- # Methoden moderner Röntgenphysik II: Streuung und Abbildung

Lecture 3

Vorlesung zum Haupt/Masterstudiengang Physik

SS 2013

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Location: Hörs AP, Physik, Jungiusstrasse

Tuesdays 12.45 – 14.15

Thursdays 8:30 – 10.00

Methoden moderner Röntgenphysik II: Streuung und Abbildung

Vorlesung: 4 SWS: Dienstag und Donnerstag
Übungen: 2 SWS: Dienstag (wenn vereinbart)
Proseminar: für Bachelor Studierende
8 Leistungspunkte für dieses Modul im Masterstudiengang

Tuesdays 12.45 – 14.15:
Thursdays 8.30 – 10.00: starting April 4, 2013

Tuesdays 14:30 – 16:00: *Tutorials/Übungen*
Organisation-1st meeting: April 9, 2013 in SemRm 4
First Tutorial: April 16 in SemRm 4

Methoden moderner Röntgenphysik II: Streuung und Abbildung

Gerhard Grübel (GG), Thomas Schneider (TS), Oliver Seeck (OS),
Stephan Roth (SR), Hermann Franz (HF)

4.4. – 16.5. Basics of X-ray physics (GG)

14.5. Site Visit

28.5.-13.6. Surfaces and Interfaces (OS)

18.6.-27.6. Biology (TS)

2.7.-11.7. Soft Matter (SR)

Site Visit: May 14, 2 pm at the DESY/Bahrenfeld site
Building 25f (Cafeteria)

■ Methoden moderner Röntgenphysik II: Streuung und Abbildung

Introduction

Overview, Introduction to X-ray scattering

X-ray Scattering Primer

Elements of X-ray scattering

Sources of X-rays, Synchrotron Radiation

Laboratory sources, accelerator based sources

Reflection and Refraction

Snell's law, Fresnel equations,

Kinematical Diffraction (I)

Diffraction from an atom, molecule, liquids, glasses, ..

Kinematical Diffraction (II)

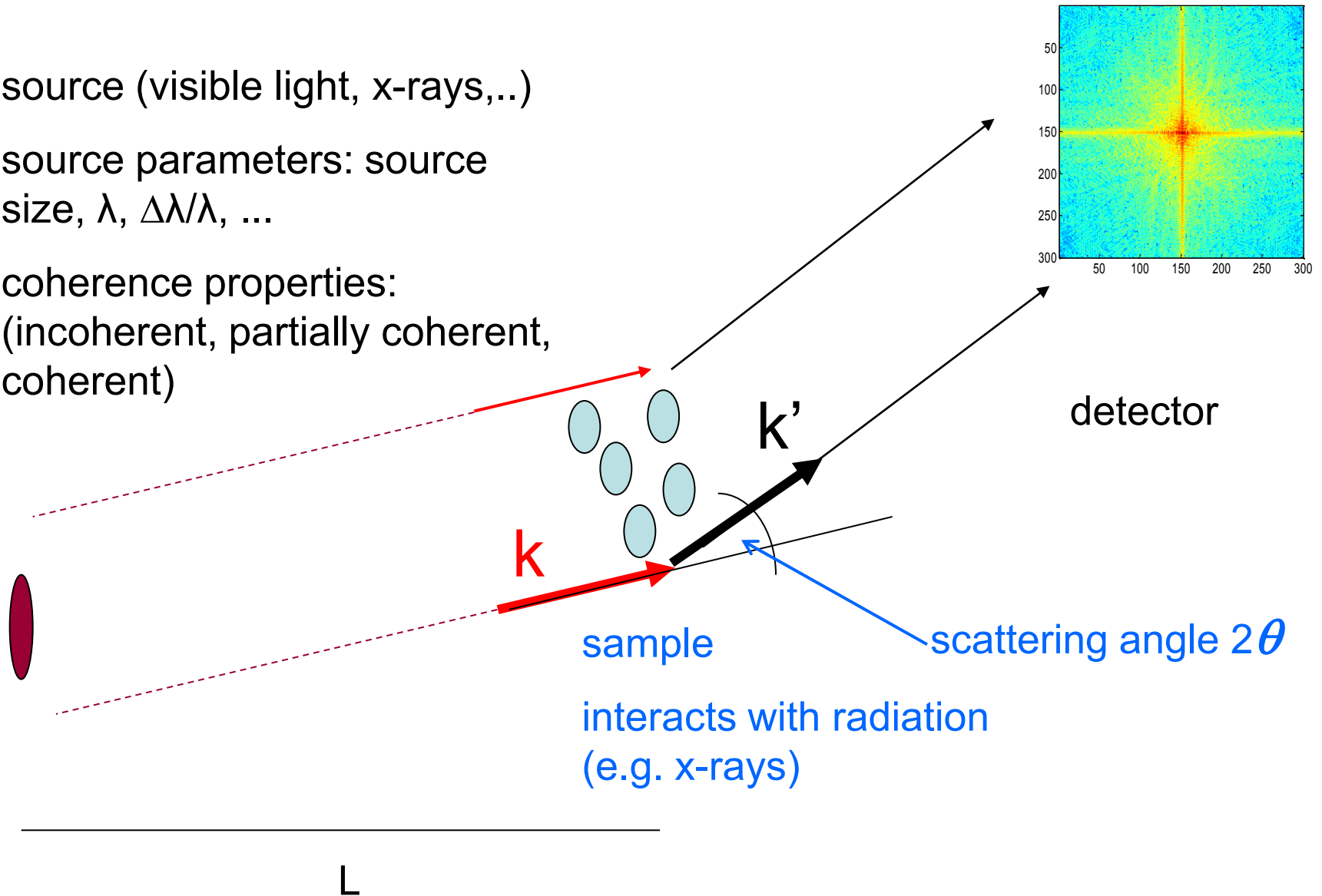
Diffraction from a crystal, reciprocal lattice, structure factor, ..

Experimental Set-Up for Scattering Experiments

source (visible light, x-rays,...)

source parameters: source size, λ , $\Delta\lambda/\lambda$, ...

coherence properties:
(incoherent, partially coherent, coherent)

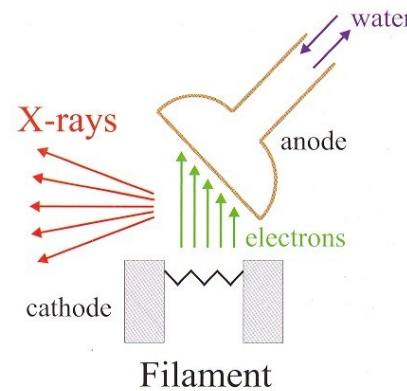


▪ Sources of X-Rays

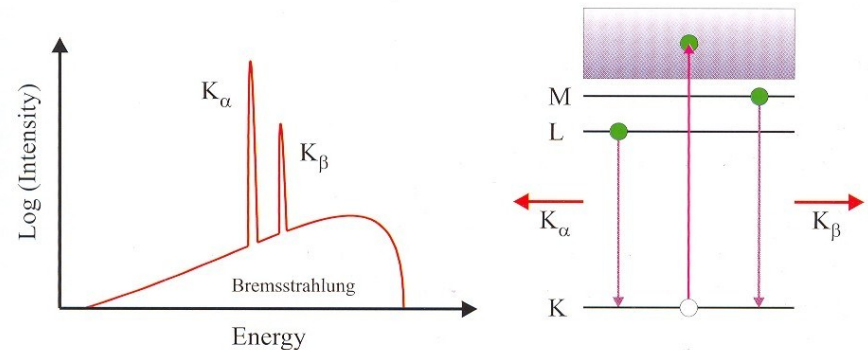
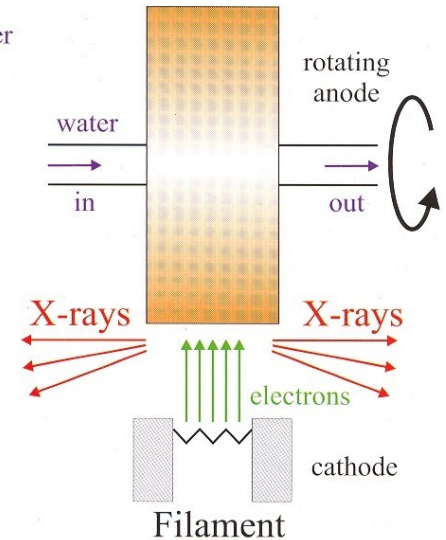
- 1895 discovered by W.C. Röntgen
- 1912 First diffraction experiment (v. Laue)
- 1912 Coolidge tube (W.D. Coolidge, GE)
- 1946 Radiation from electrons in a synchrotron, GE, Physical Review, 71,829(1947)



