

- **Coherence of light and matter:
from basic concepts to modern applications**

Part III: G. Grübel

Script IV

Coherence based X-ray techniques

Overview, Introduction to X-ray Scattering, Sources of Coherent X-rays, Speckle pattern and their analysis

Imaging techniques

Phase Retrieval, Sampling Theory, Reconstruction of Oversampled Data, Fourier Transform Holography, Applications

X-ray Photon Correlation Spectroscopy (XPCS)

Introduction, Equilibrium Dynamics (Brownian Motion), Surface Dynamics, Non-Equilibrium Dynamics

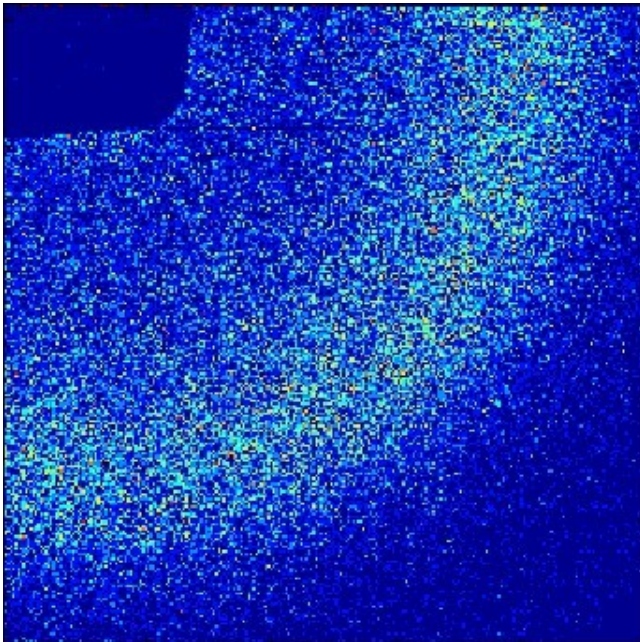
Imaging and XPCS at FEL Sources

Introduction: Scattering with coherent X-rays

If coherent light is scattered from a disordered system it gives rise to a random (grainy) diffraction pattern, known as “speckle”. A speckle pattern is an interference pattern and related to the exact spatial arrangement of the scatterers in the disordered system.

$$I(Q,t) \sim S_c(Q,t) \sim \left| \sum e^{iQR_j(t)} \right|^2$$

j in coherence volume $c = \xi_t^2 \xi_l$



Incoherent Light:

$$S(Q,t) = \langle S_c(Q,t) \rangle_{V \gg c} \quad \text{ensemble average}$$

- quantify dynamics in terms of the intensity correlation function $g_2(\mathbf{Q},t)$:

$$I(\mathbf{Q},t) = |\mathbf{E}(\mathbf{Q},t)|^2 = \left| \sum b_n(\mathbf{Q}) \exp[i\mathbf{Q} \cdot \mathbf{r}_n(t)] \right|^2$$

Note: $\mathbf{E}(\mathbf{Q},t) = \int d\mathbf{r}' \rho(\mathbf{r}') \exp [i\mathbf{Q} \cdot \mathbf{r}'(t)]$ $\rho(\mathbf{r}')$: charge density

$$g_2(\mathbf{Q},t) = \langle I(\mathbf{Q},0) \cdot I(\mathbf{Q},t) \rangle / \langle I(\mathbf{Q}) \rangle^2$$

if $\mathbf{E}(\mathbf{Q},t)$ is a zero mean, complex gaussian variable:

$$g_2(\mathbf{Q},t) = 1 + \beta(\mathbf{Q}) \langle \mathbf{E}(\mathbf{Q},0) \mathbf{E}^*(\mathbf{Q},t) \rangle^2 / \langle I(\mathbf{Q}) \rangle^2$$

$\langle \rangle$ ensemble av.; $\beta(\mathbf{Q})$ contrast

$$g_2(\mathbf{Q},t) = 1 + \beta(\mathbf{Q}) |f(\mathbf{Q},t)|^2$$

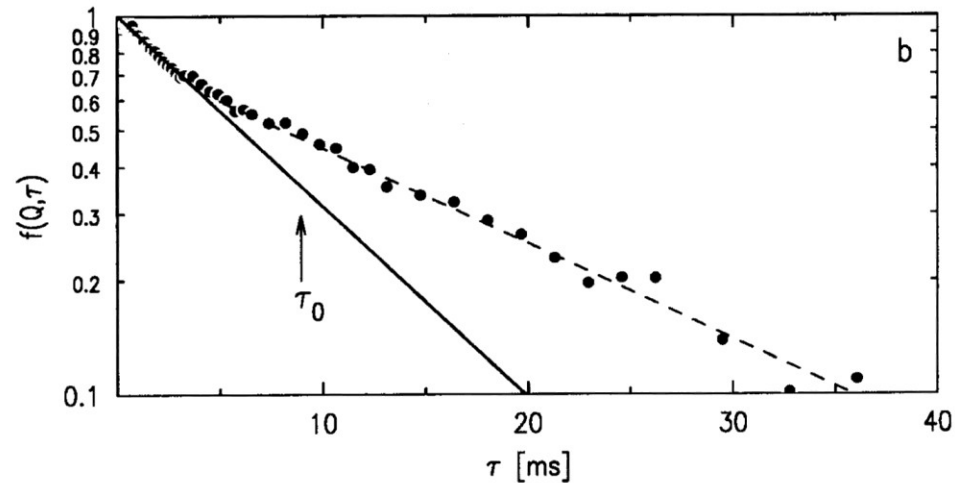
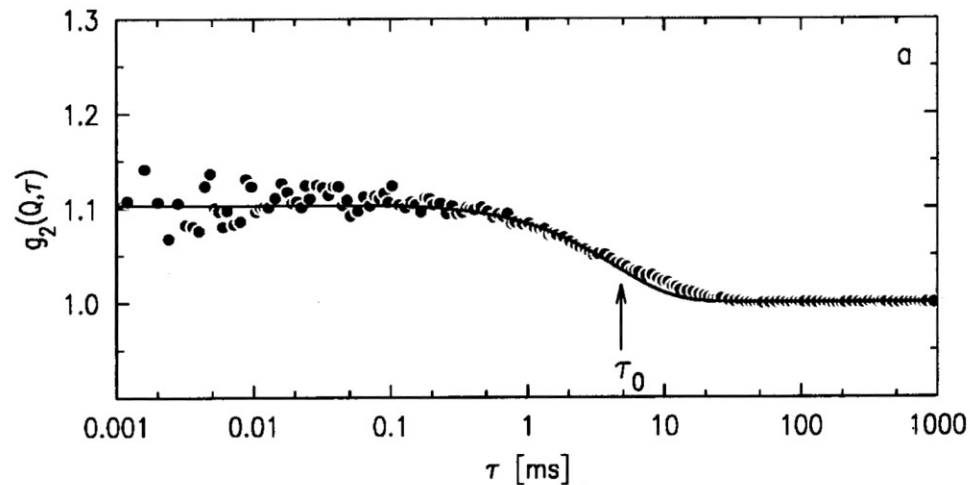
with $f(\mathbf{Q},t) = F(\mathbf{Q},t) / F(\mathbf{Q},0)$
 $F(\mathbf{Q},0)$: static structure factor
 N : number of scatterers

$$F(\mathbf{Q},t) = [1/N \{b^2(\mathbf{Q})\}] \left| \sum_{m=1}^N \sum_{n=1}^N \langle b_n(\mathbf{Q}) b_m(\mathbf{Q}) \bullet \exp\{i\mathbf{Q}[\mathbf{r}_n(0) - \mathbf{r}_m(t)]\} \rangle \right|$$

- A time correlation function $g_2(\mathbf{Q}, \tau)$

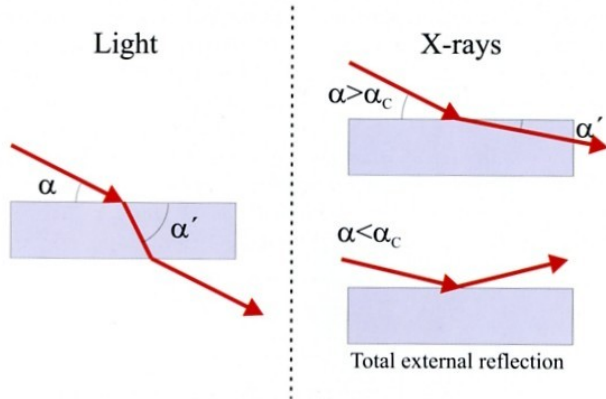
$$g_2(\mathbf{Q}, t) = 1 + \beta(\mathbf{Q}) |f(\mathbf{Q}, t)|^2 \quad \text{and} \quad f(\mathbf{Q}, t) = \exp(-\Gamma t) = \exp(-t/\tau)$$

$\beta(\mathbf{Q})$ \updownarrow

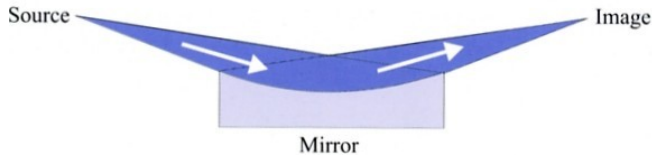


Surface Sensitivity of X-Rays

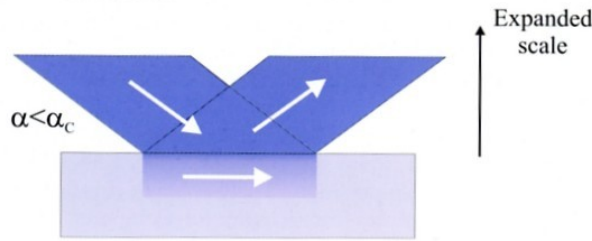
(a) Refraction and reflection of light and X-rays



(b) Focussing X-ray mirror



(c) Evanescent wave



refractive index:

$$n = 1 - \delta + i\beta = 1.5 - 1.8 \text{ visible light}$$

X-rays: $\delta = 10^{-5}$ (solid matter)

$$\beta \ll \delta$$

$$n < 1$$

$$\cos \alpha = n \cdot \cos \alpha'$$

total external reflection for: $\alpha < \alpha_c$

for $\alpha' = 0$ (and $\beta=0$):

$$\alpha_c = \sqrt{2\delta} \text{ [mrad]}$$

$\alpha < \alpha_c$: evanescent wave
with nanometer penetration

Dynamics at Surfaces: Capillary Waves

Thermally excited capillary waves decorate the surfaces of all liquids depending on the surface tension σ and the viscosity η .

Capillary wave:

$$\zeta(r,t) = \zeta_0 \exp(iq_{\parallel}r - \omega t)$$

$$\omega = \omega_0 + i\Gamma$$

The linear Navier-Stokes equation for simple liquids yields:

Low viscosity: $\sigma\rho/4\eta^2q_{\parallel} > 1$

$$\Gamma = (2\eta/\rho) q_{\parallel}^2$$

$$\omega_0 = \sqrt{(\sigma/\rho)} q_{\parallel}^{3/2}$$

propagating wave

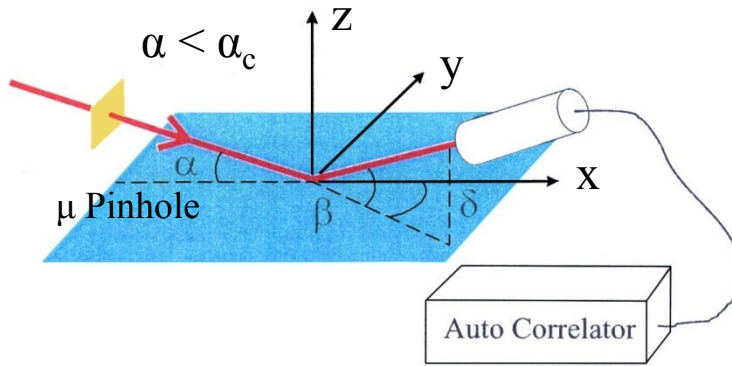
High viscosity: $\sigma\rho/4\eta^2q_{\parallel} < 1$

$$\Gamma = (\sigma/\eta) q_{\parallel}$$

overdamped wave

Dynamics at Surfaces: Capillary Waves

X-Ray Photon Correlation Spectroscopy in Surface Geometry



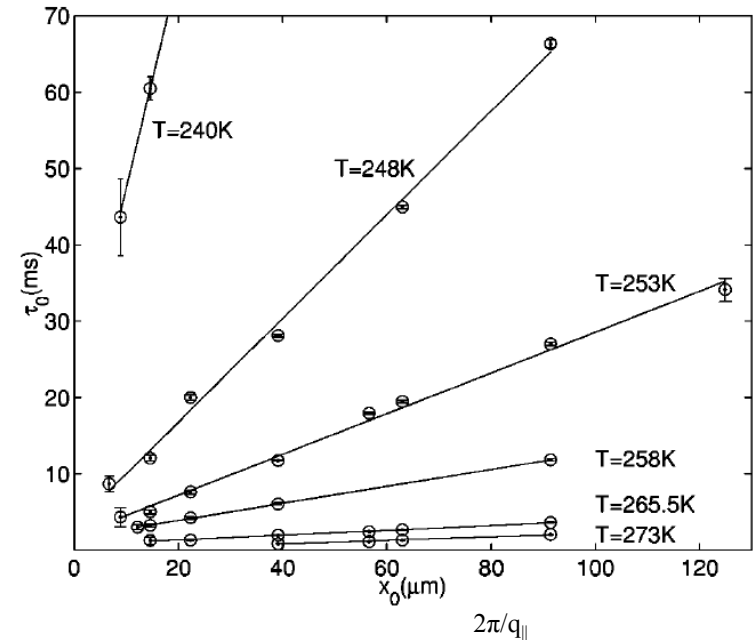
$\alpha = \beta$: reflectivity $q_z = (2\pi/\lambda) 2 \sin \alpha$

$\alpha_i = \alpha_c$: 50 - 100Å penetration surface sensitivity

$\alpha \neq \beta$: $q_{||} = 2\pi/\lambda(\cos \beta - \cos \alpha)$ in-plane correlations

Glycerol: a „prototypical“ glassformer

Seydel, Madsen, Tolan, Grübel, Press, PRB63, 73409 (2001)

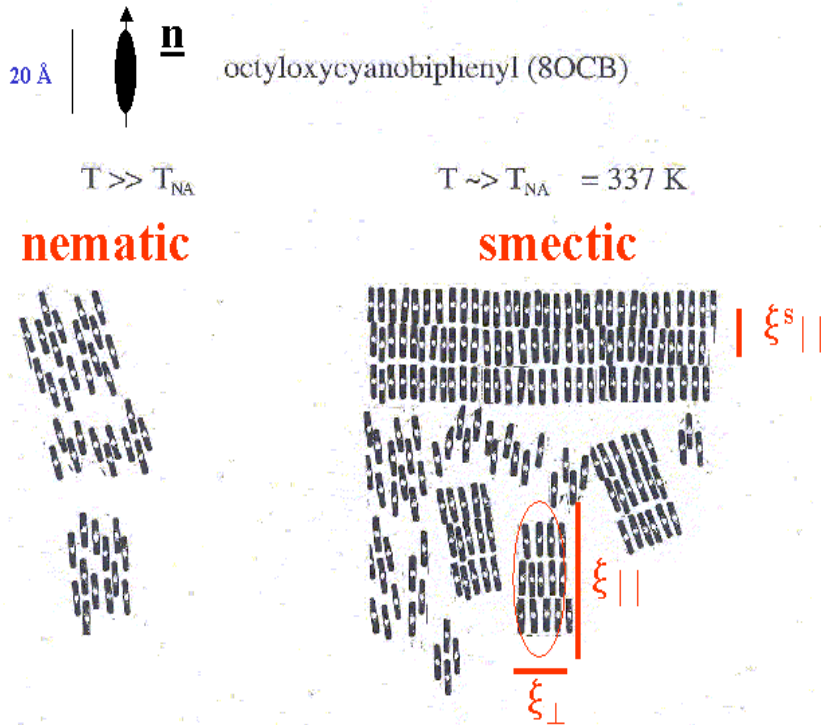


overdamped regime:

$$\Gamma = (\gamma/2\eta)k; \quad \tau_0 = \{\eta(T)/\pi \gamma(T)\}x_0$$

Viscosity of a liquid crystal near the nematic to smecticA transition

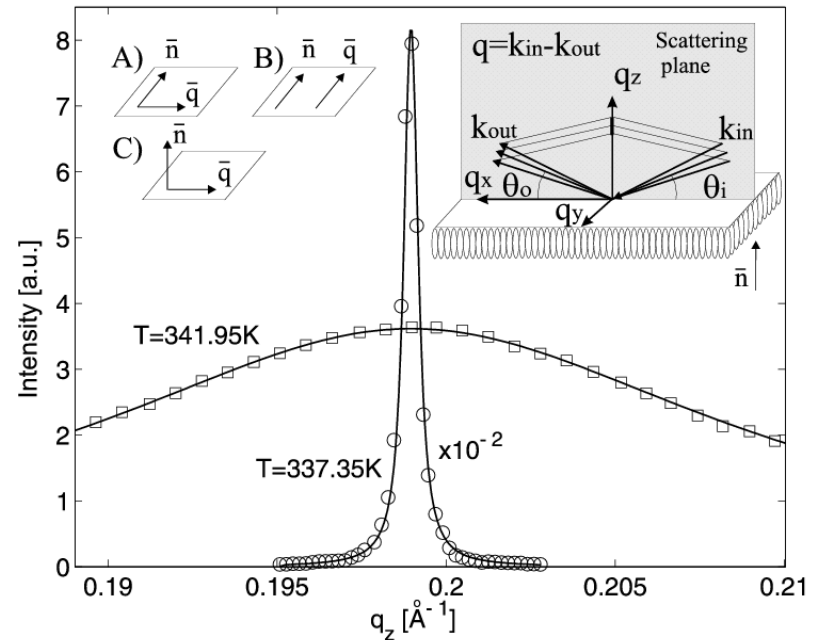
A. Madsen, J. Als-Nielsen and G. Grübel, PRL,90,85701(2003)



$$\xi_{\parallel} \sim \left\{ \frac{T - T_{NA}}{T_{NA}} \right\}^{\nu_{\parallel}} = \xi_{\parallel}^S$$

$$\xi_{\perp} \sim \left\{ \frac{T - T_{NA}}{T_{NA}} \right\}^{\nu_{\perp}}$$

$$\nu_{\parallel} = 0.71; \nu_{\perp} = 0.58$$



References: Pershan and Als-Nielsen, PRL52,759(1984)
 Pershan, Braslau, Weiss, Als-Nielsen, PRA35,4800(1987)