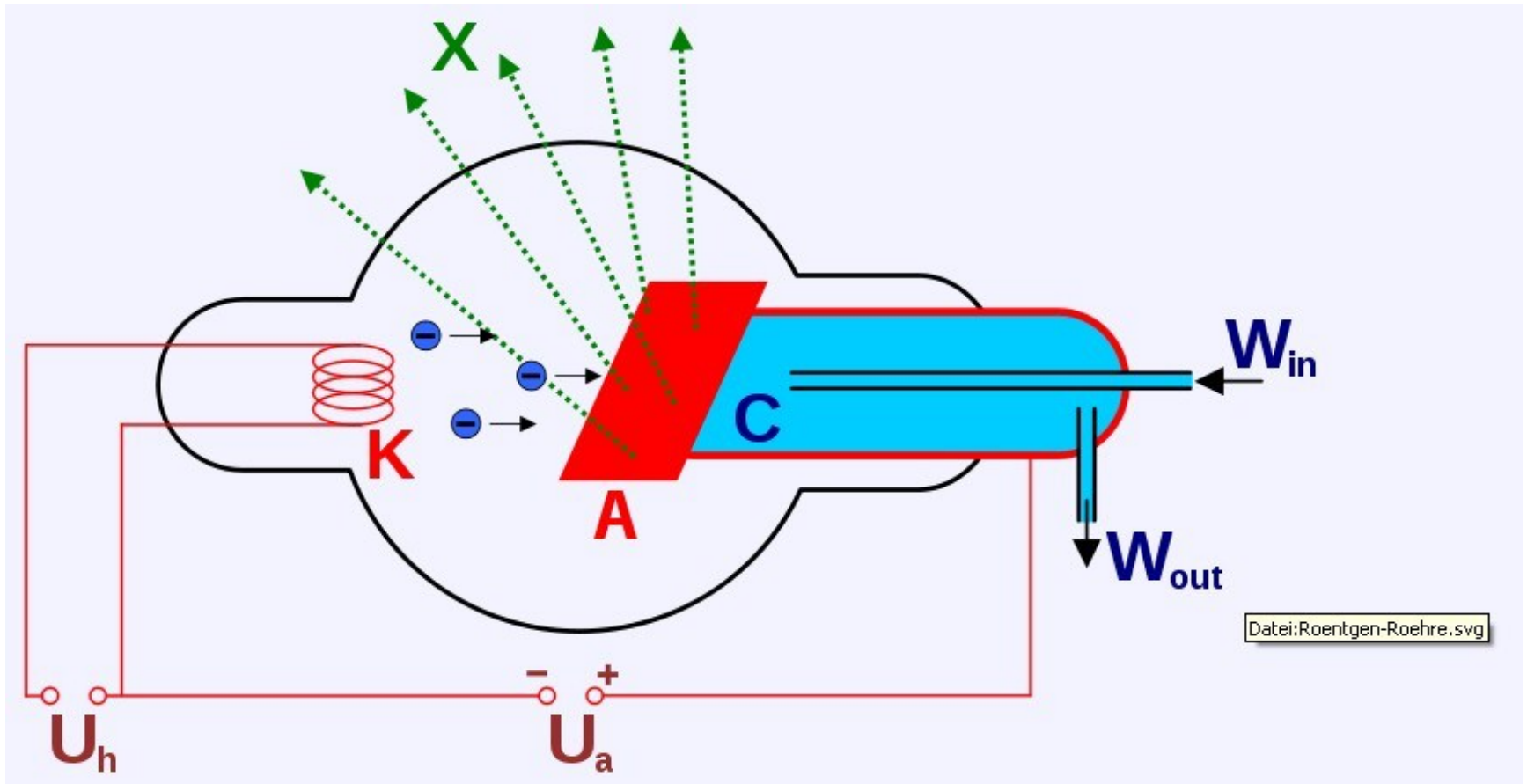
Coolidge Tube



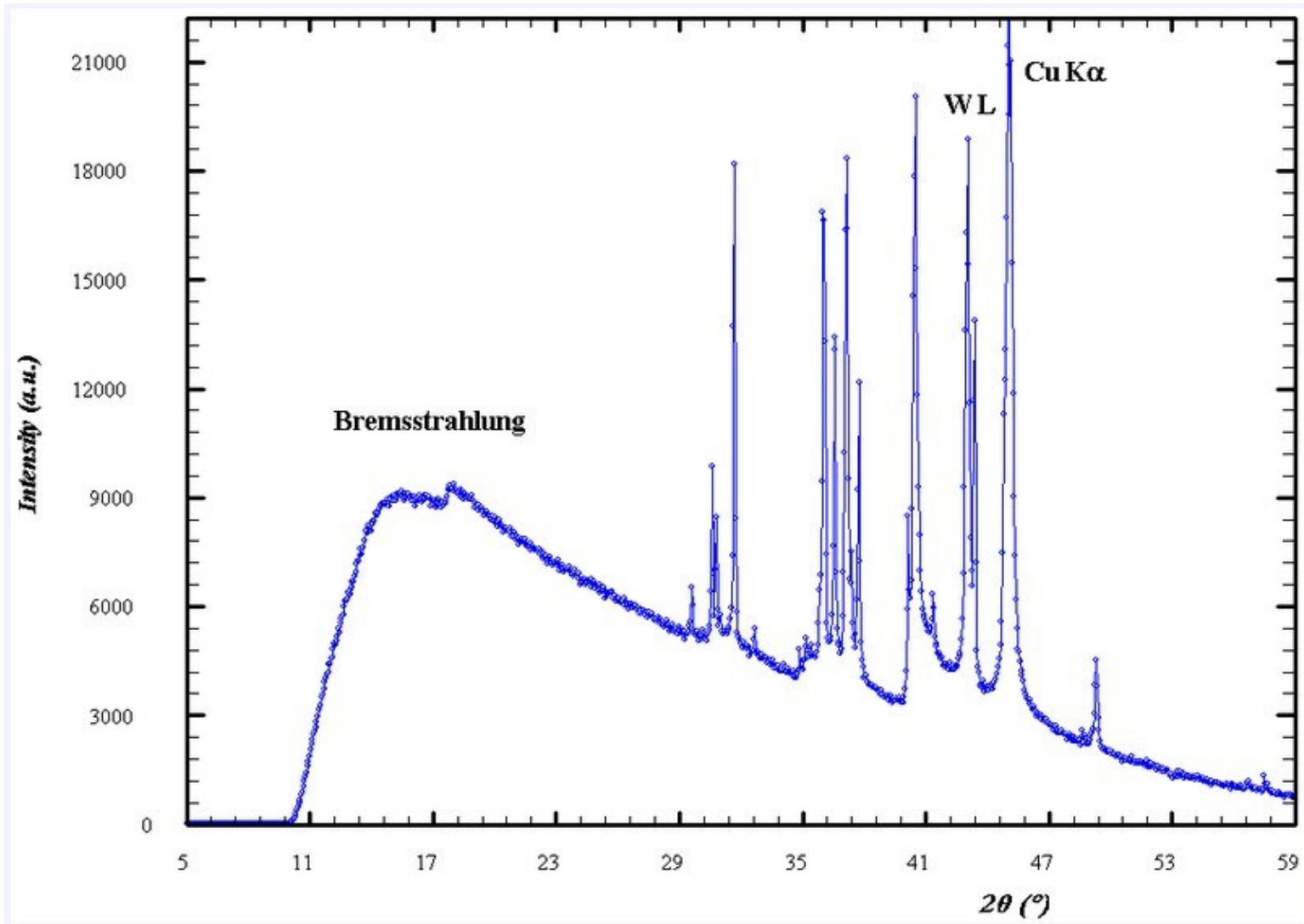
Rotating Anode



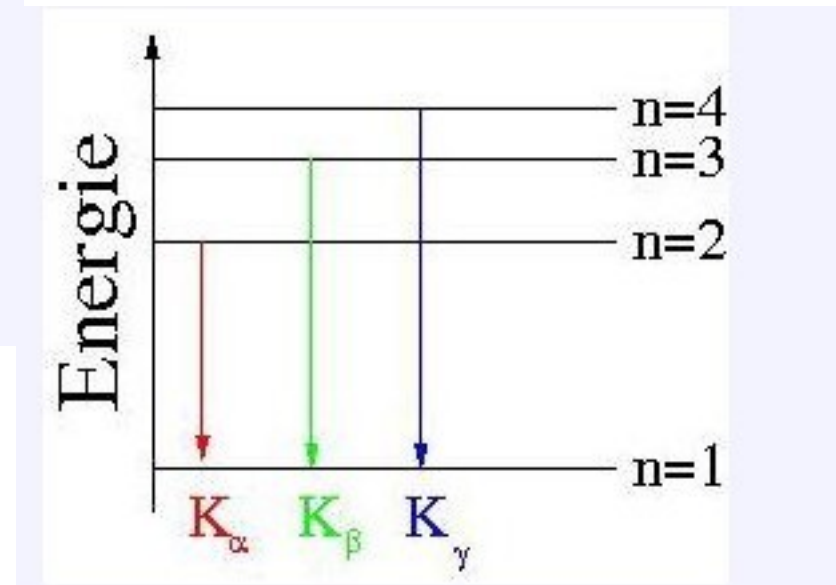
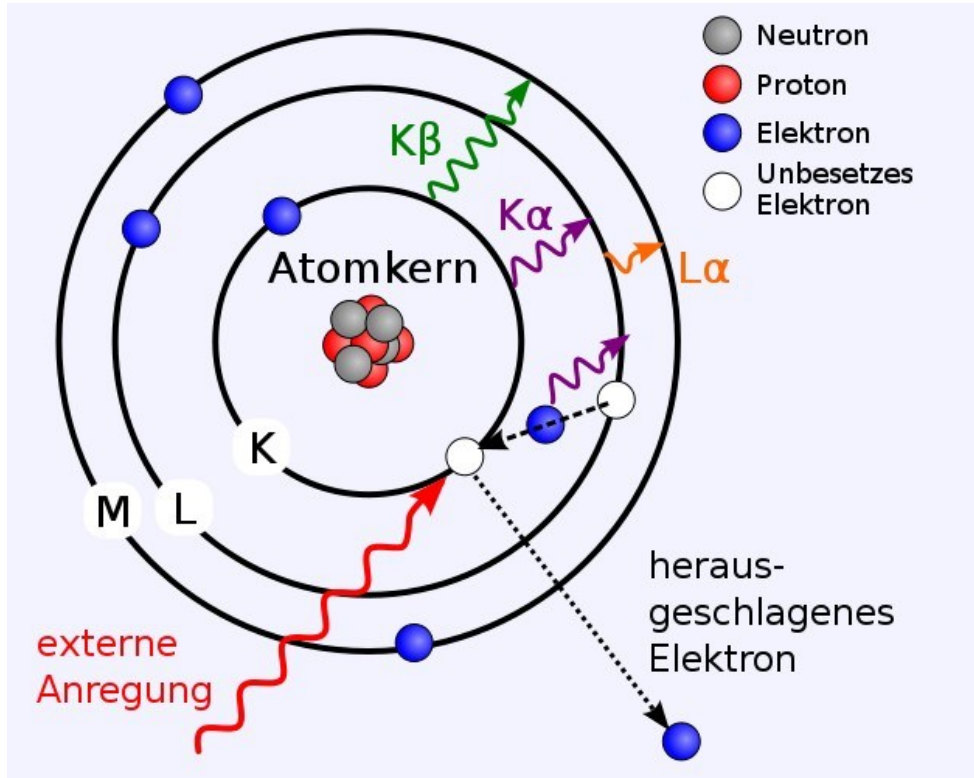
- X-ray Tube



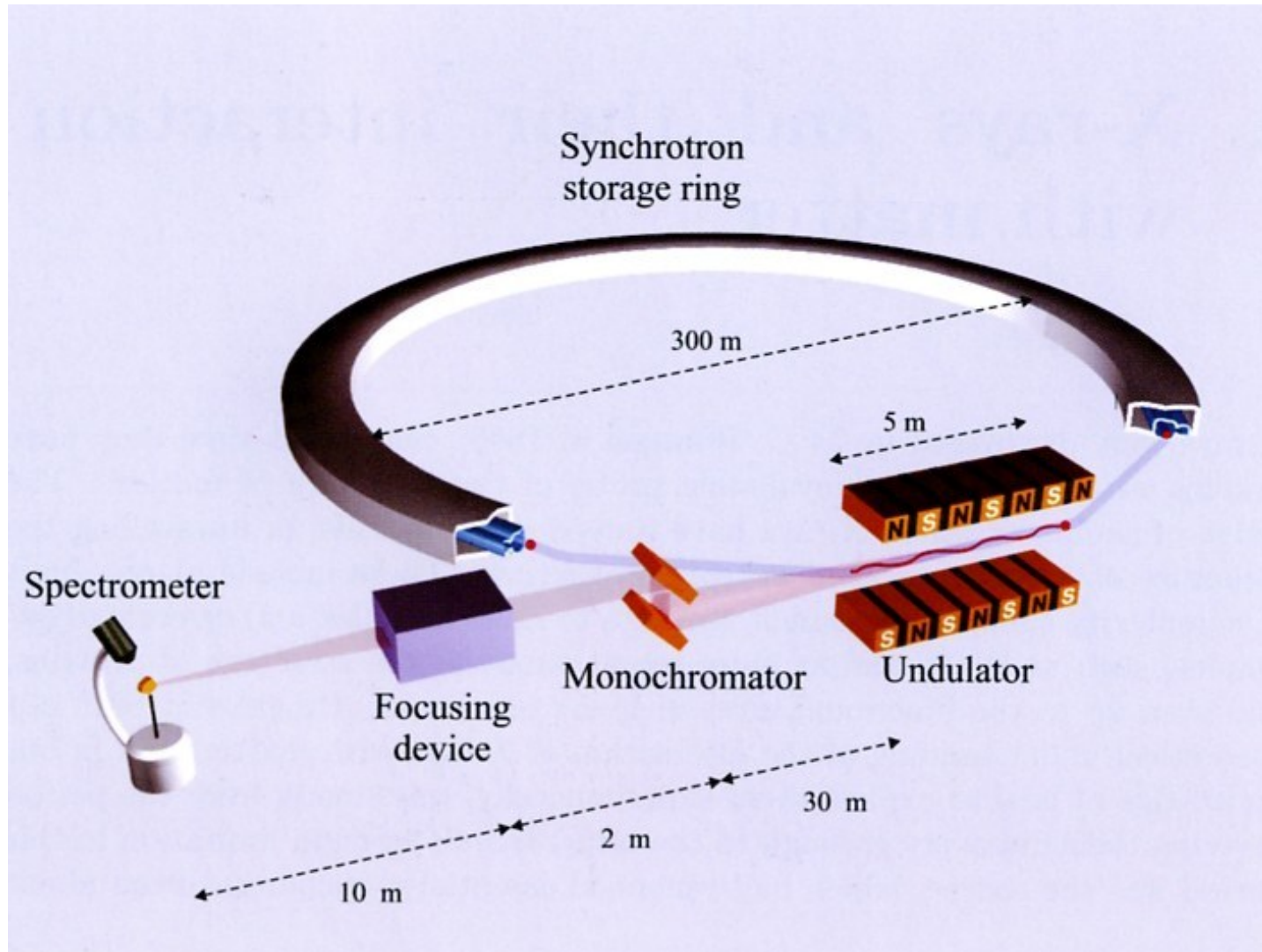
- X-ray Tube



X-ray Tube



▪ Synchrotron Radiation Storage Ring



- **Circular Accelerators**

Cyclotron

Microtron

Synchrotron

Storage Ring

▪ Cyclotron

- o Proposed in 1930 by E.O. Lawrence
- o Electrons circulate in a homogeneous magnetic field B
- o Frequency for one cycle is given by