▪ Viscosity of a Liquid Crystal near the Nematic-SmecticA Transition

Dynamics:

viscosity is anisotropic: η_1, η_2, η_3

depending on the relative orientations:

$\mathbf{n}, \mathbf{v}, \nabla \cdot \mathbf{v}$

described by Leslie coefficients $\alpha_1 \dots \alpha_5$,
or parameters $v_1 \dots v_5$ (Harvard notation)
($v_4 = v_5 = 0$ for incompressible fluids)

Theory: (N-SmA transition)

$$\eta_1 \sim \exp(E_A/kT)$$

$$\eta_2 \sim \exp(E'_A/kT)$$

$$\eta_3 = c(T - T_{NA}/T_{NA})^{-\beta} + \text{non.div.}$$

Predictions:

$$\beta = 3v_{\parallel} - 2v_{\perp} \quad [1]$$

$$\beta = 1/3 \quad [2]$$

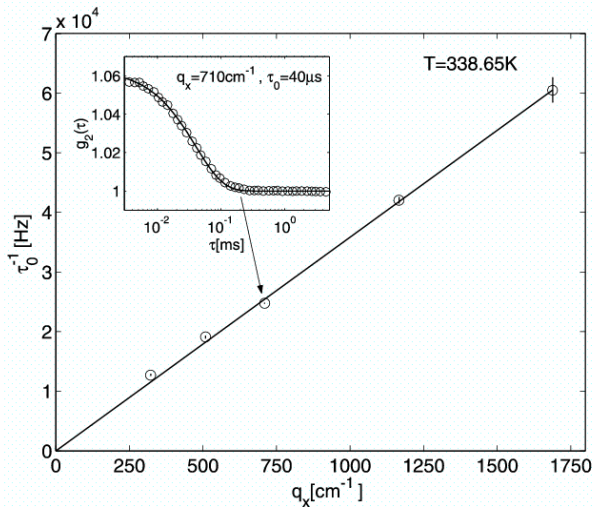
$$\beta = 1/2 \quad [3]$$

[1] Hossein, Swift, Chen,
Lubensky, PRB19,432(1979)

[2] Jähnig, Brochard,
J.Phys., 35, 301(1974);
deGennes, Sol. State
Comm., 10, 753 (1972)

[3] Langevin, J.Phys. 37,
101(1975)

Viscosity of a Liquid Crystal near the Nematic-SmecticA Transition



XPCS: overdamped modes

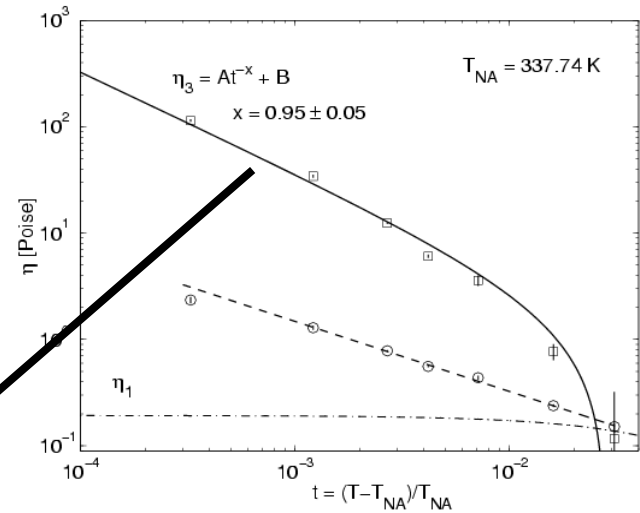
$$\eta_3 \sim (T - T_{NA} / T_{NA})^{-\beta}$$

$$\beta = 0.95(2)$$

in agreement with theory [1]:

$$\beta = 3v_{\parallel} - 2v_{\perp} = 0.94$$

$$v_{\parallel} = 0.70; v_{\perp} = 0.58 \text{ (static data)}$$



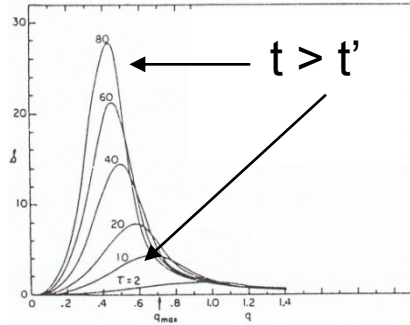
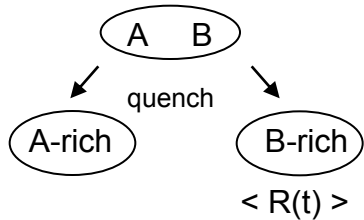
η / σ diverging

σ constant [1]; data reflect temperature dependence of viscosity

- Out-plane movements ($\zeta \parallel n$) within the smectic layers are strongly damped (η_3 critical).
- The smectic layers remain viscous in-plane (η_2 non-critical)

Non-Equilibrium Dynamics

Domain coarsening in phase separating systems (glasses, alloys,...) e.g. after quenching, aging...



Phase – separating Glass

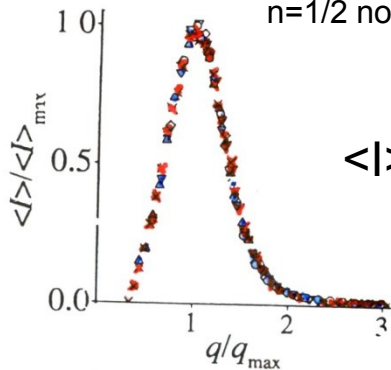
Malik *et al.*, PRL 81, 5832, 1998



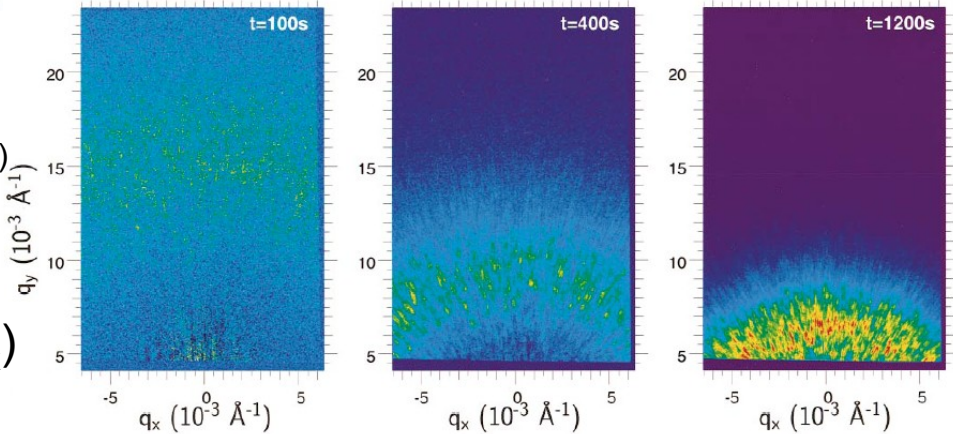
quench
 (B_2O_3) -rich (SiO_2) -rich $943\text{K} < T < 963\text{K}$

Dynamic Scaling: $\langle R(t) \rangle \sim t^n$

$n=1/3$ conserved order parameter (model B)
 $n=1/2$ non-cons. order parameter (model A)



$$\langle I \rangle(q,t) / \langle I \rangle_{max} = F(q/q_{max})$$



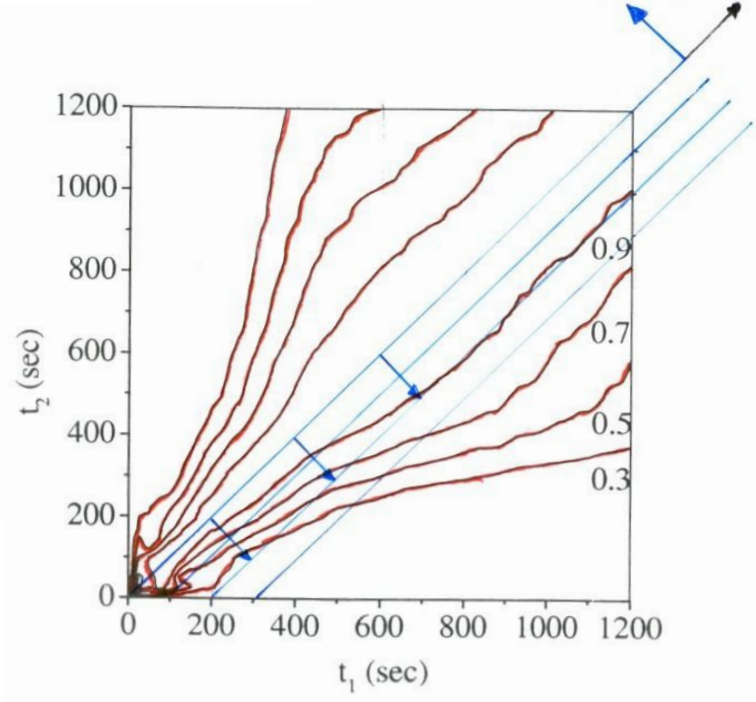
$$\langle R(t) \rangle \sim t^n \quad n=1/3$$

XPCS: investigate fluctuations around the average scaling behaviour: $\tau = \tau(q,t)$

Two time correlation function:

$$C(q, t_1, t_2) = \frac{\langle I(t_1) I(t_2) \rangle - \langle I(t_1) \rangle \langle I(t_2) \rangle}{[\langle I^2(t_1) \rangle - \langle I(t_1) \rangle^2]^{1/2} [\langle I^2(t_2) \rangle - \langle I(t_2) \rangle^2]^{1/2}}$$

$$\Delta t = t_1 - t_2 \quad \bar{t} = \frac{t_1 + t_2}{2}$$



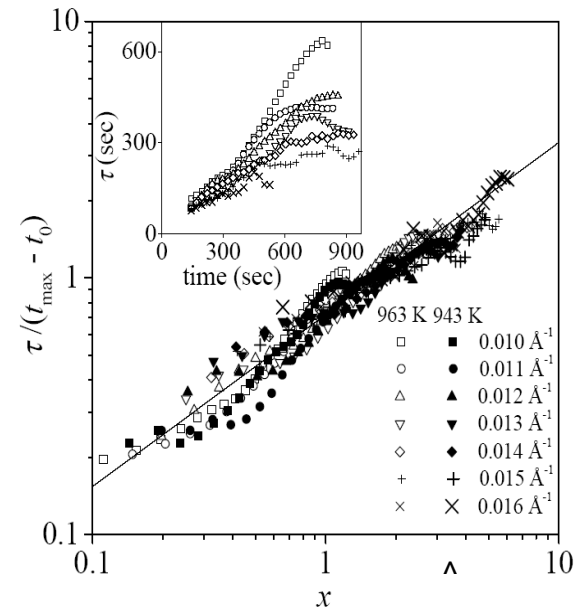
$$\tau = \tau(t)$$

Fluctuations $\tau = \tau(q, t)$

Prediction: Brown et al. PRE 56,6601,1997

$$\tau(q, t) = [t_{\max}(q) - t_0] \left\{ a \frac{[t - t_0]}{[t_{\max}(q) - t_0]} \right\}^{(1-n)}$$

$$\sim 1/q \ t^{2/3}$$

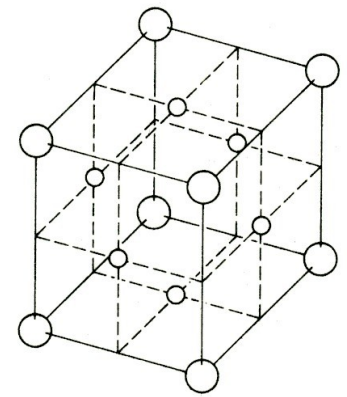


$$a = 0.72(2) \quad (1-n) = 0.65(4) = 1 - 1/3$$

Phase-Ordering in Cu_3Au

high T: fcc sites occupied by either Cu or Au

$T \leq T_c = 383 \text{ C}$: ordering with Au on corner and Cu on face sites
4-fold degenerate ground-state

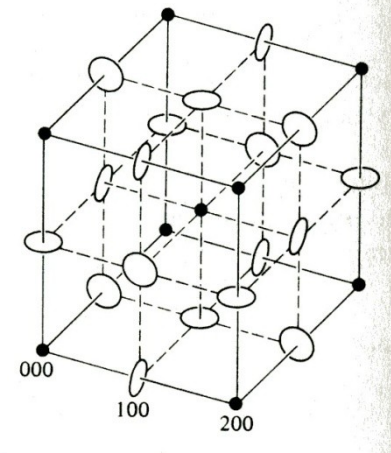


(corners can be chosen in 4 different ways)

groundstates separated by domain walls

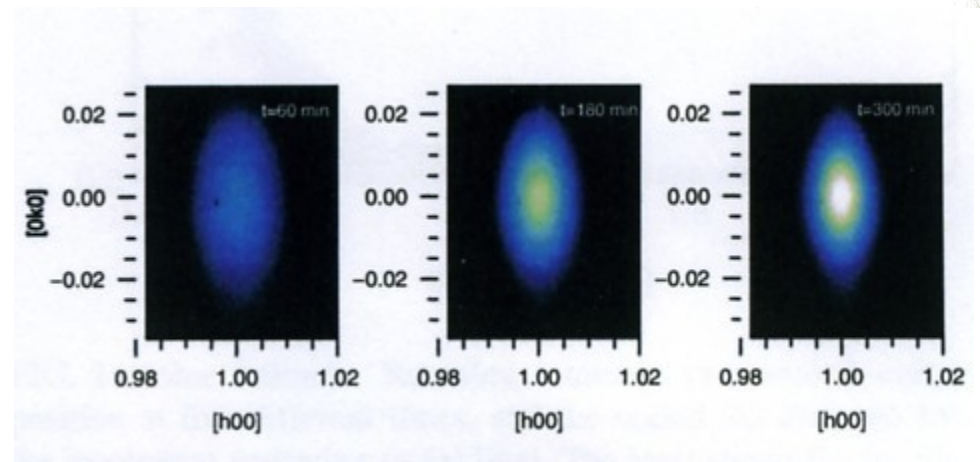
domain walls give rise to ellipse shaped superlattice reflections of type $[100]$

quench: domain formation and growth in disordered phase
domain coarsening with $R \sim t^{1/2}$



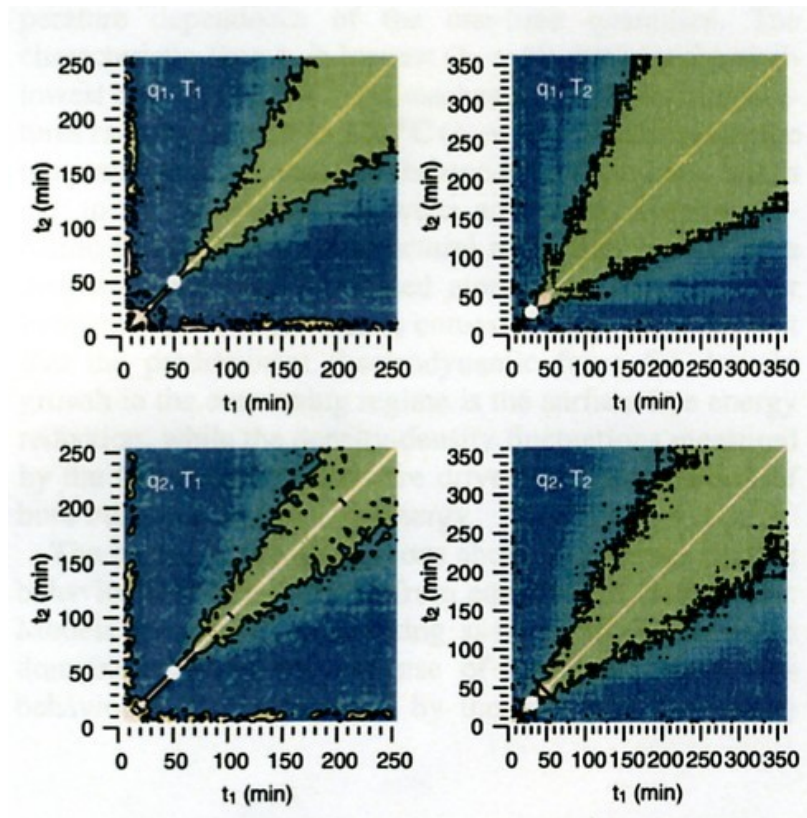
$[100]$ superlattice reflection
after quench from 425 C to
370 C

Fluerasu&Sutton, PRL94(2005)55501



- Study fluctuations about the average behaviour: XPCS characterize by two-time correlation function

$$C(q, t_1, t_2) = \frac{\langle I(t_1) I(t_2) \rangle - \langle I(t_1) \rangle \langle I(t_2) \rangle}{[\langle I^2(t_1) \rangle - \langle I(t_1) \rangle^2]^{1/2} [\langle I^2(t_2) \rangle - \langle I(t_2) \rangle^2]^{1/2}}$$



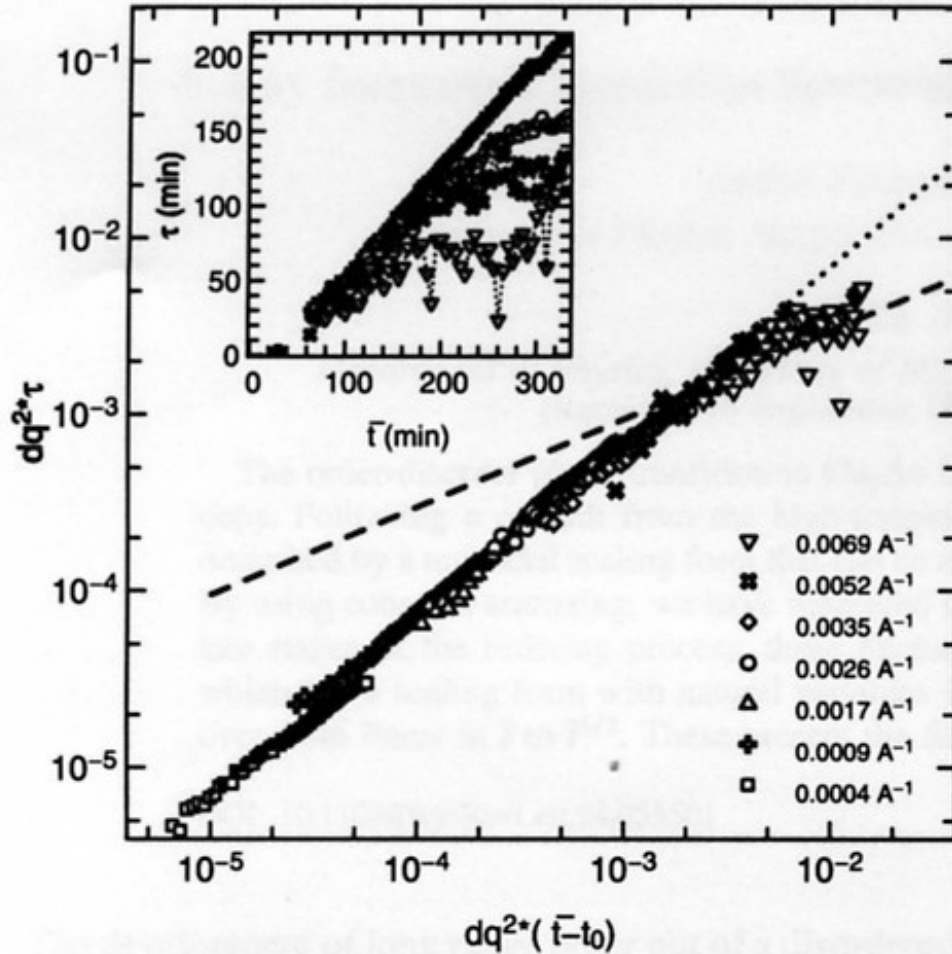
rescaled correlation time $\tau \sim t_{\text{mean}}$

in the low t_{mean} limit

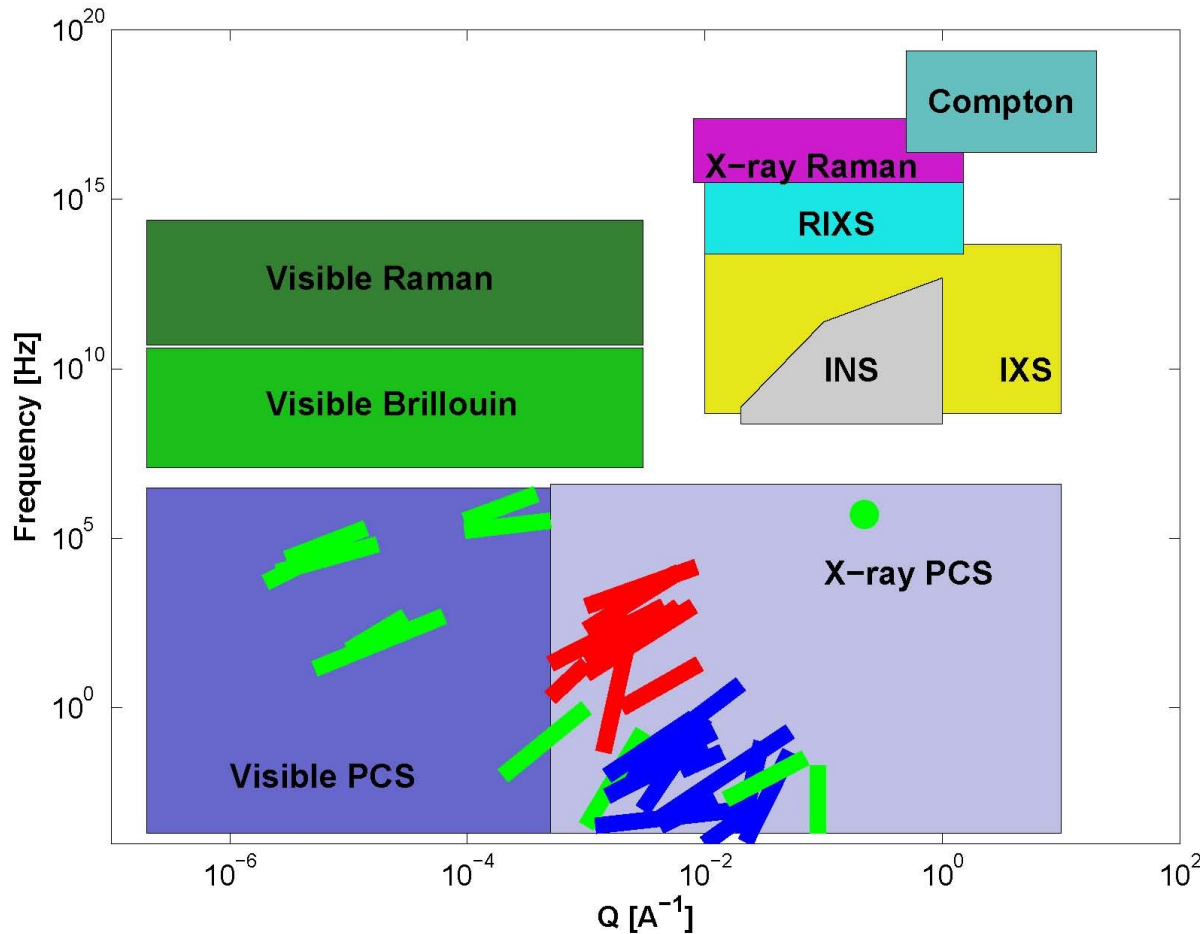
$$\sim t_{\text{mean}}^{1/2}$$

in the high t_{mean} limit

with $dq = Q - [100]$



▪ XPCS – operating range



XPCS

- access to large momentum transfers ($Q_{\text{max}}=2\pi \cdot \sin\theta/\lambda$) or short lengthscales
- not subject to multiple scattering
- can be combined with the surface sensitivity of X-rays

▪

Imaging and XPCS at an FEL source

$$F_c = (\lambda/2)^2 \bullet B(\text{rilliance}) \quad B(\text{storage ring}) \approx 10^{20} - 10^{21}$$

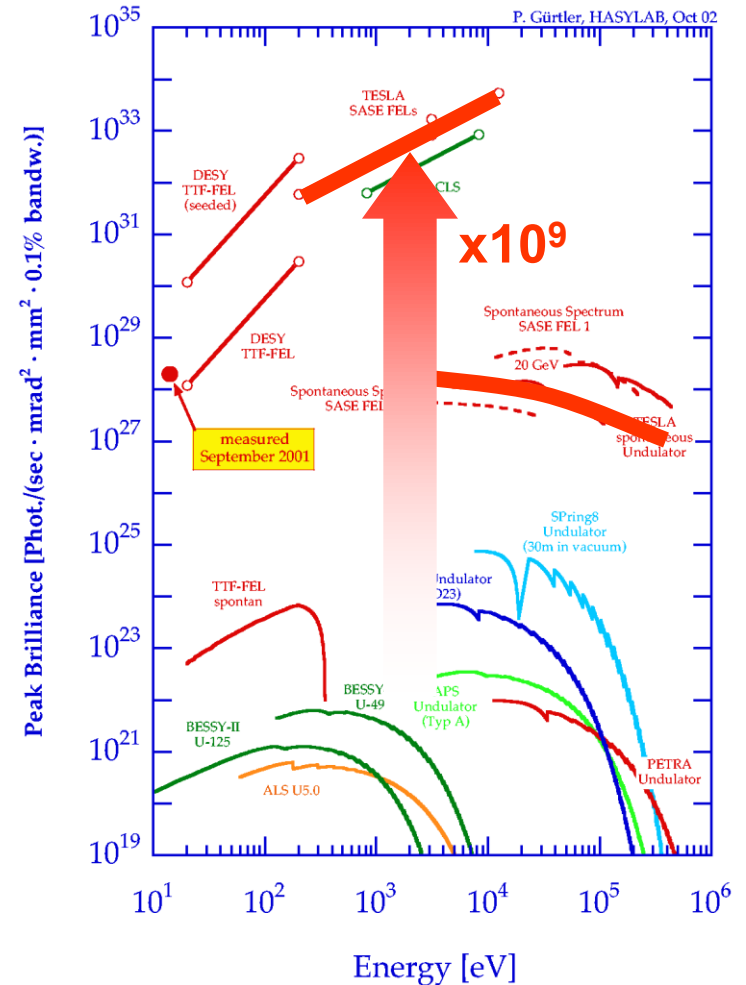
Length of a storage ring pulse \approx 50-100 ps

(X)FEL key parameters

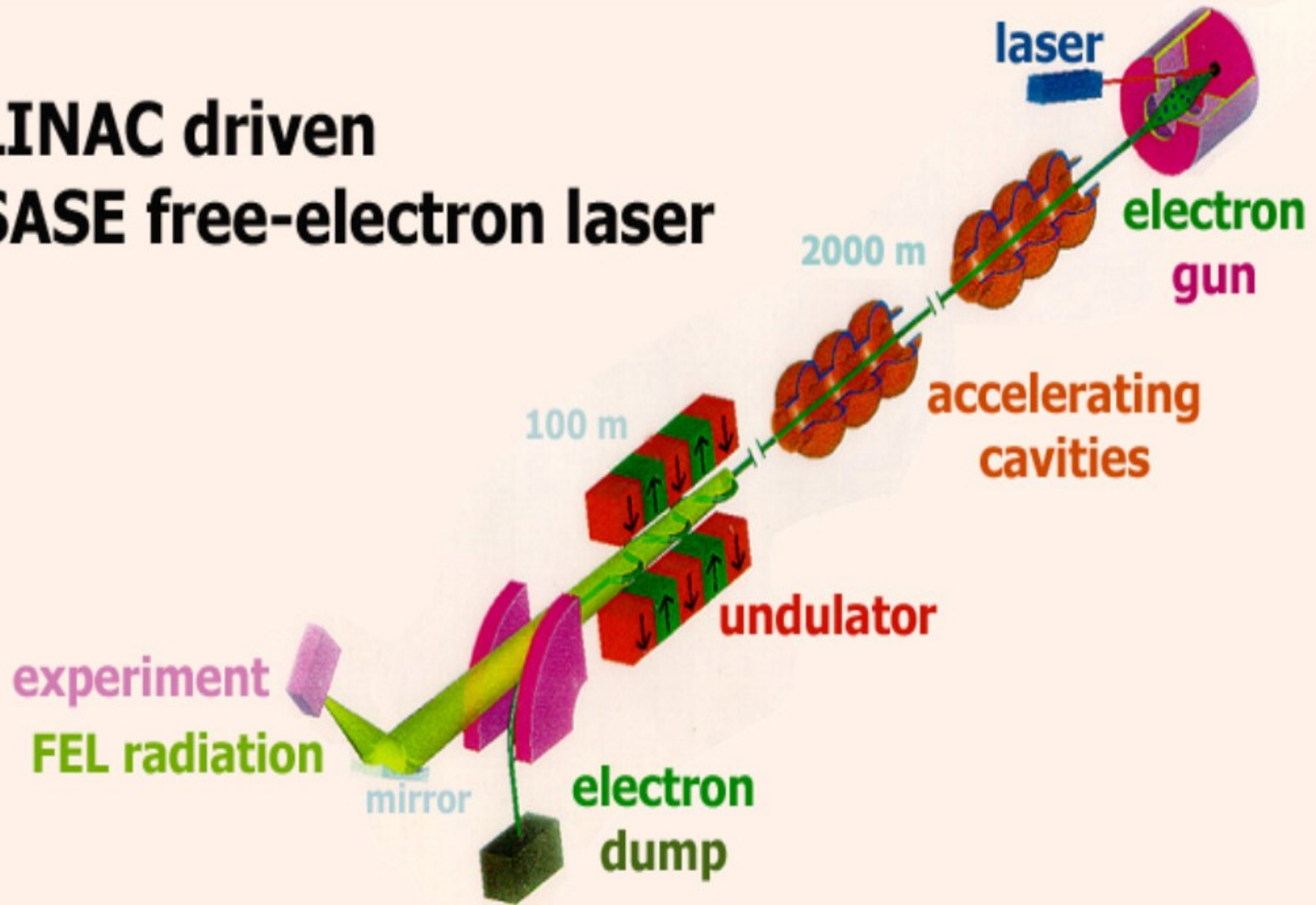
- X-ray FEL radiation (0.2 - 14.4 keV)

- ultrashort pulse duration **100 fs**
- extreme pulse intensities **10^{12} - 10^{14} ph**

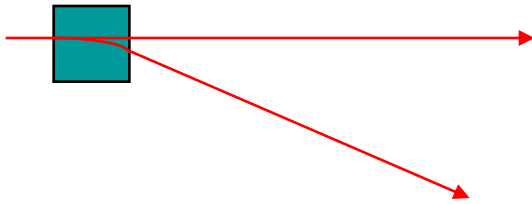
- coherent radiation **$\times 10^9$**
- average brilliance **$\times 10^4$**



LINAC driven SASE free-electron laser

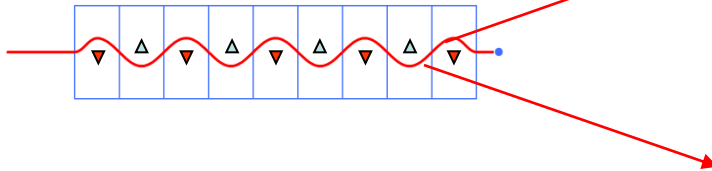


Dipole magnet synchrotron radiation

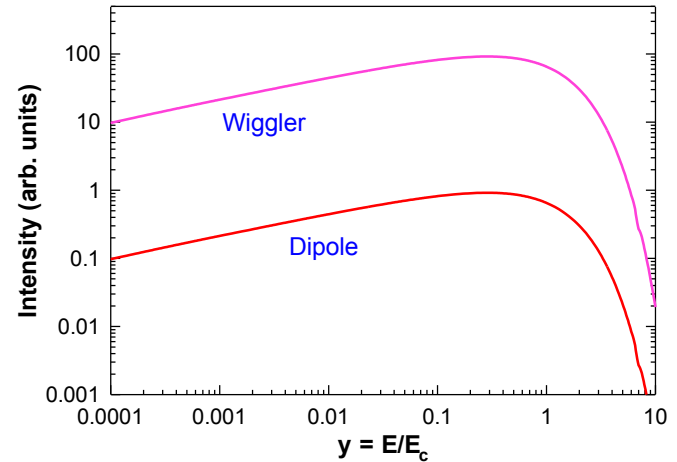


$$dF \sim E^2 I$$

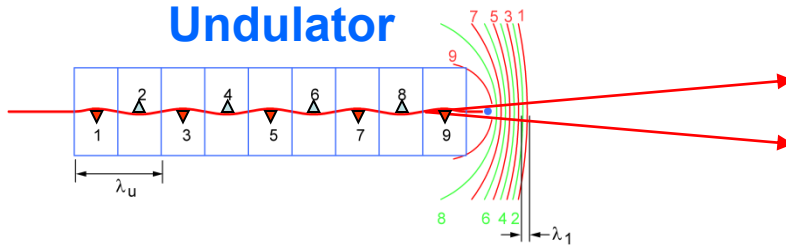
Wiggler



$$\sim 2N E^2 I$$



Undulator

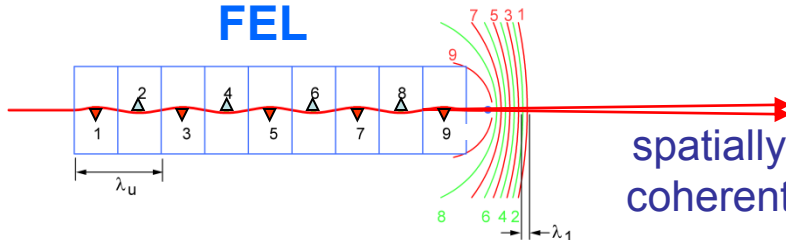


$$\sim N^2 E^2 I$$

$$\sim n_e$$

The radiation emitted by a single electron in subsequent oscillations in an undulator is in phase. Radiation from different electrons is NOT (positional disorder in bunch).