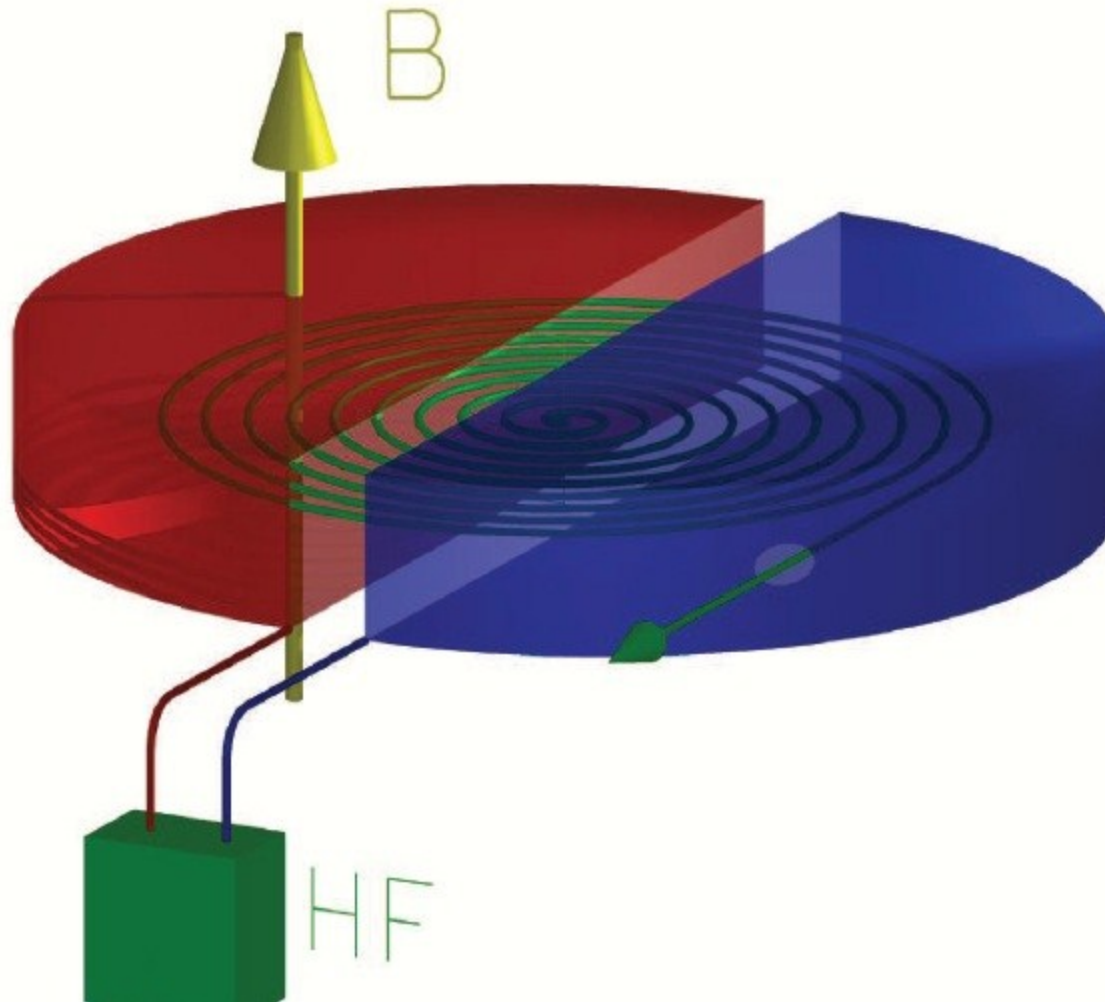
$$\omega_c = (e/m) B_z$$

- o For non-relativistic electrons ω_c is independent of the velocity v ($v/c < 0.15$)
- o At high energies the mass changes and the frequency of the field needs to be adapted.

$$\text{Example: } E_{\text{kin}} = 10 \text{ keV} = eU = m_e v_e^2 / 2 \Rightarrow v_e / c = 0.2!$$

- o Electrons at 10 keV are already relativistic!

- Cyclotron



▪ Cyclotron



Zyklotron der
Uni Bonn

▪ Microtron

- o Acceleration with a linear accelerator
- o Circular bend similar to a cyclotron
- o Bending radius R in magnetic field B for relativistic particles