FEL



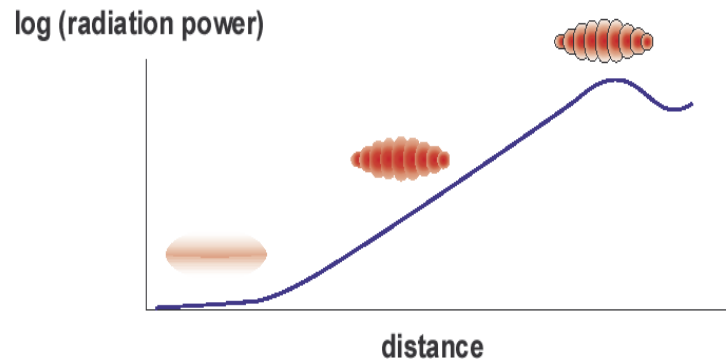
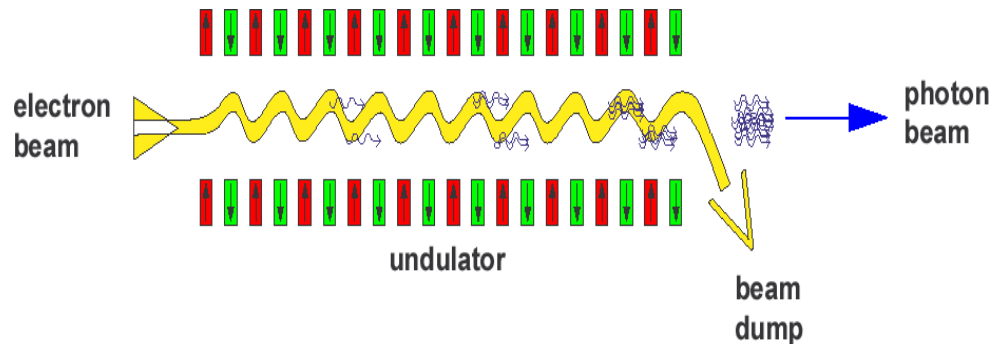
$$\sim n_e^2$$

spatially
coherent

“Phasing” is achieved via positional order in the bunch (micro-bunching) with a period equal to the x-ray wavelength.

SASE (self-amplified spontaneous emission)

Idea: Send a perfect electron beam through a very long undulator using the spontaneous radiation as a seed



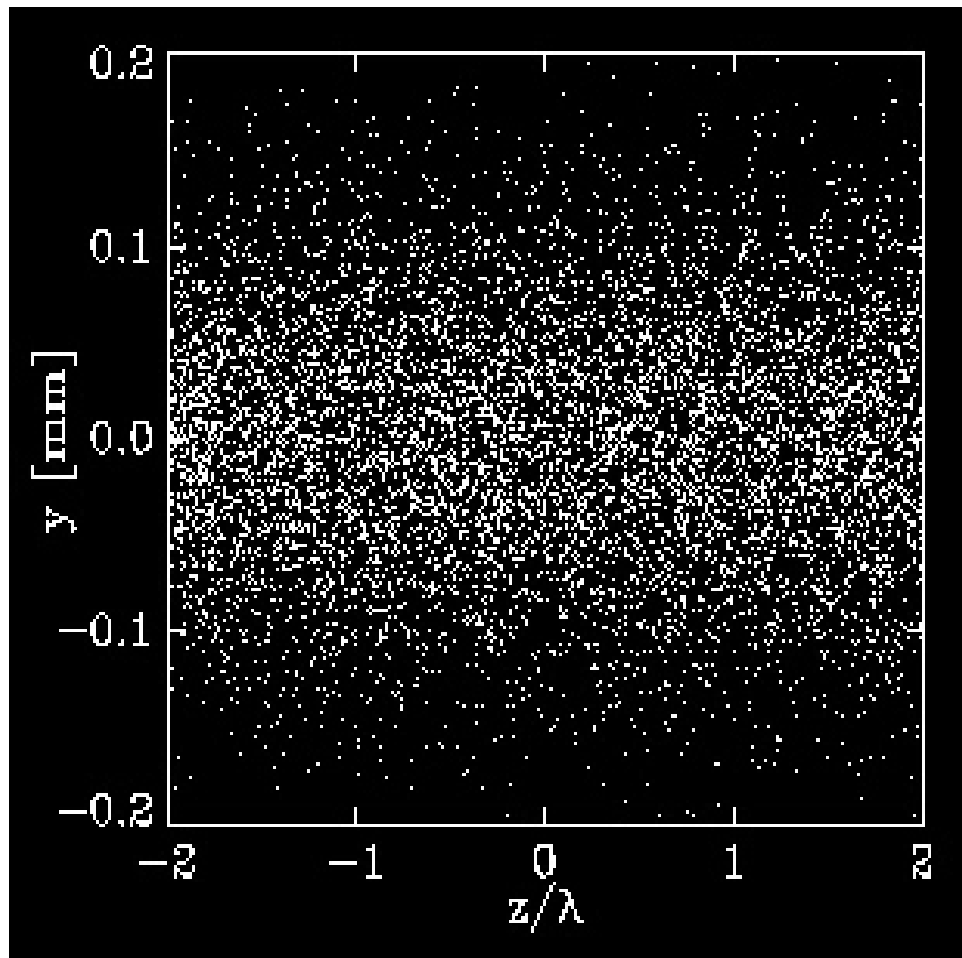
A.M. Kondratenko, E.L. Saldin, Part. Acc. **10**, 207 (1980)

Generation of coherent radiation by a relativistic electron beam in an undulator

R. Bonifacio, C. Pellegrini, L. Narducci, Opt. Commun. **50**, 373 (1984)

Collective instabilities and high-gain regime in a free electron laser

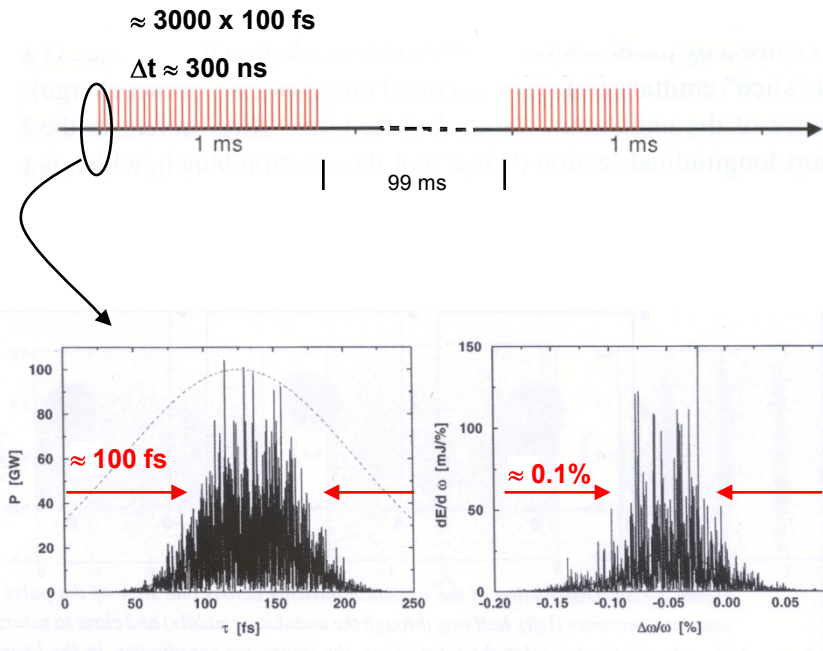
Simulation



GENESIS – simulation for TTF parameters

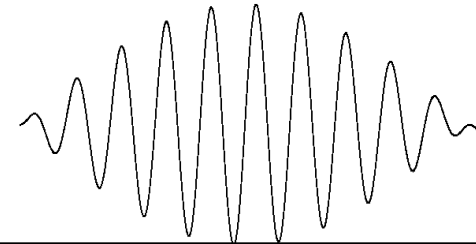
Courtesy Sven Reiche (UCLA)

Time Structure and Coherence Properties of XFEL radiation



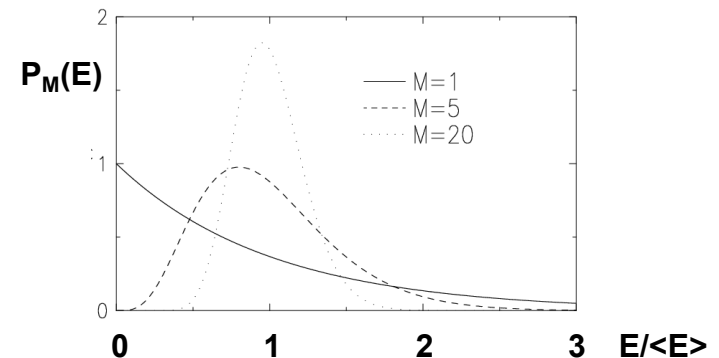
for $\Delta\lambda/\lambda / N(=1000) = 10^{-6}$:

Get single spatially and temporally coherent wave packet carrying about 10^9 photons.



Pulse-to-pulse statistics:

- M=1:** single mode
 strong pulse-to-pulse fluctuations
M>1: multi-mode



$$\begin{aligned}
 \#ph/mode &\approx 10^9 = \#ph/bunch / M(\# \text{ modes}) \\
 &= 10^{12} / 1000 \\
 M &< 1000
 \end{aligned}$$

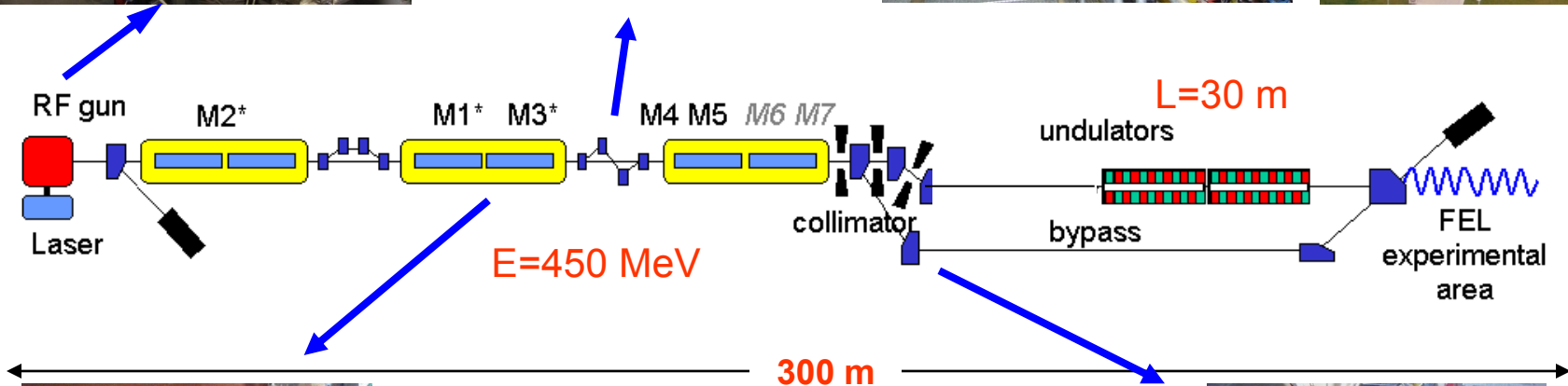
The FLASH facility

$$\lambda \geq 4.5 \text{ nm}$$



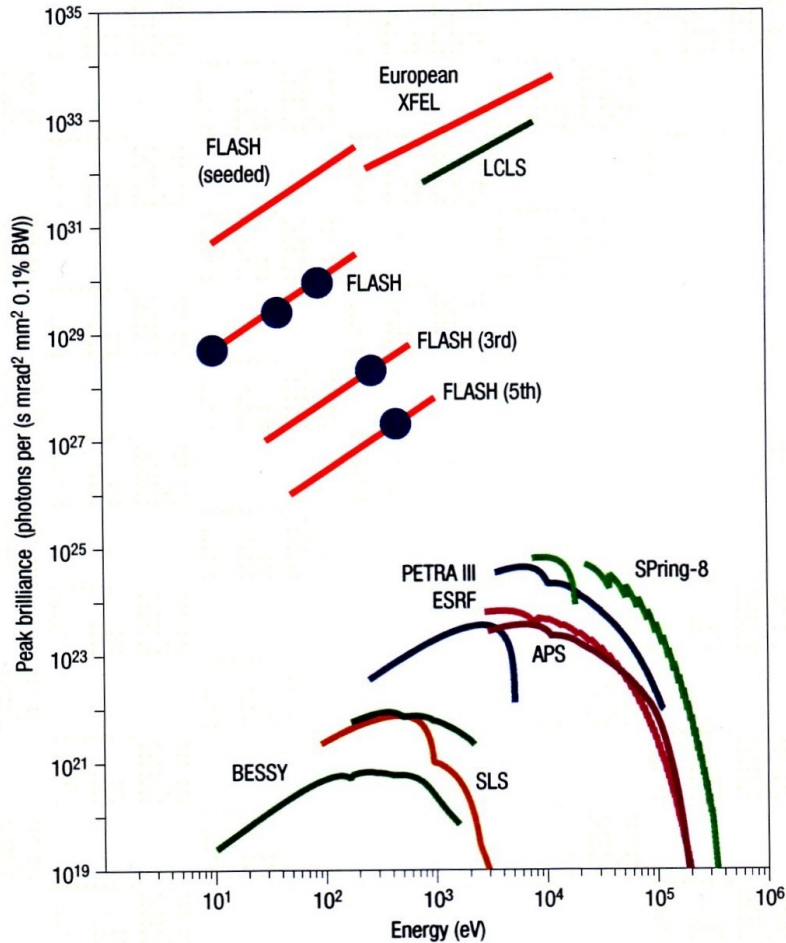
Commissioning: 2004/5
User experiments: 2005

FLASH Overview



- Jan 2005: first lasing at 32 nm
- Aug 2005: first user exper.