Lorentz Force = Radial Force

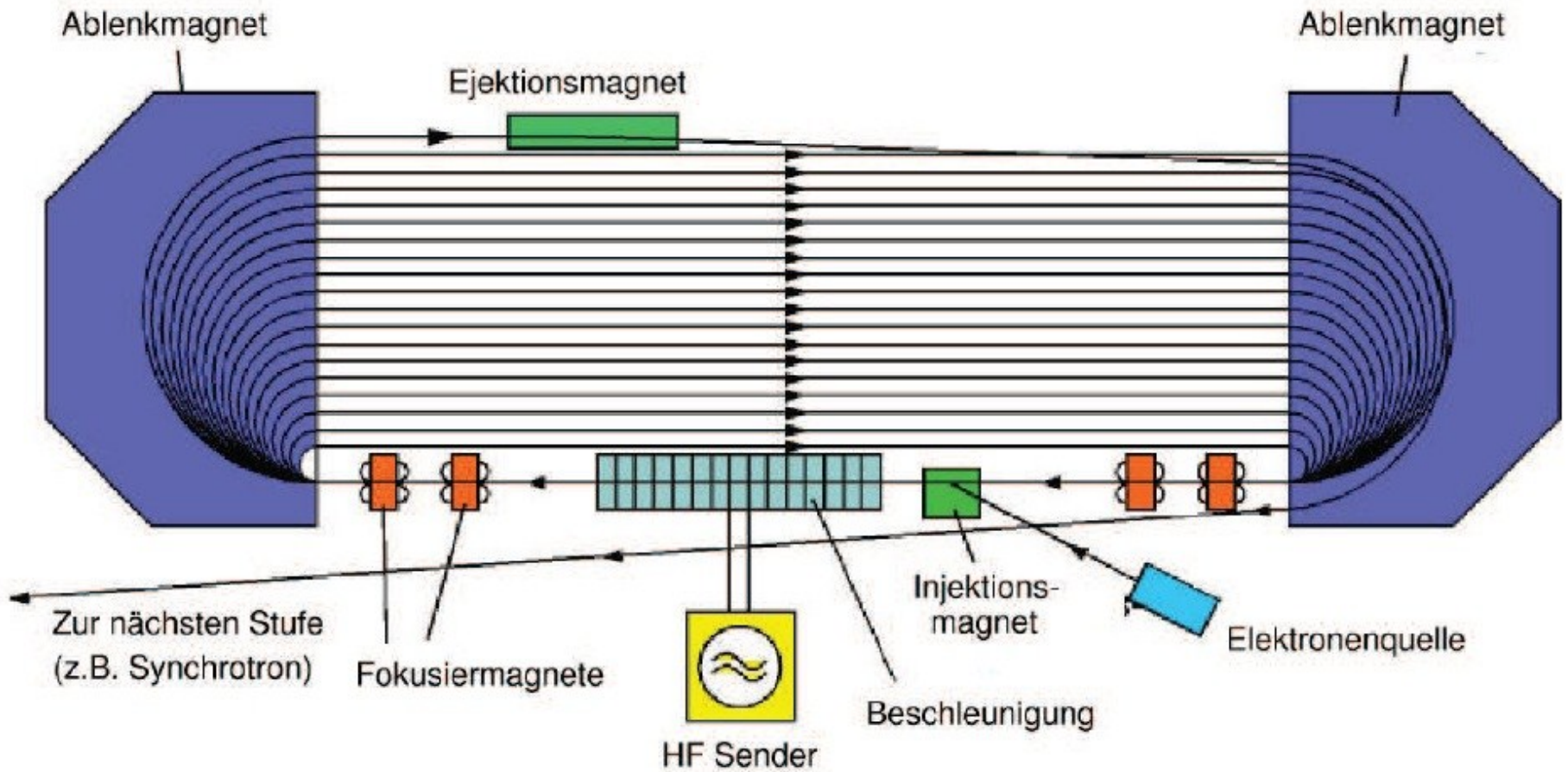
$$evB = m v^2/R$$

$$\Rightarrow R = mv/eB = vmc^2/ec^2B = (v/ec^2B) E$$

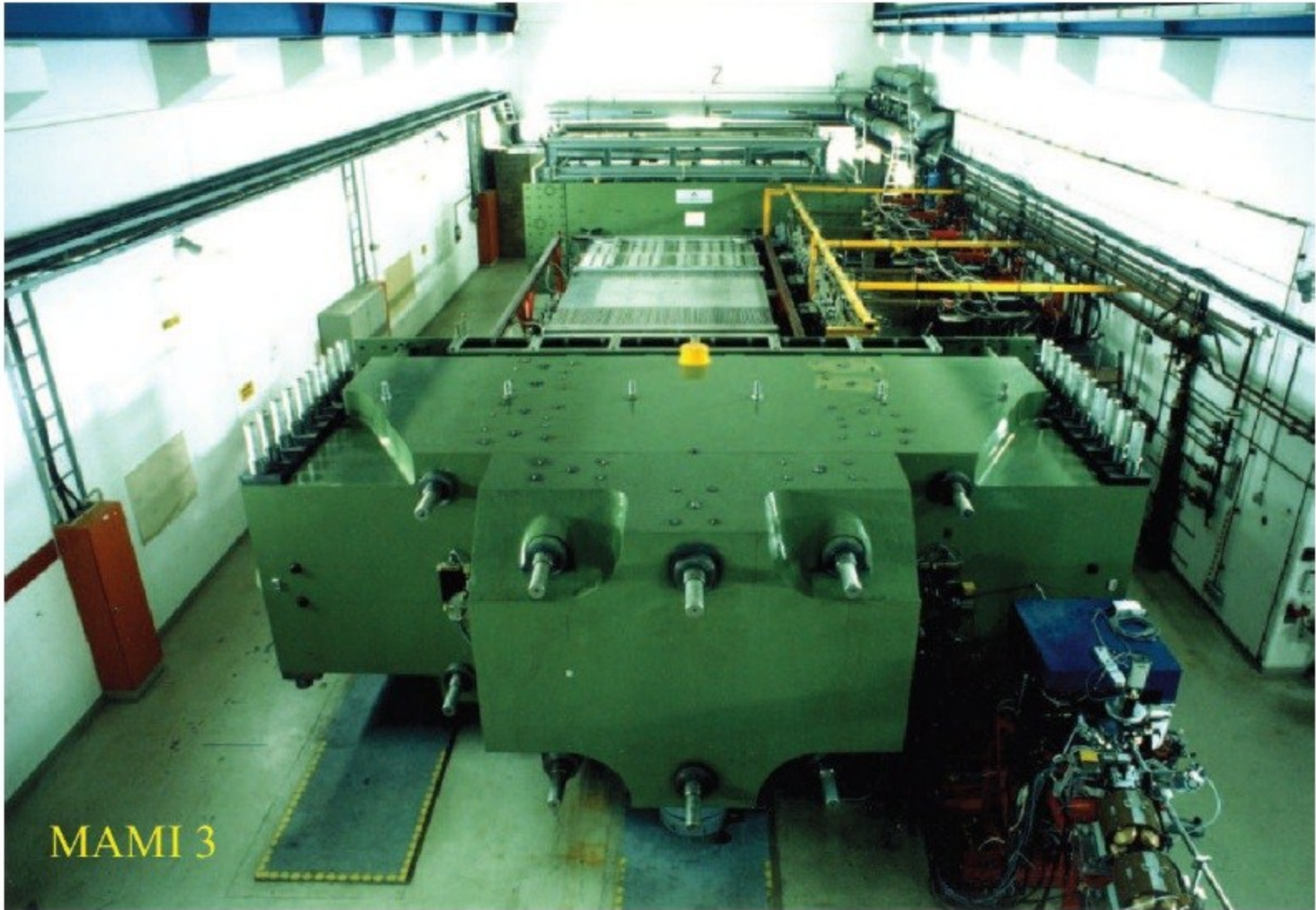
- o Acceleration such that electrons are in phase with RF field
- o Energies up to 100 MeV can be reached

Example: BESSY II

▪ Microtron



- **Microtron**



▪ Synchrotron

o For relativistic particles $v \cong c$ in a B field the radius is given by

$$R = E/ecB$$

o For $E > 1$ GeV and $B = 5$ T: $R >$ several meter.

o Technically difficult.

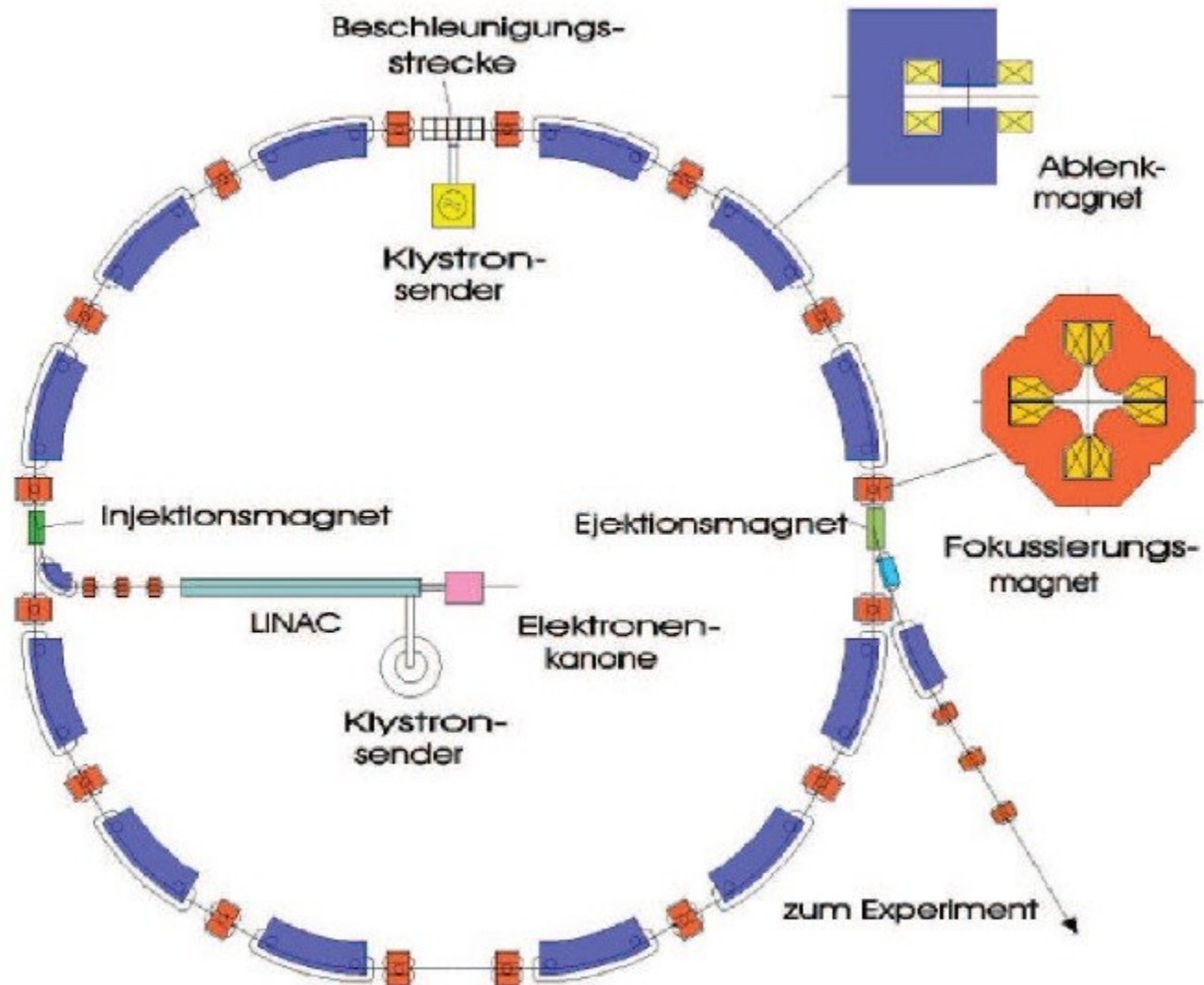
o Enforce trajectory with constant radius.

Bends in small, local magnets.