FLASH Performance



FLASH performance:

pulse duration ≤ 20 fs

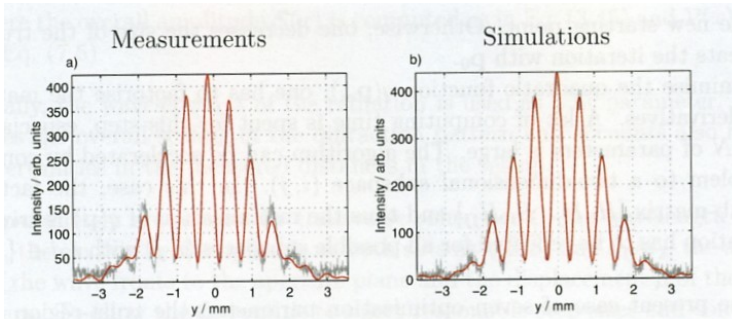
1st 13.7 nm $B_{\text{peak}} = 6 \cdot 10^{29}$

3rd 4.6 nm $B_{\text{peak}} = 2 \cdot 10^{28}$

5th 2.75 nm $B_{\text{peak}} = 2 \cdot 10^{27}$

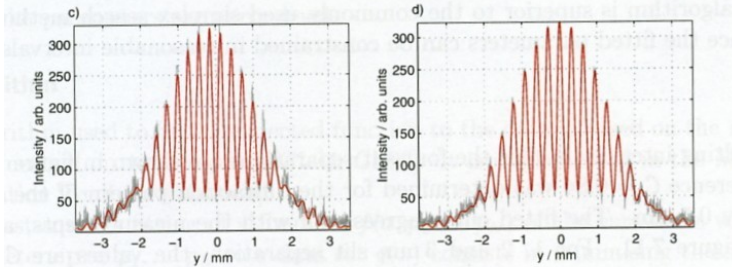
Double Slit Diffraction at FLASH

$d=0.5\text{mm}$

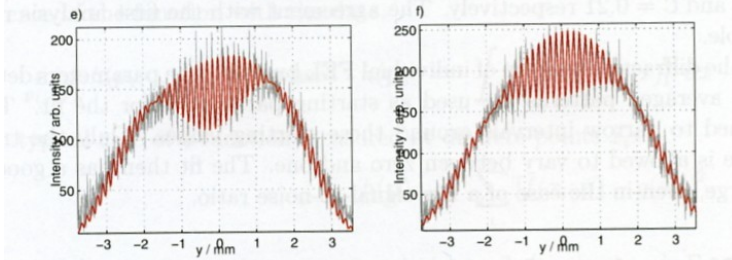


$\lambda = 100\text{ nm}$

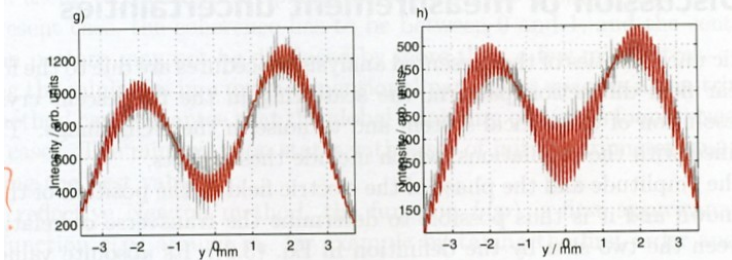
$d=1\text{mm}$



$d=2\text{mm}$



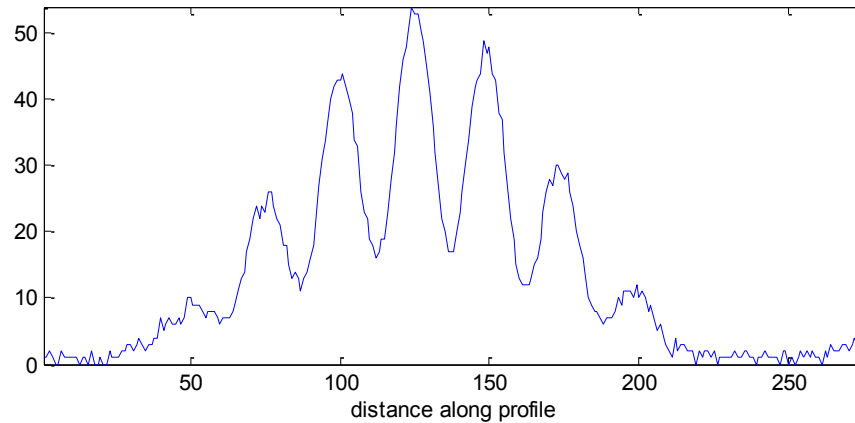
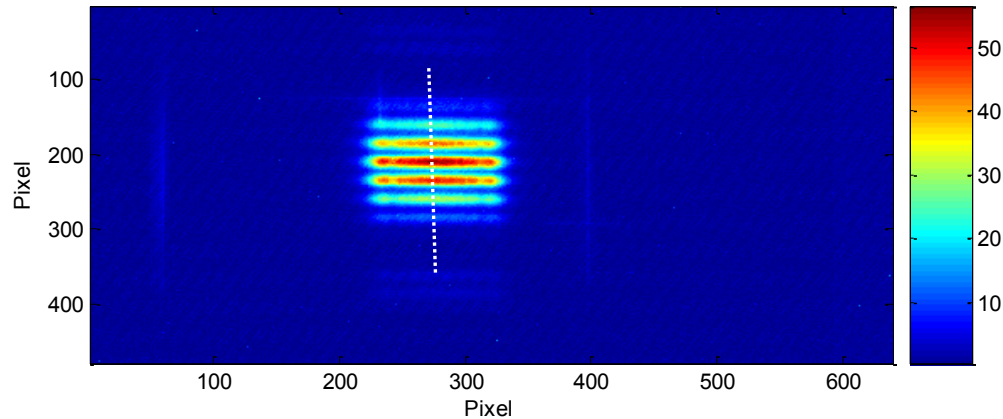
$d=3\text{mm}$



Thesis Rasmus Ischebeck

Double Slit Diffraction at FLASH

: pixel x-axis binning, bunch(es), mm encoder position, aperture, slits_0p15mm_hori_1147339465.tif - None,

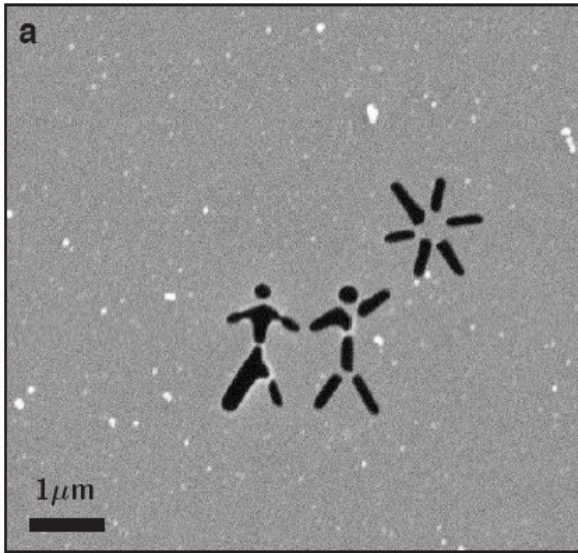


R. Treusch
M. Kuhlmann

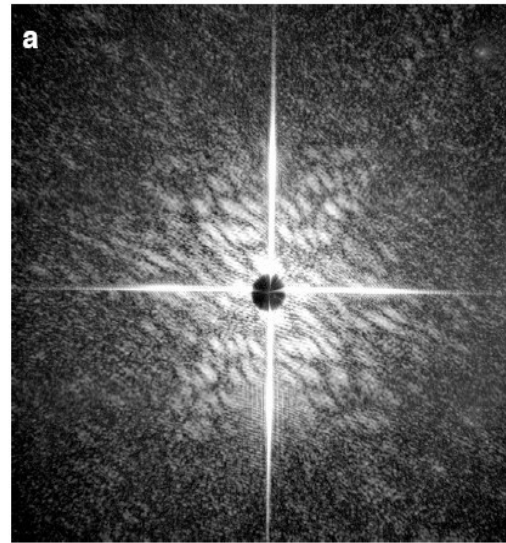
Wavelength $\lambda = 25.6$ nm, Slits separation 0.15 μm , Image is a sum of 10 FEL pulses

■

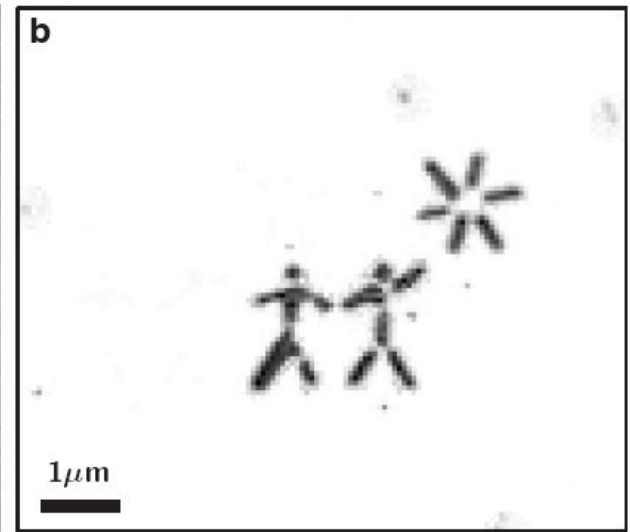
Imaging at a FEL



Model structure in 20 nm SiN membrane

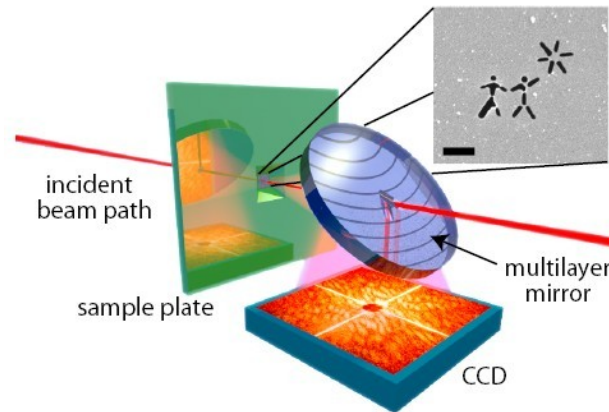


Speckle pattern recorded with a single (25 fs) pulse

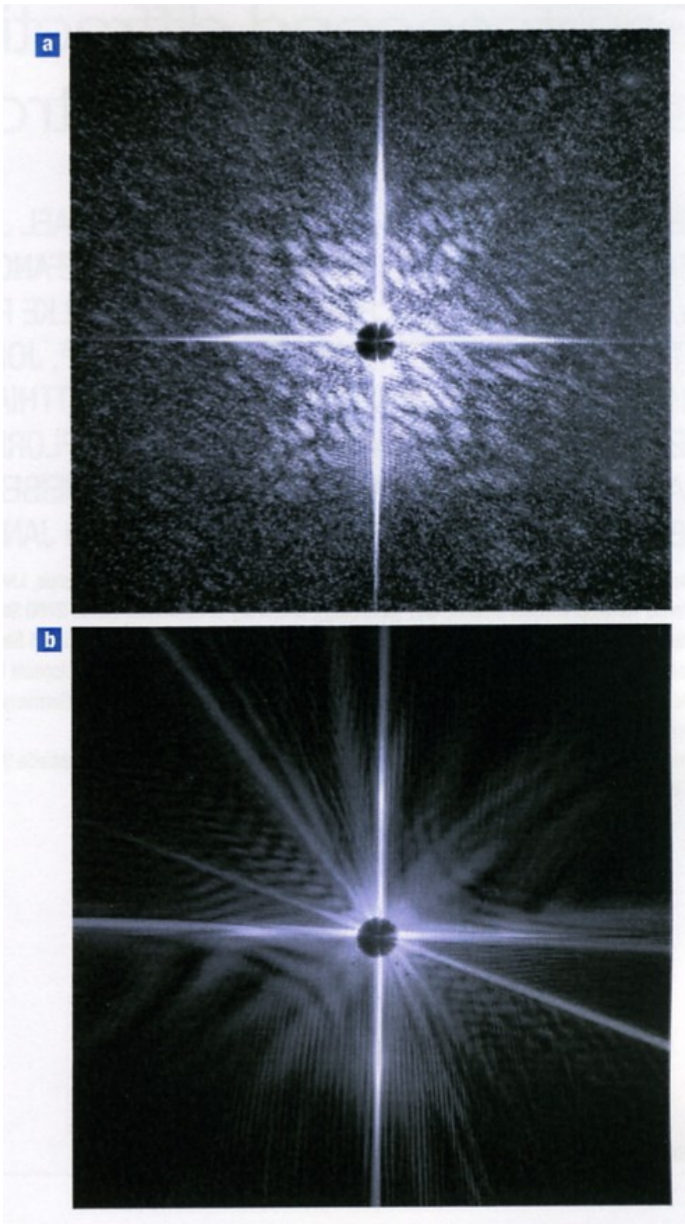


Reconstructed image

Incident FEL pulse:
25 fs, 32 nm,
 $4 \times 10^{14} \text{ W cm}^{-2}$ (10^{12}
ph/pulse)



H. Chapman et al.,
 Nature Physics,
 2,839 (2006)



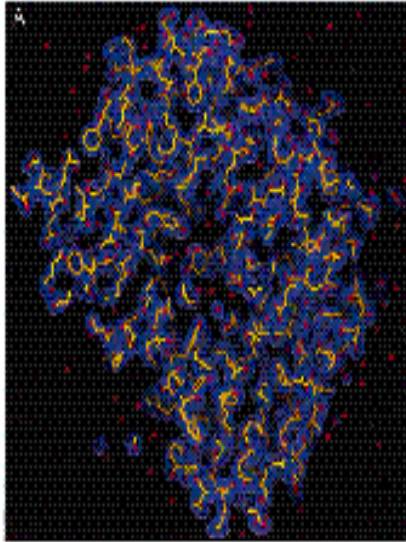
model structure

**First FEL pulse:
25 fs, 32 nm,
 $4 \times 10^{14} \text{ W cm}^{-2}$
(10^{12} ph/pulse)**

“destroyed” model structure

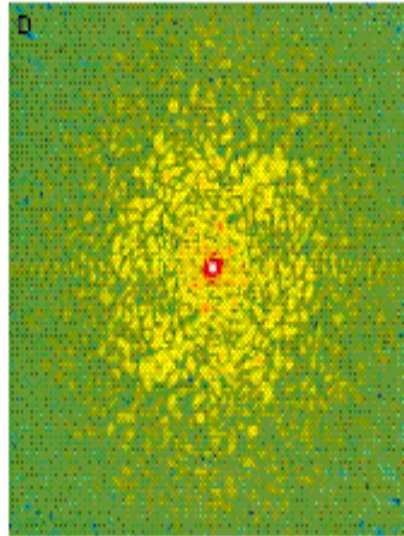
**Second FEL pulse:
25 fs, 32 nm,
 $4 \times 10^{14} \text{ W cm}^{-2}$
(10^{12} ph/pulse)**

An approach to three-dimensional structures of biomolecules by using single-molecule diffraction images: **A simulation**

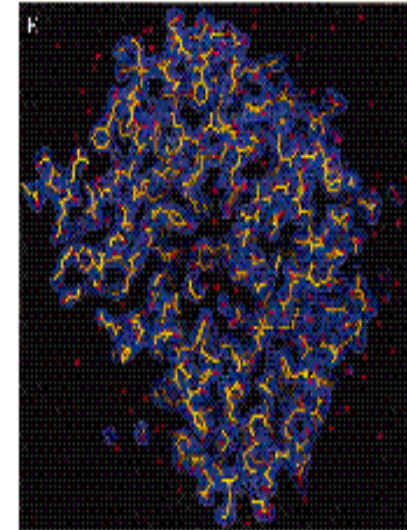


3-D structure (2.5 Å resolution) of rubisco molecule.

(106 kDa)



Top view of a section ($kz=0$) of 3-D scattering pattern from 10^6 single molecules (of known relative orientation) each “exposed” by a single 10 fs XFEL pulse ($\lambda=1.5\text{\AA}$, $0.1\mu\text{m}$ beamsize) containing $2\cdot 10^{12}$ photons.



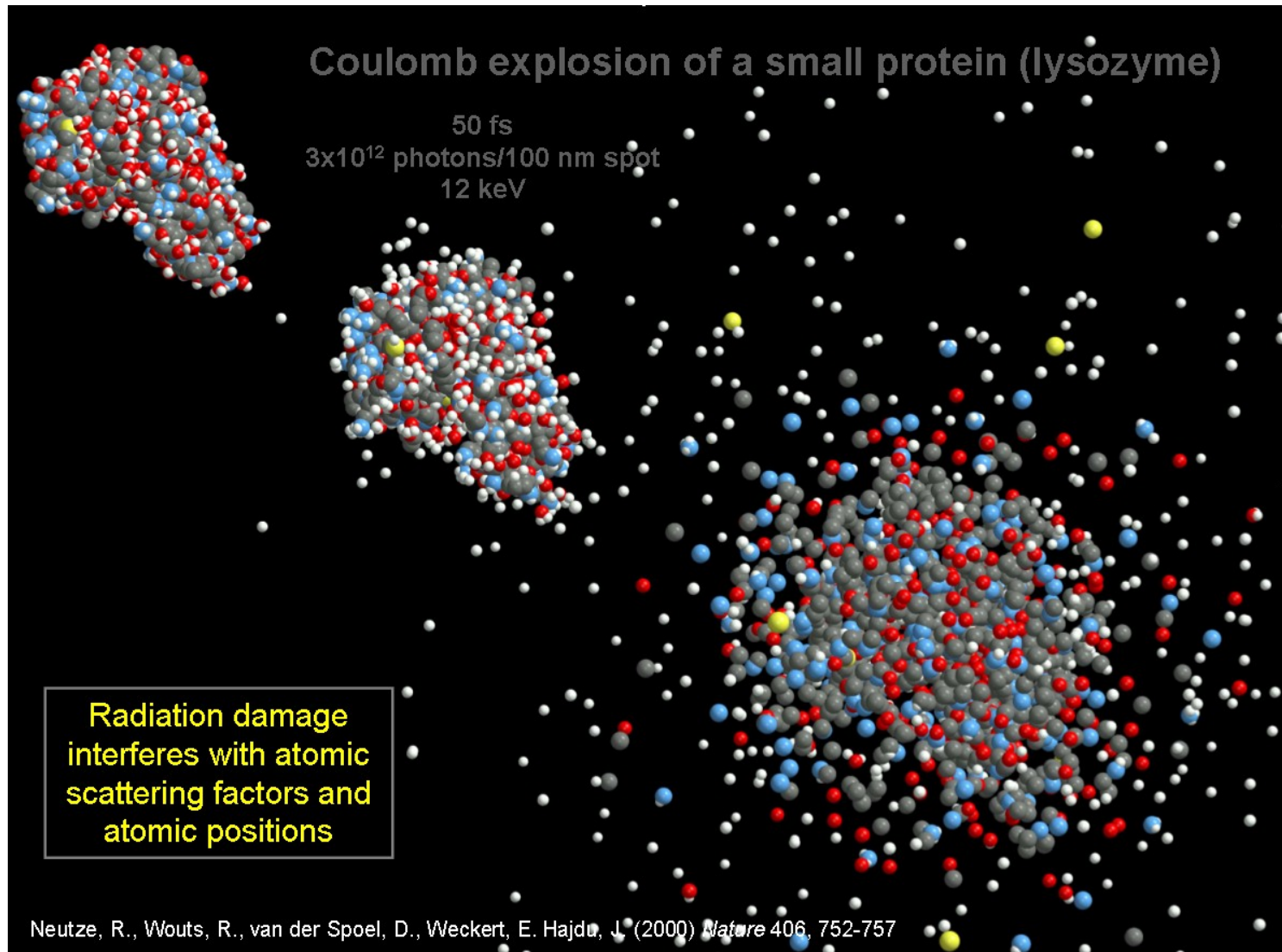
Reconstructed 3-D pattern (from 250 2-D projections). Phasing by “oversampling” technique.

■

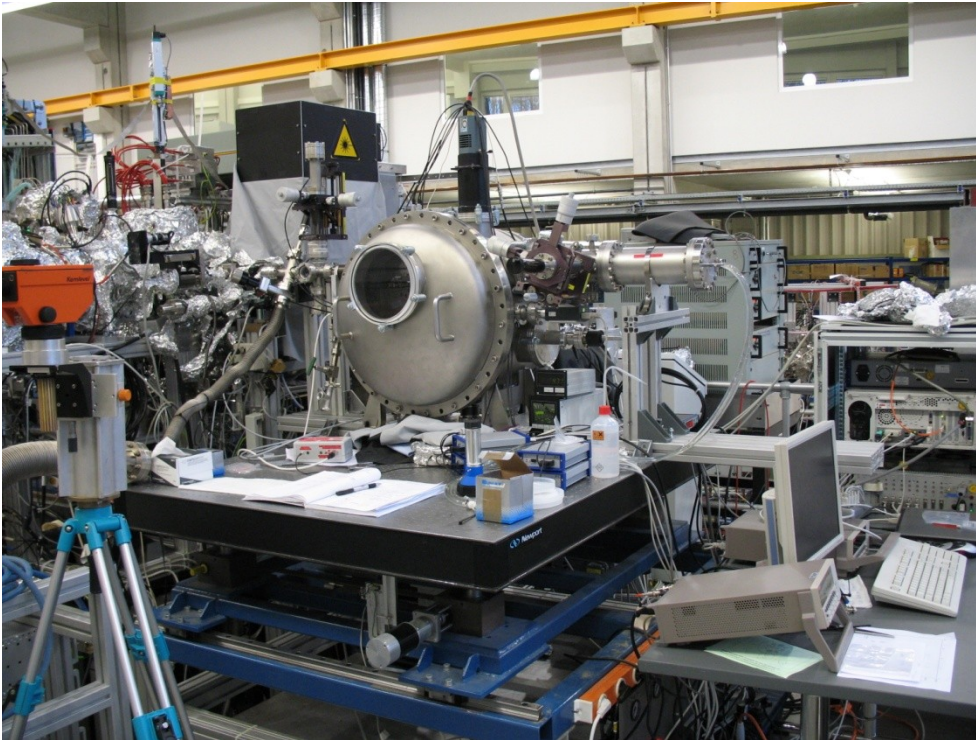
J. Miao, K.O. Hodgson and D. Sayre, PNAS, 98, 6641 (2001)

NOTE: Radiation Damage

Beam – Sample Interaction



■ Magnetic Small Angle Scattering at FLASH (1)



Magnetic Small Angle Scattering:

50 [Co(4Å) / Pt(7Å)] sputtered on 20nm Pt layer on Si₃N₄ membrane, capped with 2nm Pt

FLASH operation 19.12.2007:

SASE on 5th harmonic of 7.97 nm = 1.59 nm \equiv 778 eV

L.-M. Stadler¹, S. Streit-Nierobisch¹, C.Gutt¹, A. Mancuso¹, A. Schropp¹, J. Gulden¹, B. Reime¹, I.Vartaniants¹, E. Weckert¹, B. Pfau², C.M. Günther², R. Könnecke², S. Eisebitt², O. Hellwig³, F. Staier⁴, A. Rosenhahn⁴, T. Wilhein⁵, D. Stickler⁶, H. Stillrich⁶, R. Frömter⁶, H.P. Oepen⁶, R. Treusch¹, N. Guerassimova¹, M. Martins⁷, K. Honkavaara¹, B. Faatz¹, S. Schreiber¹, M.V. Yurkov¹, E.A. Schneidmiller¹, A. Brenger¹, and G. Grübel¹

¹ Deutsches Elektronen-Synchrotron DESY, Notkestraße 85, D-22607 Hamburg, Germany

² Berliner Elektronenspeicherring BESSY – Gesellschaft für Synchrotronstrahlung, Albert-Einstein-Straße 15, D-12489 Berlin, Germany

³ Hitachi Global Storage Technology, 650 Harry Road, San Jose, CA 95120, USA

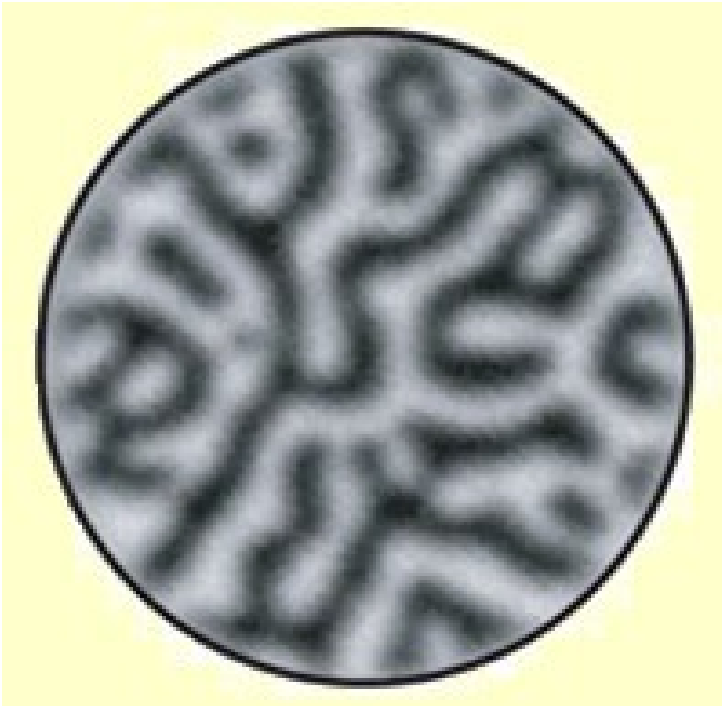
⁴ Institut für Physikalische Chemie, Universität Heidelberg, Im Neuenheimer Feld 229, D-69120 Heidelberg, Germany

⁵ Institute for Xray-Optics, RheinAhrCampus Remagen, FH Koblenz, Südallee 2, 53424 Remagen, Germany

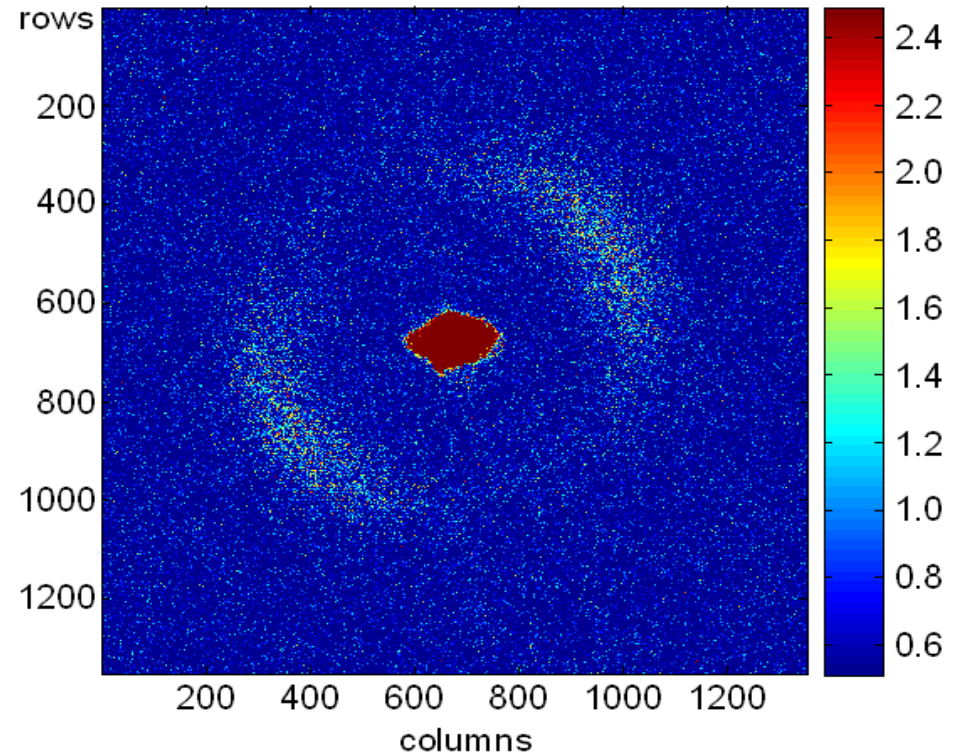
⁶ Institut für Angewandte Physik, Universität Hamburg, Jungiusstraße 11, D-20355 Hamburg, Germany

⁷ Institut für Experimentalphysik, Universität Hamburg, Luruper Chaussee 149, D-22761 Hamburg, Germany

▪ Magnetic Small Angle Scattering at FLASH (2)



meandering magnetic stripe domain of a CoPt multilayer



Magnetic SAXS pattern:

$E = 778.1 \text{ eV}$

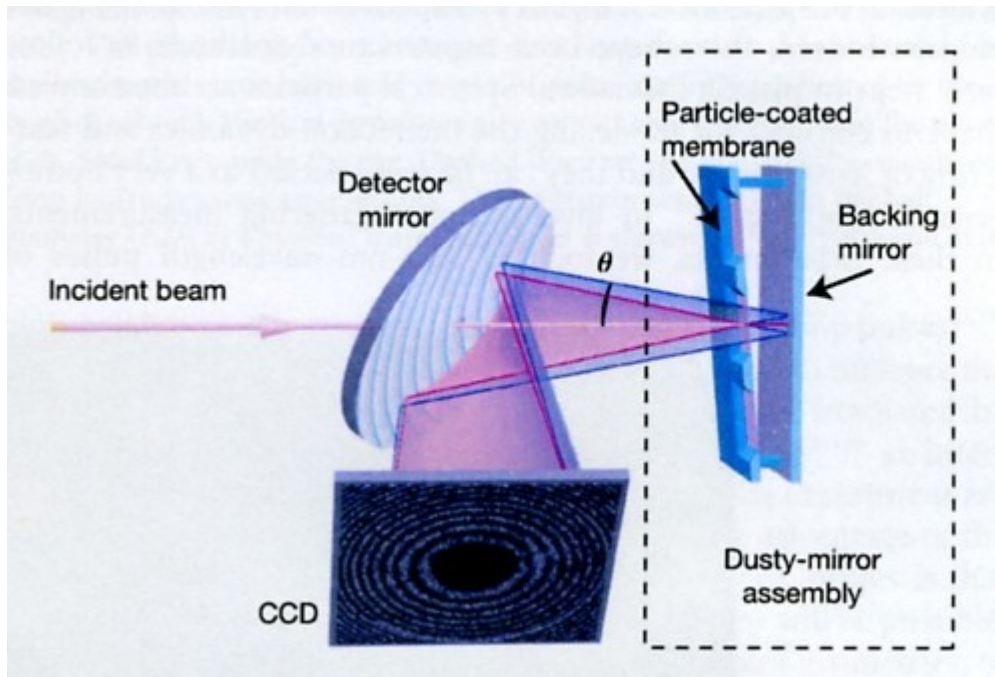
ON (Co L_{III}) resonance

■

Time-dependent Imaging and XPCS at a FEL

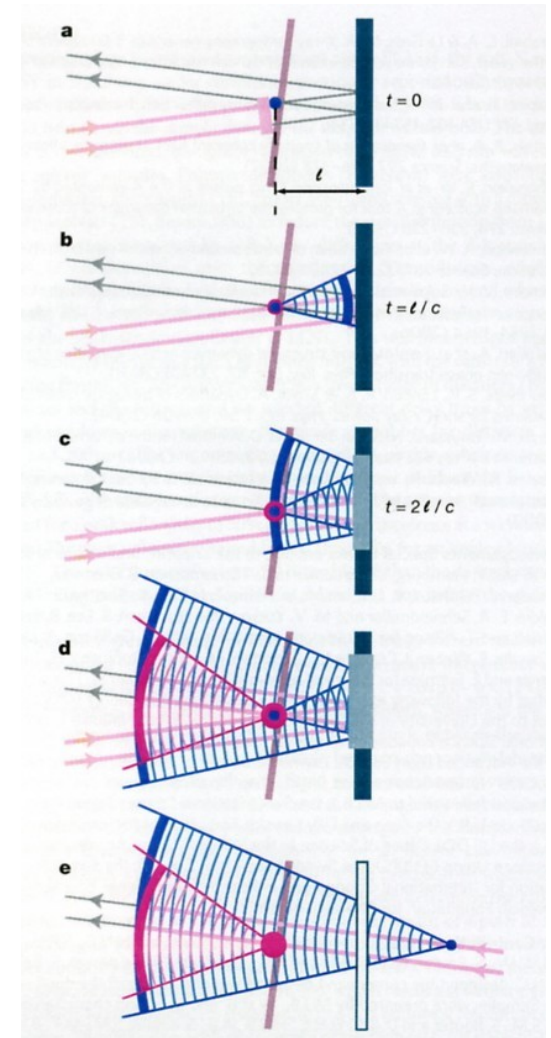
Time delay holography (1)

H.N. Chapman et al.,
Nature 448(2007)676



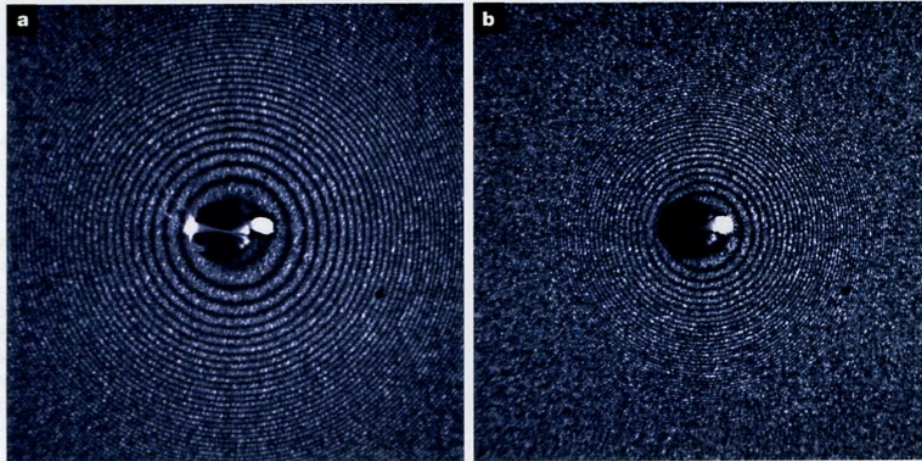
25 fs long single pulses ($\approx 10^{12}$ ph/pulse) of 32.5 nm light from FLASH impinging on a particle coated (140 nm polystyrene PS particles) Si_3N_4 membrane.

The primary diffraction pattern (blue) interferes with the secondary (red) diffraction pattern arising from the (incident) beam being backreflected by a mirror and diffracted after a time delay Δt from the “exploding” PS particles.