$E/B = \text{const.} \Rightarrow$ synchronous ramping of E and B

\Rightarrow Synchrotron

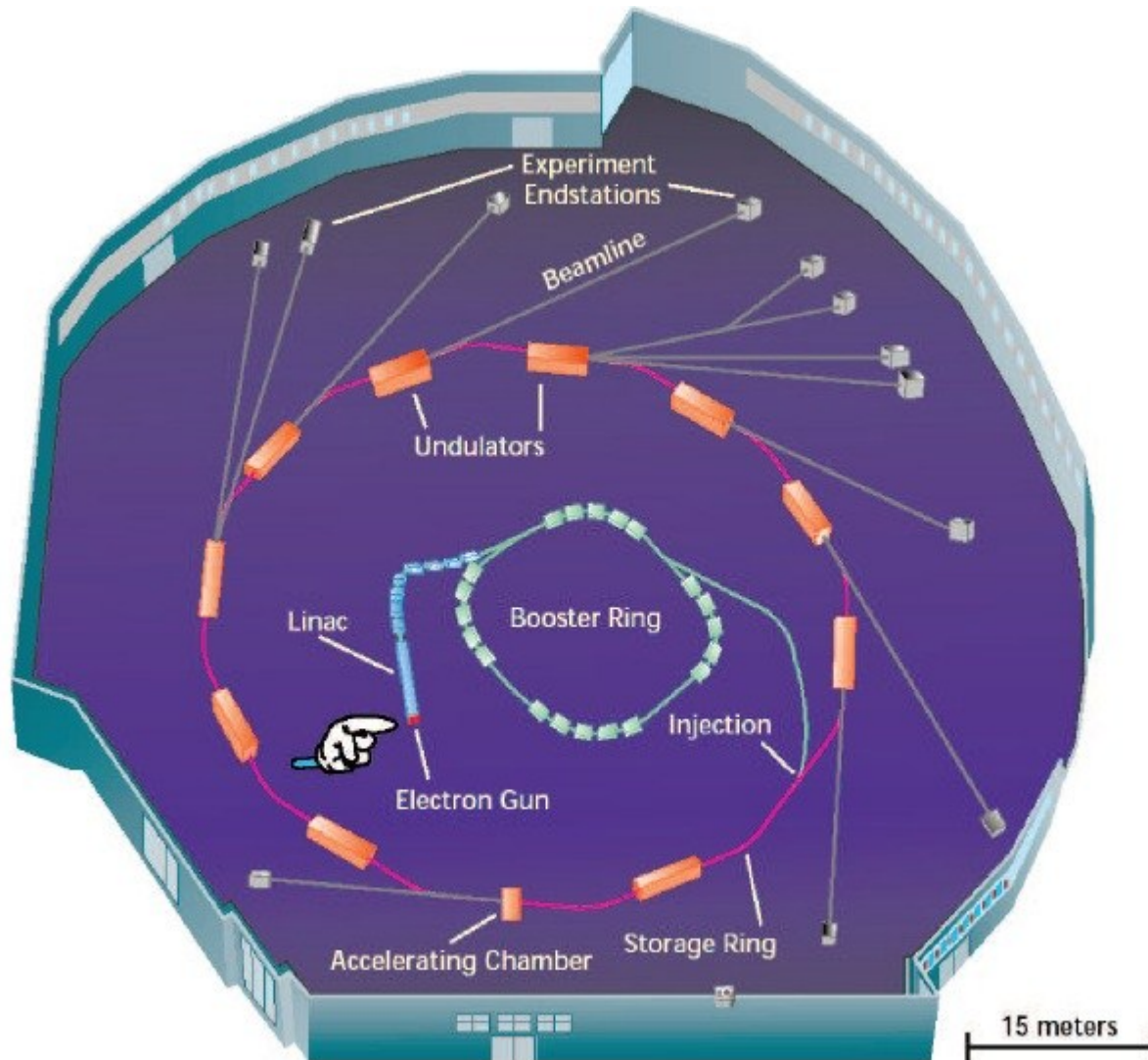
▪ Synchrotron



▪ Storage Rings

- o Modern synchrotron radiation sources are built as storage rings
- o Synchrotron cannot operate at $E=0$ since it requires $B=0$.
 - ⇒ Use LINAC or Microtron as pre-accelerator
 - Use synchrotron to reach the final energy E
 - Use storage ring to keep electrons at energy
- o The storage ring supplies the energy lost by radiation in each turn.
- o Typical parameters:
 - Lifetime: up to 30 h
 - Current: 100 – 500 mA
- o Current losses through interaction with residual gas ⇒ UHV
- o Current supplied in bunches.

Storage Rings



Storage Rings



▪ Photos machines

The three largest and most powerful synchrotrons in the world



APS, USA



ESRF, Europe-France



Spring-8, Japan



▪ Synchrotron Radiation Primer

Radiation of a non-relativistic, accelerated particle:

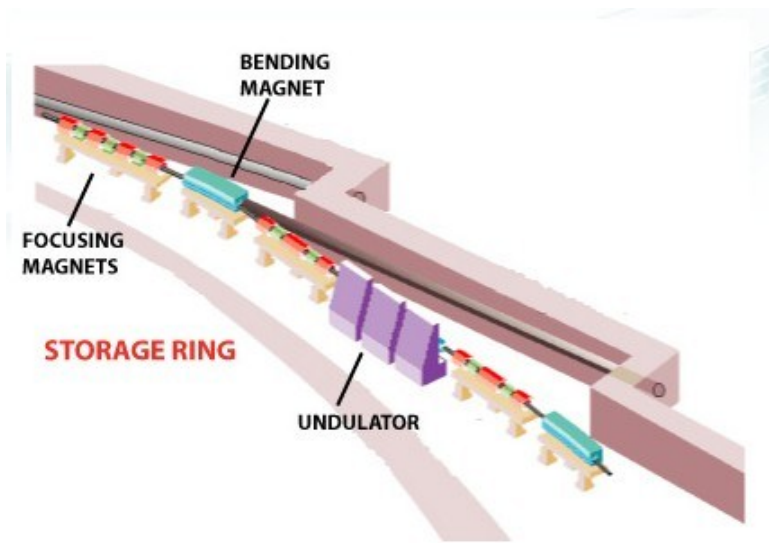
$$P = (e^2/6\pi\epsilon_0 m_0^2 c^3) (dp/dt)^2$$

Angular distribution resembles the one of a Hertz dipol:

$$(dP/d\Omega) = (e^2/16\pi^2\epsilon_0 m_0^2 c^3) (dp/dt)^2 \sin^2(\psi)$$

Radiation is emitted (similar to the dipole) in the direction perpendicular to the acceleration

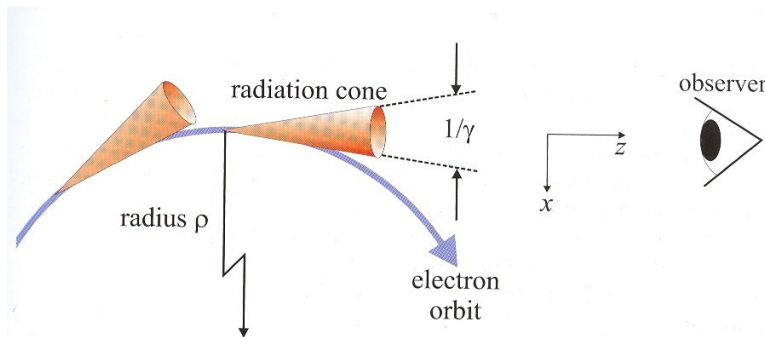
▪ Synchrotron Radiation Primer



Energy E_e of an electron at speed v :

$$E_e = mc^2 / \sqrt{1 - (v/c)^2} = \gamma mc^2$$

For 5 GeV and $mc^2 = 0.511$ MeV get $\gamma \approx 10^4$



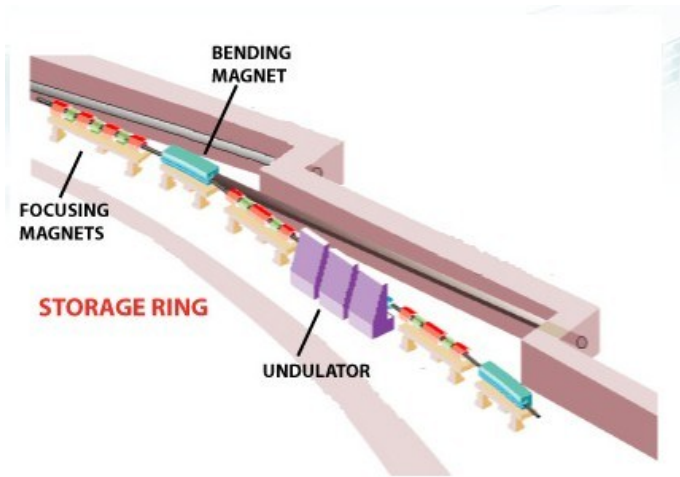
Centrifugal=Lorentz force yields for radius:

$$\rho = \gamma mc / eB = 3.3 E[\text{GeV}] / B[\text{T}] \approx 25 \text{ m}$$

$$E_e \approx 6 \text{ GeV}, B = 0.8 \text{ T}$$

Opening angle is of order $1/\gamma \approx 0.1$ mrad

Bending magnets



Characteristic energy $\hbar\omega_c$ for bend or wiggler:

$$\hbar\omega_c [\text{keV}] = 0.665 E_e^2 [\text{GeV}] B(\text{T}) \approx 20 \text{ keV}$$

$$\text{Flux} \sim E^2$$

Energy loss by synchrotron radiation per turn:

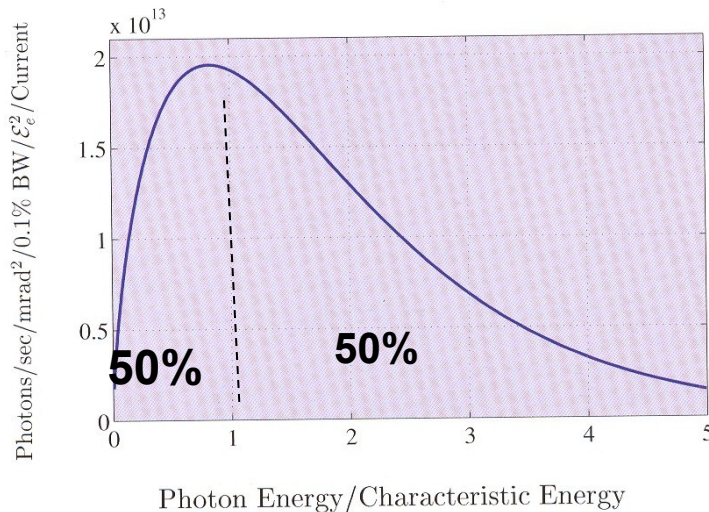
$$\Delta E [\text{keV}] = 88.5 E^4 [\text{GeV}] / \rho [\text{m}]$$

For 1 GeV and $\rho=3.33 \text{ m}$: $\Delta E = 26.6 \text{ keV/turn}$

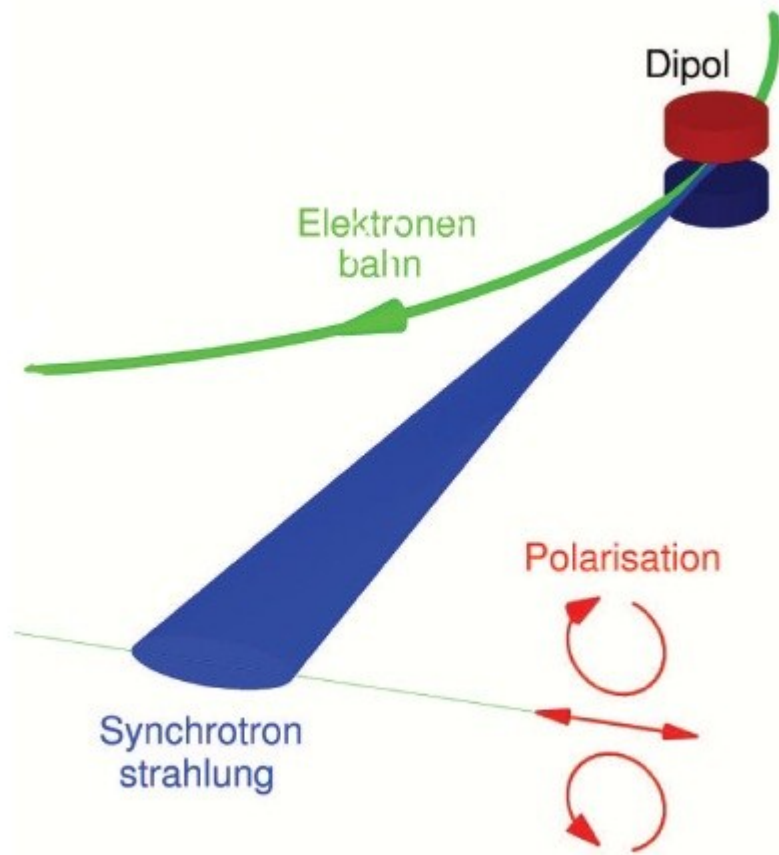
For $I=500 \text{ mA} \equiv 0.5 \text{ Cb/s} = 0.5 \times 6.25 \times 10^{18} \text{ e/s}$

$$\rightarrow P = 0.5 \times 6.25 \times 10^{18} \text{ e/s} \times 26.6 \text{ keV}$$

$$= 8.3125 \times 10^{22} \times 1.6 \times 10^{-19} = 13.3 \text{ KJ/s} = 13.3 \text{ KW}$$



▪ Polarisation



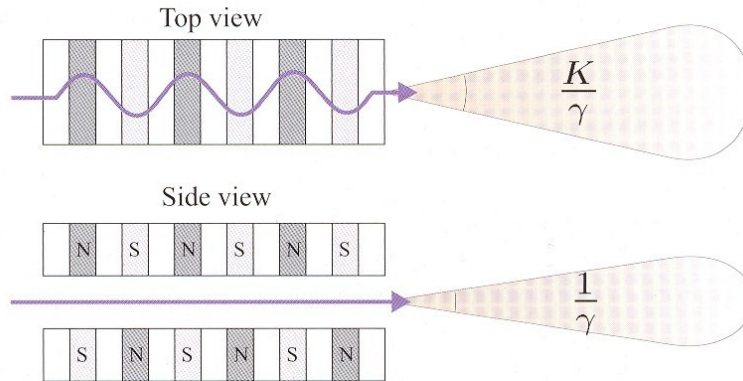
o Synchrotron radiation is polarised linearly in the plane of the orbit.

o Above and below the orbital plane the polarisation is circular.

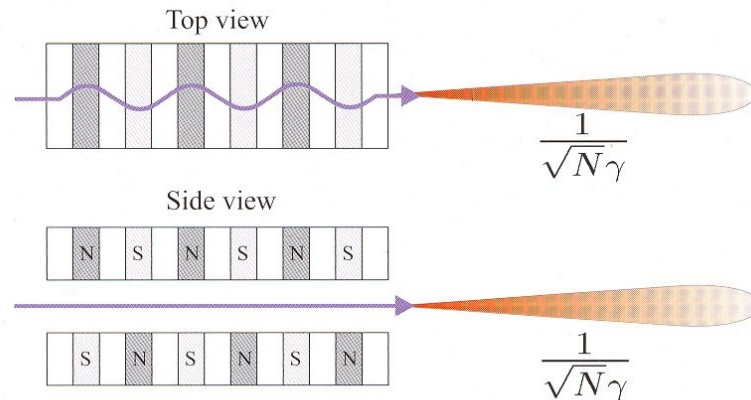
o Important applications for magnetic X-ray scattering.

Insertion Devices (wigglers and undulators)

(a) Wiggler



(b) Undulator



Wiggler:

$$P[\text{kW}] = 0.633 E_e^2 [\text{GeV}] B^2 [\text{T}] L [\text{m}] I [\text{A}]$$

$$\text{Flux} \sim E^2 \times N$$

N: number poles

Undulator:

$$k = eB / mc \quad k_u = 0.934 \lambda_u [\text{cm}] B_0 [\text{T}]$$

with λ_u undulator period

undulator fundamental:

$$\lambda_0 = \lambda_u / 2\gamma^2 \{ (1 + k^2/2 + (\gamma\theta)^2) \}$$

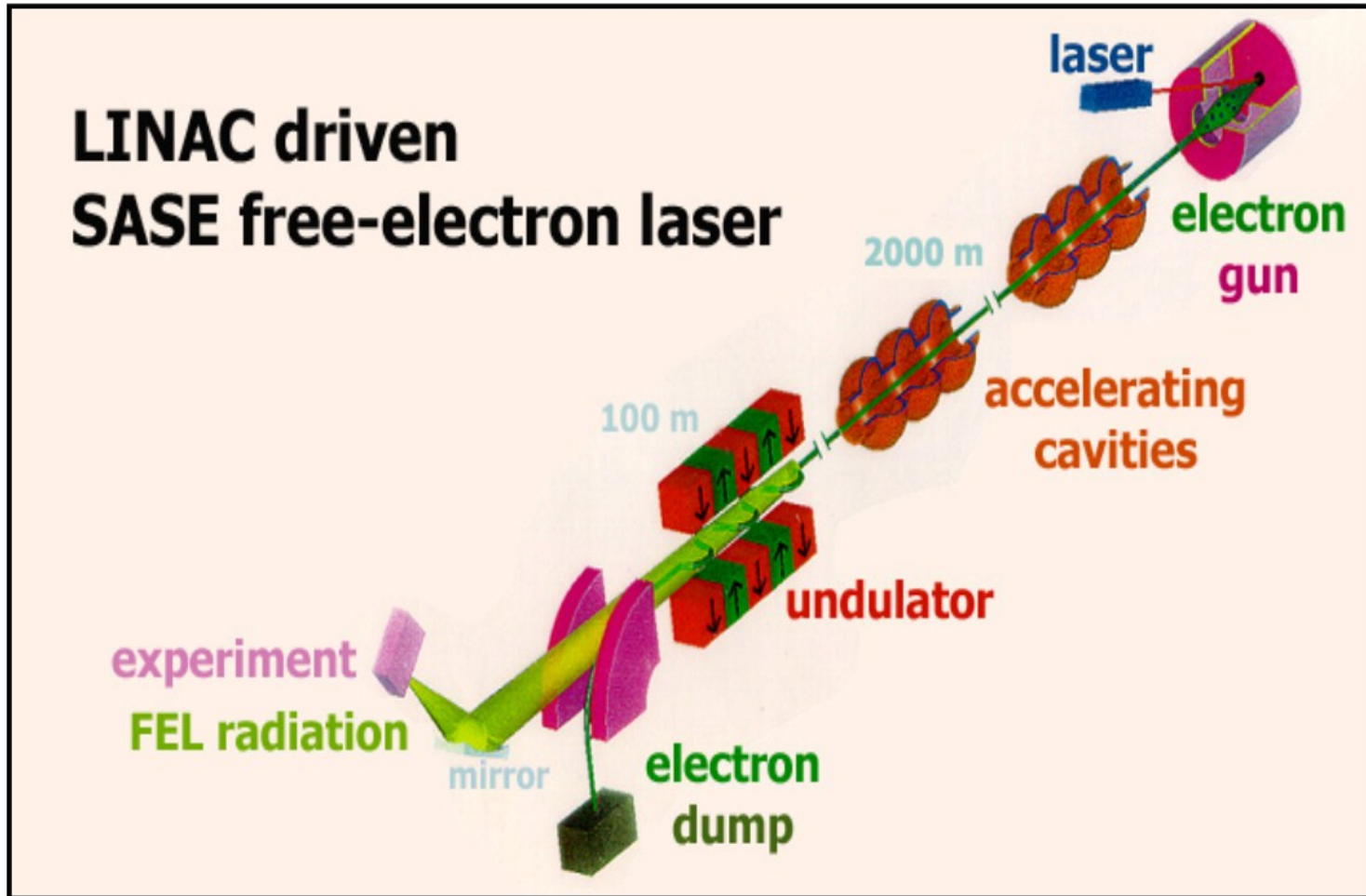
~~on axis~~

$$\text{Flux} \sim E^2 \times N^2$$

bandwidth:

$$\Delta\lambda/\lambda \sim 1/nN$$

- Free Electron Lasers (FELs)



Synchrotron and FEL sources

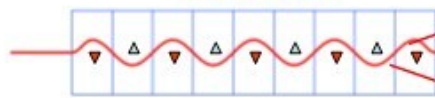
Dipole magnet
Synchrotron radiation

$$dF \sim E^2 I$$



$$\sim 2N E^2 I$$

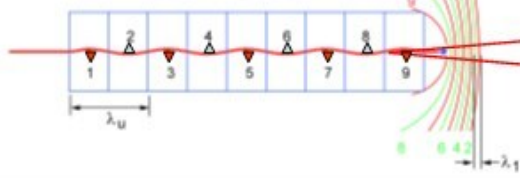
Wiggler



$$\sim N^2 E^2 I$$

$$\sim n_e$$

Undulator

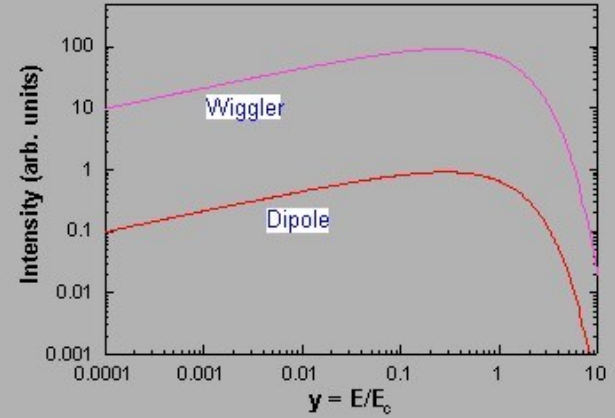


$$\sim n_e^2$$

FEL



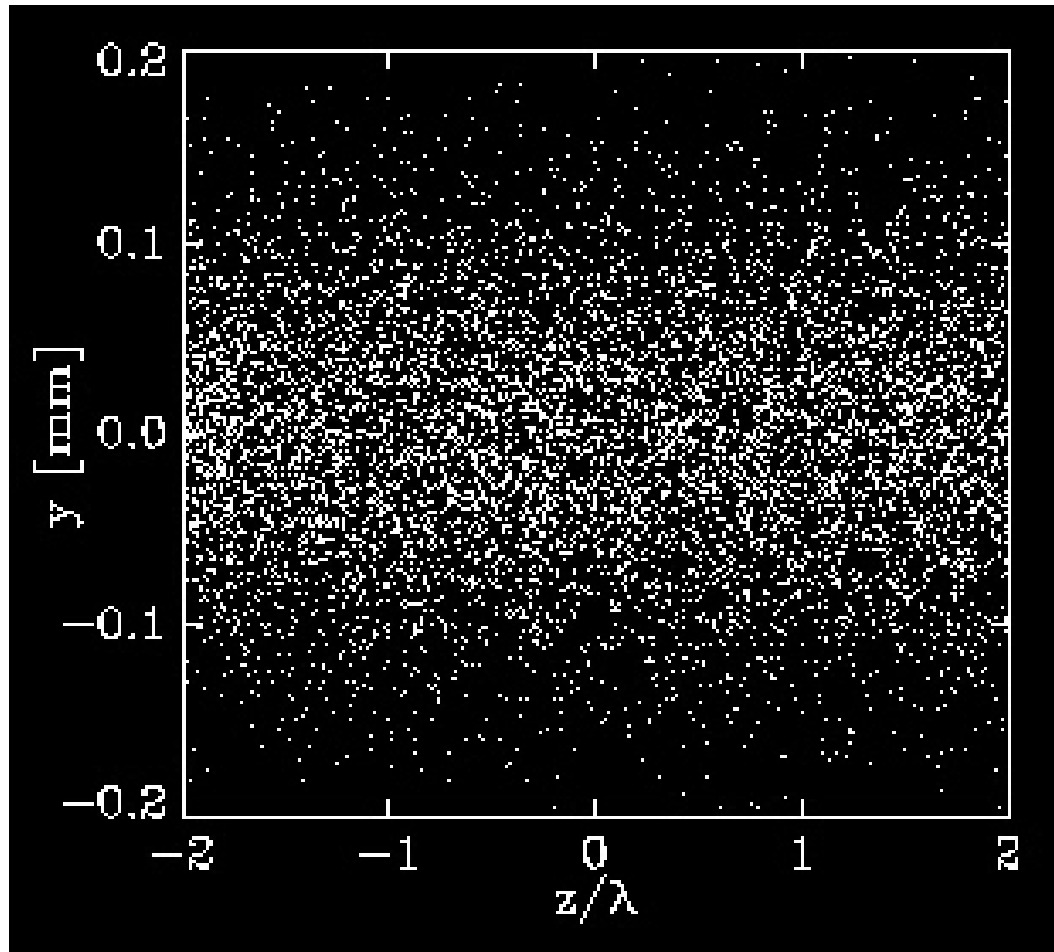
spatially coherent



The radiation emitted by a single electron in subsequent oscillations in an undulator is in phase. Radiation from different electrons is NOT (positional disorder in bunch).

“Phasing” is achieved via positional order in the bunch (micro-bunching) with a period equal to the x-ray wavelength.

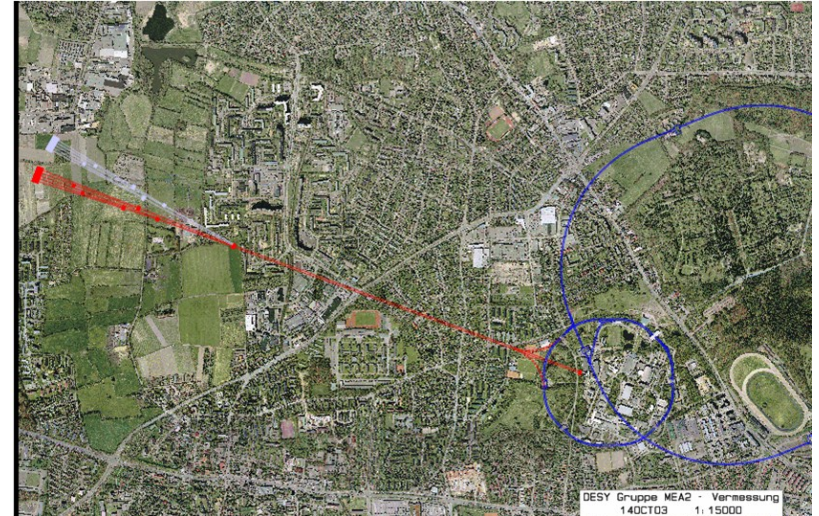
- Electron bunching



GENESIS – simulation for TTF parameters

Courtesy Sven Reiche
(UCLA)

VUV and X-Ray FELs



Brilliance

$$B = \frac{\text{photons}}{\text{s mm}^2 \text{ mrad}^2 \text{ 0.1\% BW}}$$