Time delay holography (2)

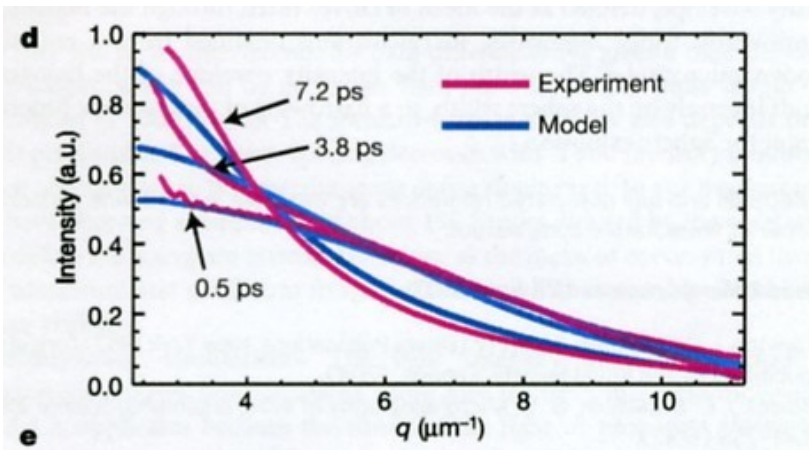
$\lambda = 32.5 \text{ nm}$
 25 fs duration
 $0.5 \cdot 10^{14} \text{ W/cm}^2$
 $\Delta t = 348 \text{ fs}$



$\Delta t = 733 \text{ fs}$

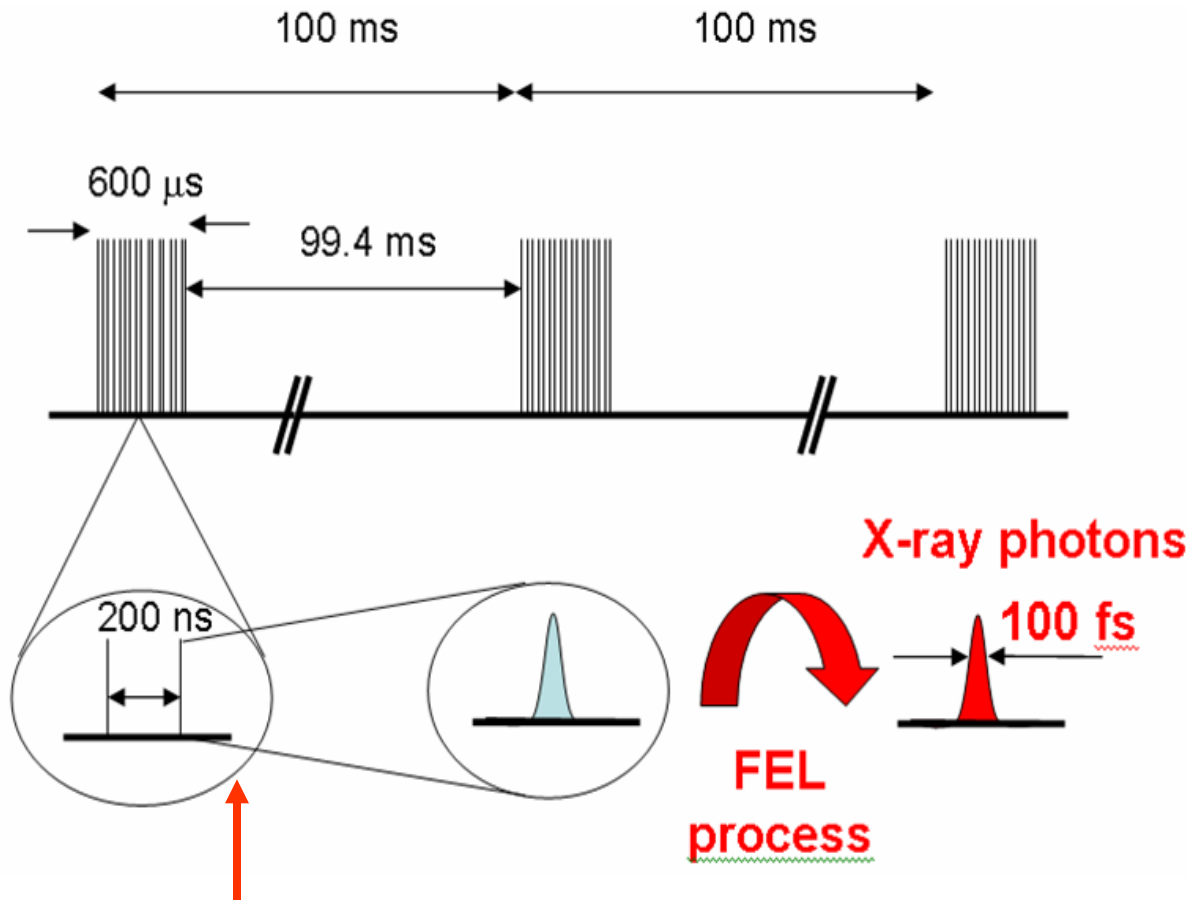
Figure 2 | Time-delay X-ray holograms of 140-nm-diameter polystyrene spheres. The time delays were $348 \pm 1 \text{ fs}$ (a) and $733 \pm 2 \text{ fs}$ (b). The pulses were 32-nm wavelength, 25-fs duration with intensities $(0.5 \pm 0.2) \times 10^{14} \text{ W cm}^{-2}$. The intensities of the holograms are shown on a

linear greyscale, to a half-width of $4.5 \mu\text{m}^{-1}$. We derive the time delays and the change in optical path through the exploding particles from the fringe pattern. The particle sizes are determined from the envelope of the intensity.



The intensity envelope of the interference pattern evolves to lower q indicating increasing radii of the particles.

XPCS at a XFEL source:



$t > 0.1$ s
 $200\text{ ns} < t < 600$
 μs :

” movie” mode

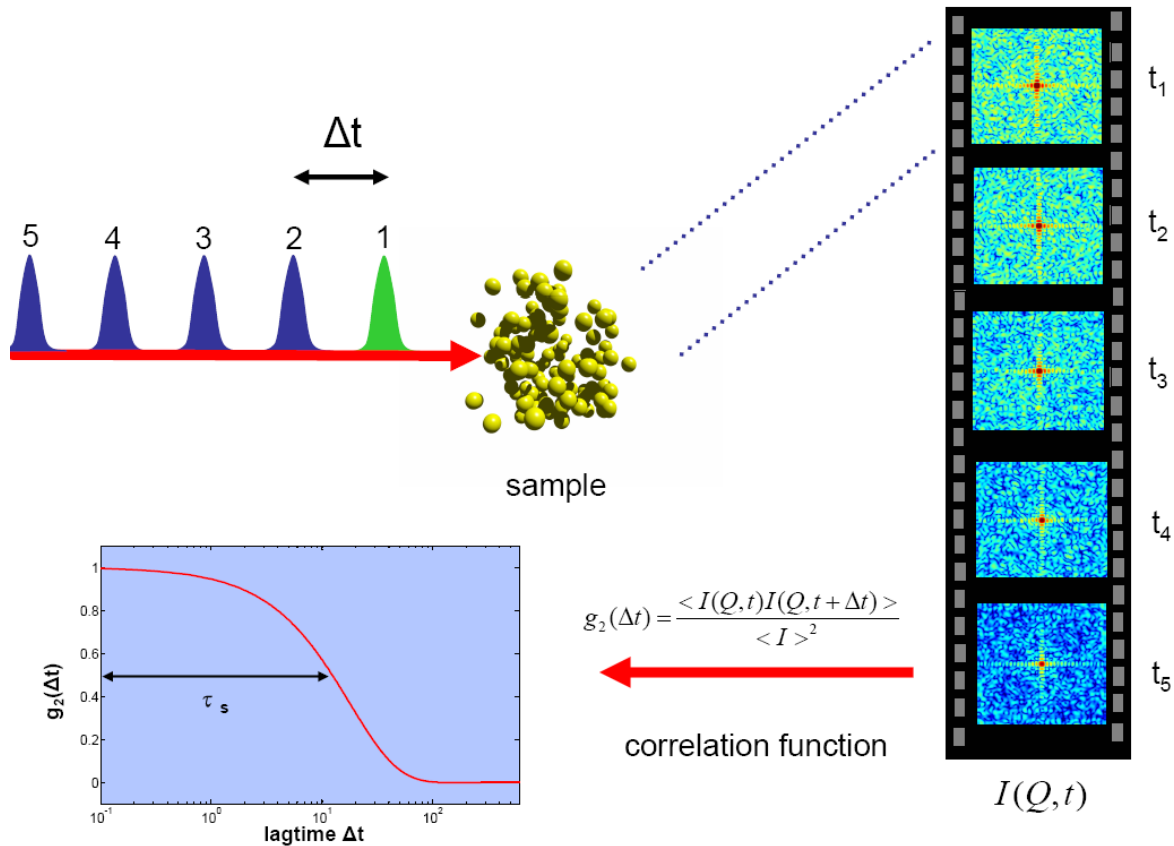
$1\text{ ps} < t < 10$ ns:

for “all” times:

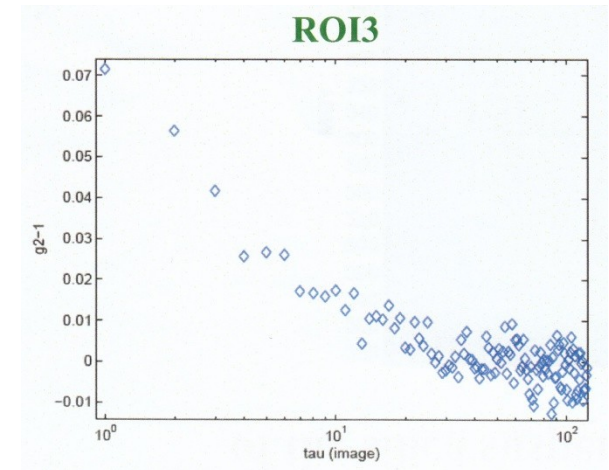
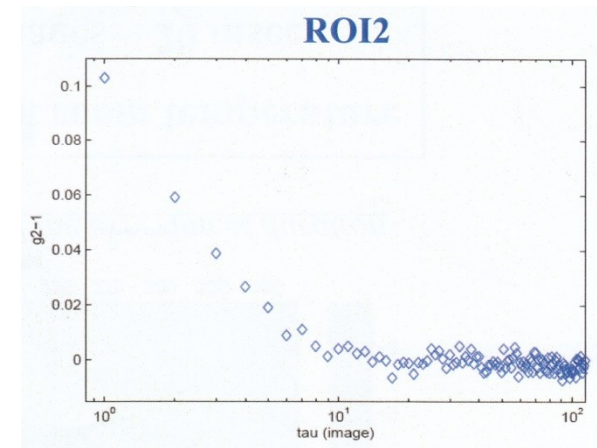
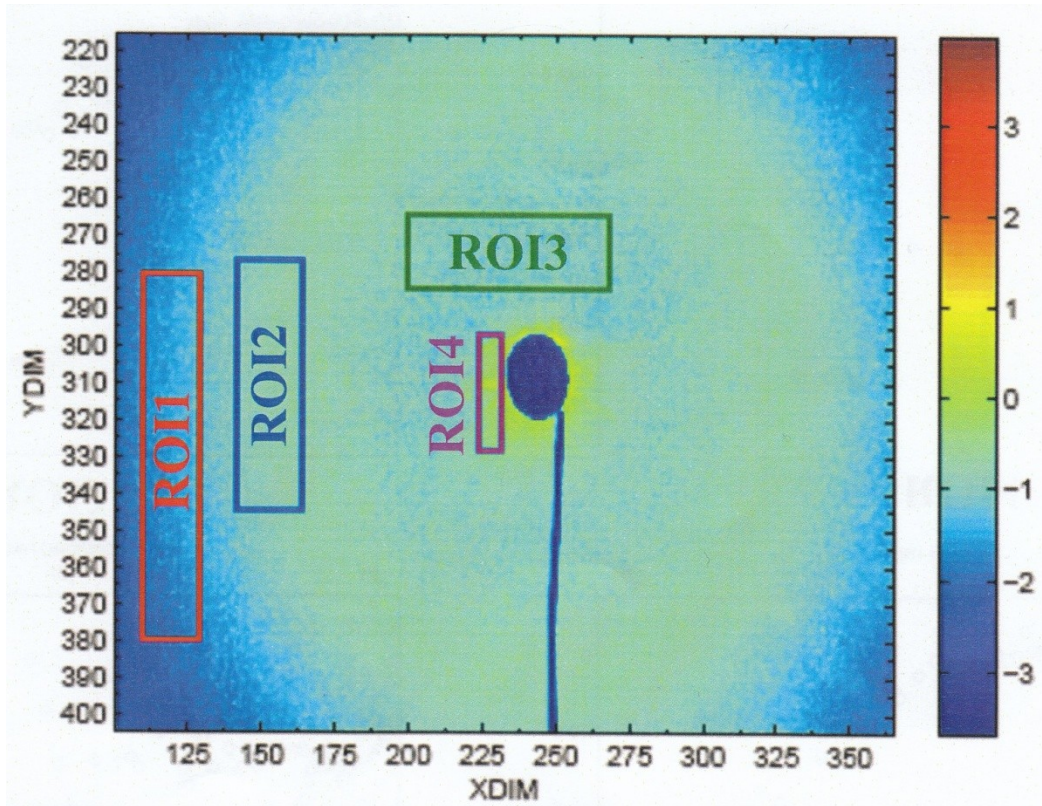
“delay-line” mode

“pump-probe” mode

- XPCS at a FEL source: Movie Mode



The Quest for fast 2-D Detectors



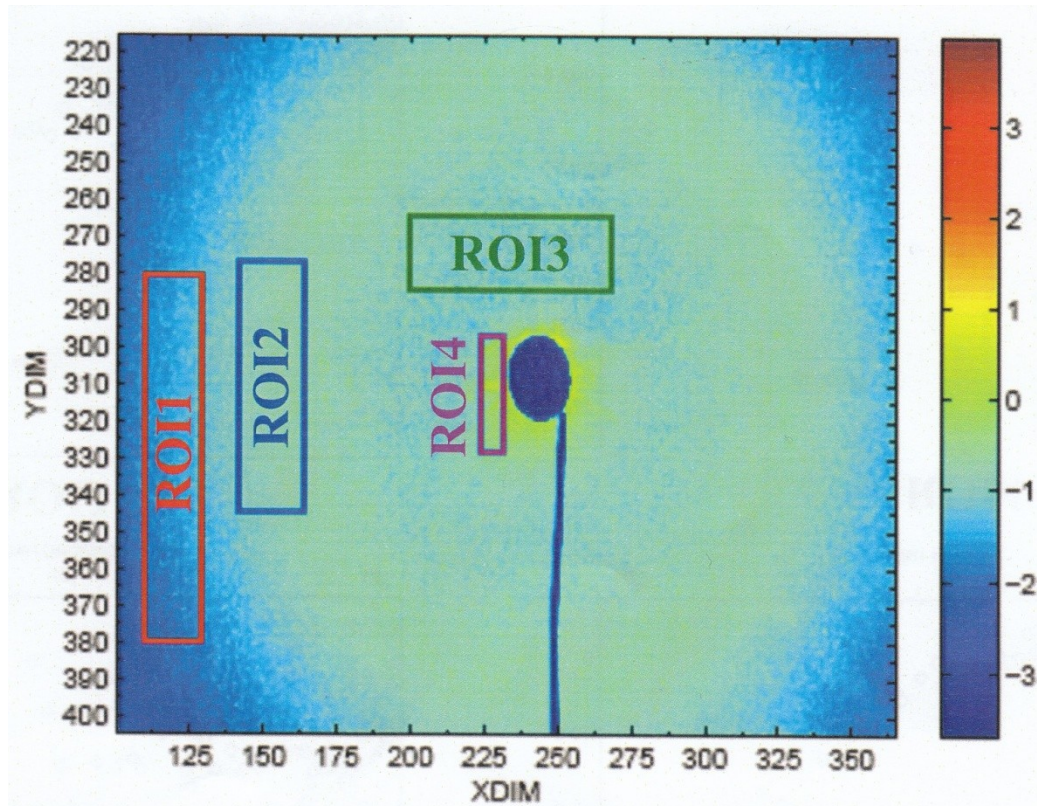
Pilatus detector module: 172 μm pixel size

Silica particles suspended in PPG

sum of 5000 frames with 30 ms exposure

data taken at cSAXS/SLS

▪ The Quest for fast 2-D Detectors



Data acquisition time with 0-D detector: $\approx 2\text{-}3$ days

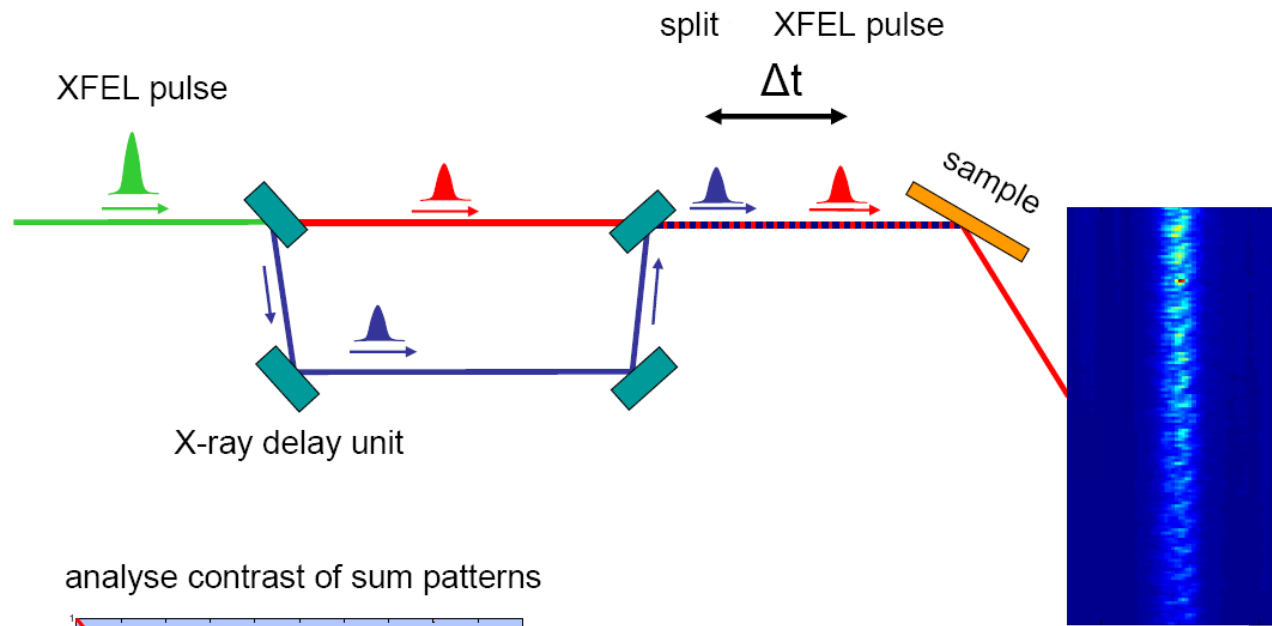
Pilatus: sum of 5000 frames with 30 ms exposure separated by a 20 ms delay btw. 2 frames:

Total data acquisition time:
 $5000 \times 50 \text{ ms} = 250 \text{ s}$

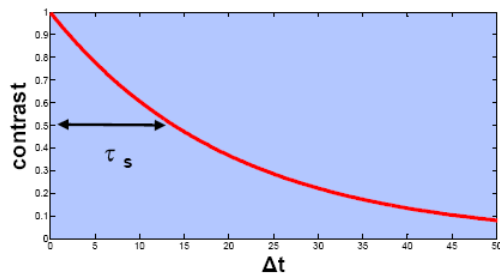
Estimated data acq. at XFEL:

$3000 \times (100 \text{ fs} + 200 \text{ ns}) \approx 600 \mu\text{s}$

Delay Line Mode



analyse contrast of sum patterns



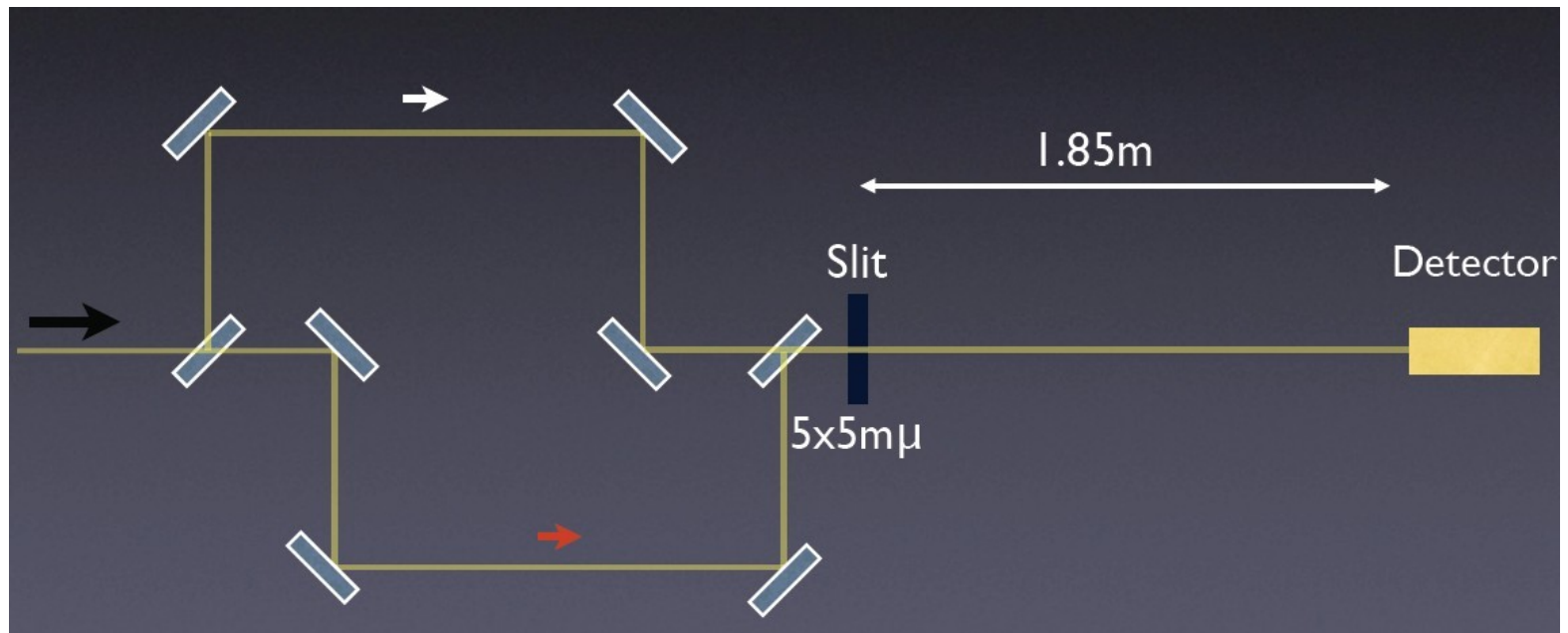
sum of speckle patterns
from prompt and delayed
pulse recorded on CCD

“Delay Line” Mode: $1\text{ ps} < \Delta t < 10\text{ ns}$ ($1\text{ ps} \Leftrightarrow 0.3\text{ mm}$; $1\text{ ns} \Leftrightarrow 300\text{ mm}$) “luminosity limited”.

Delay Line for “hard” X-Rays

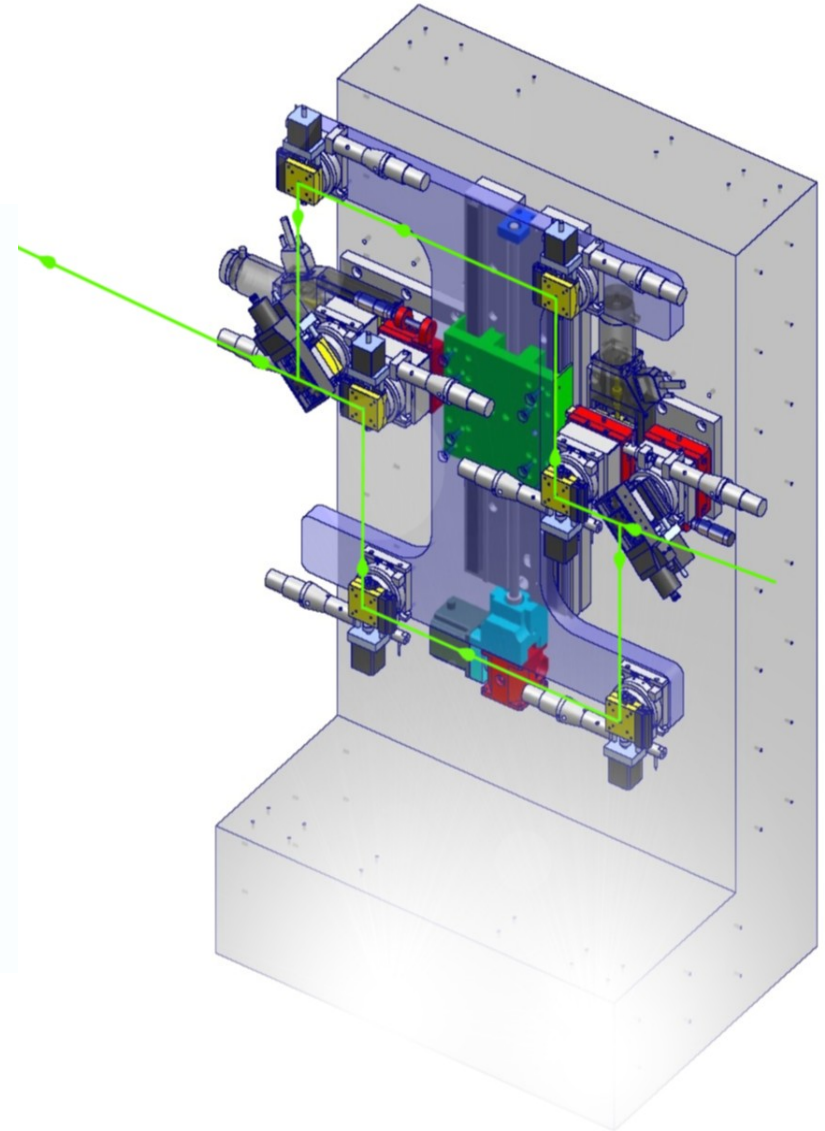
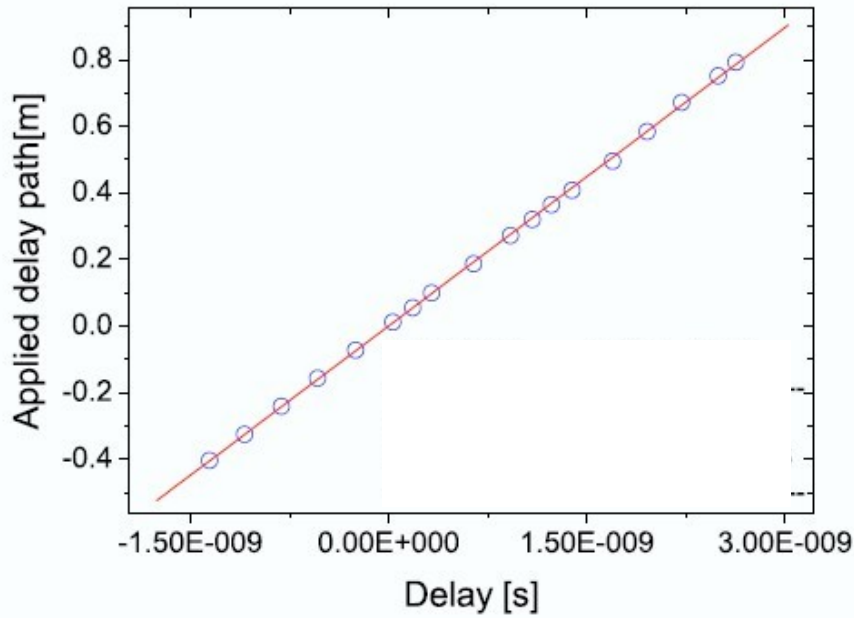
Si(511) at 8.38 keV

$\Delta t_{\text{max}} = 2.8 \text{ ns}$

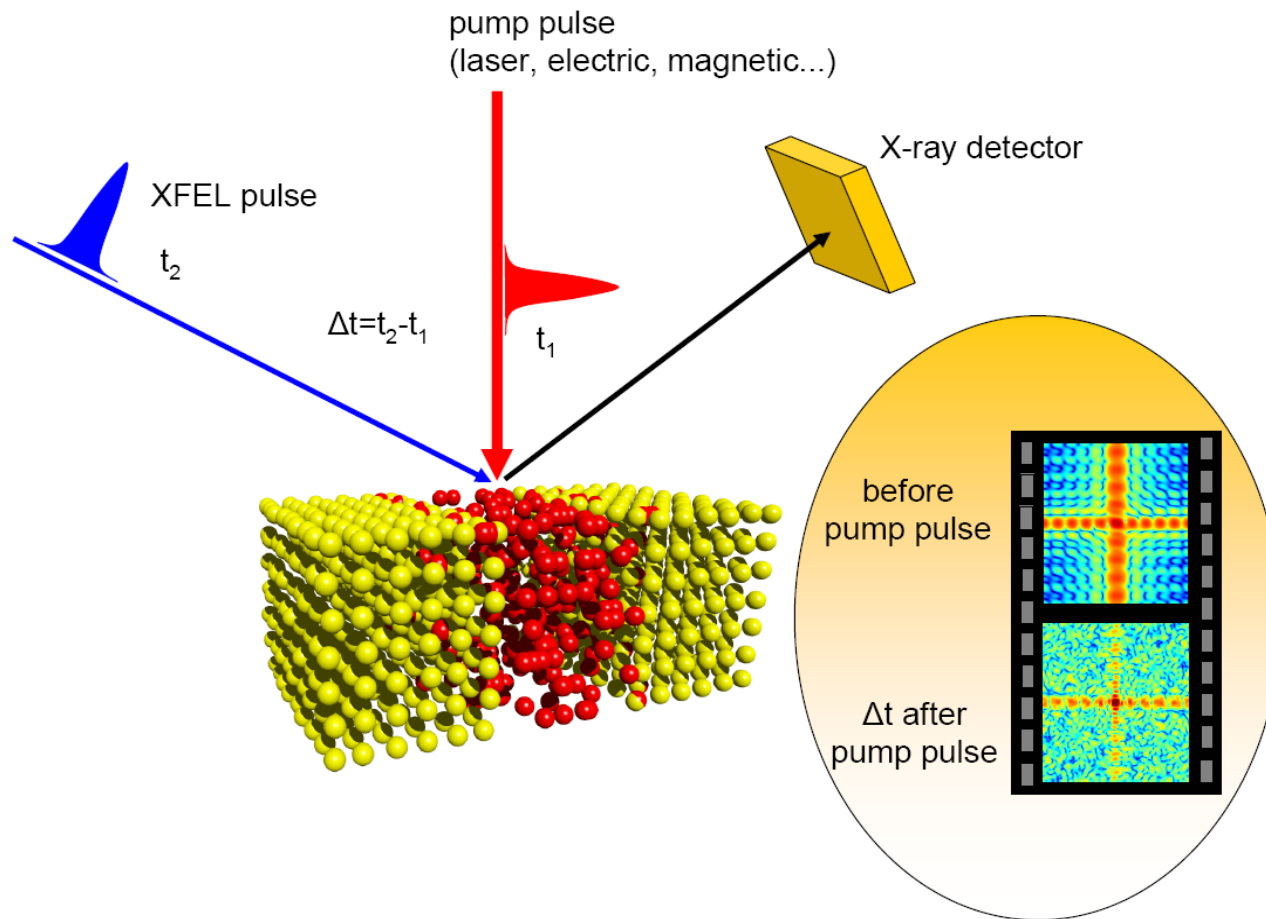


W. Roseker, H. Schulte-Schrepping,
A. Ehnes, H. Franz, O. Leupold and
G. Grübel

- X-Ray delay line



XPCS at a FEL source: pump-probe mode



▪

XPCS at a FEL source

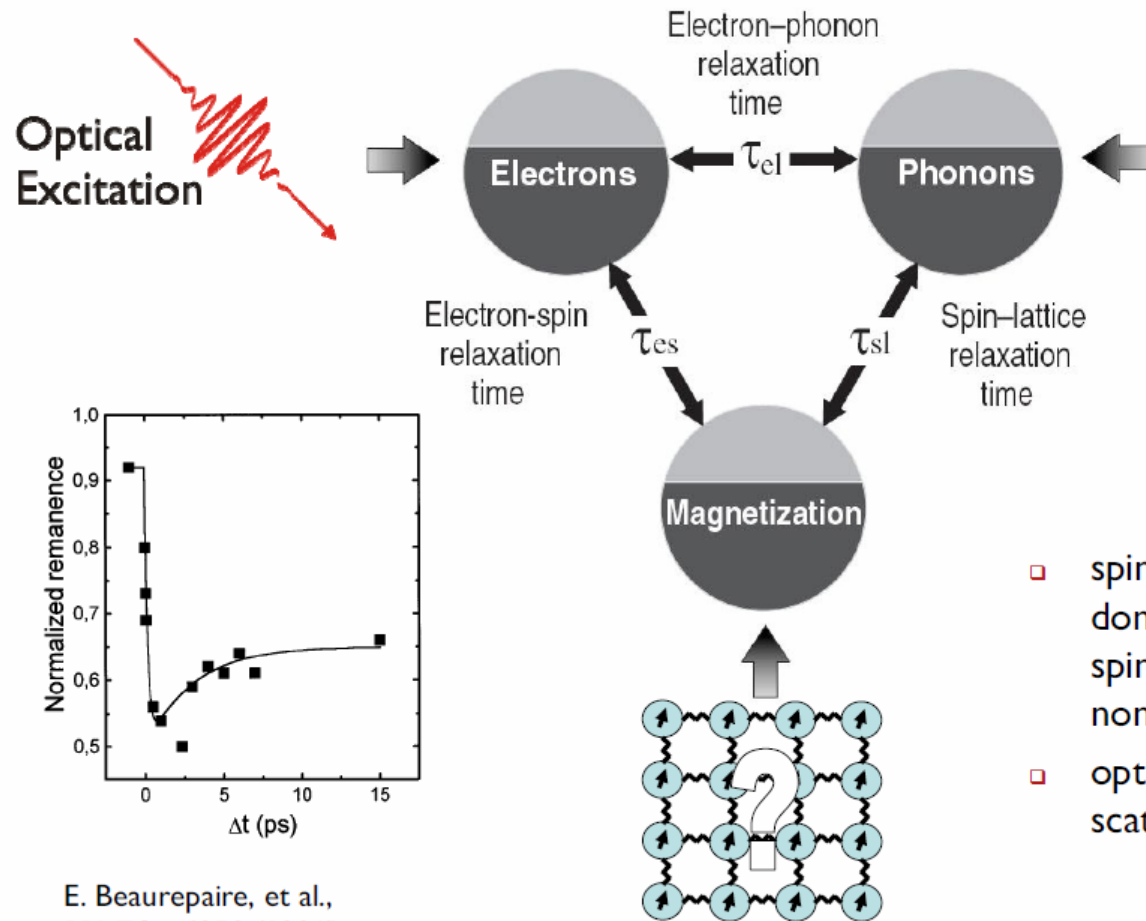
Magnetization Dynamics

Ferroelectrics

Ultrafast dynamics at surfaces and interfaces of liquids

.....

Ultrafast demagnetization



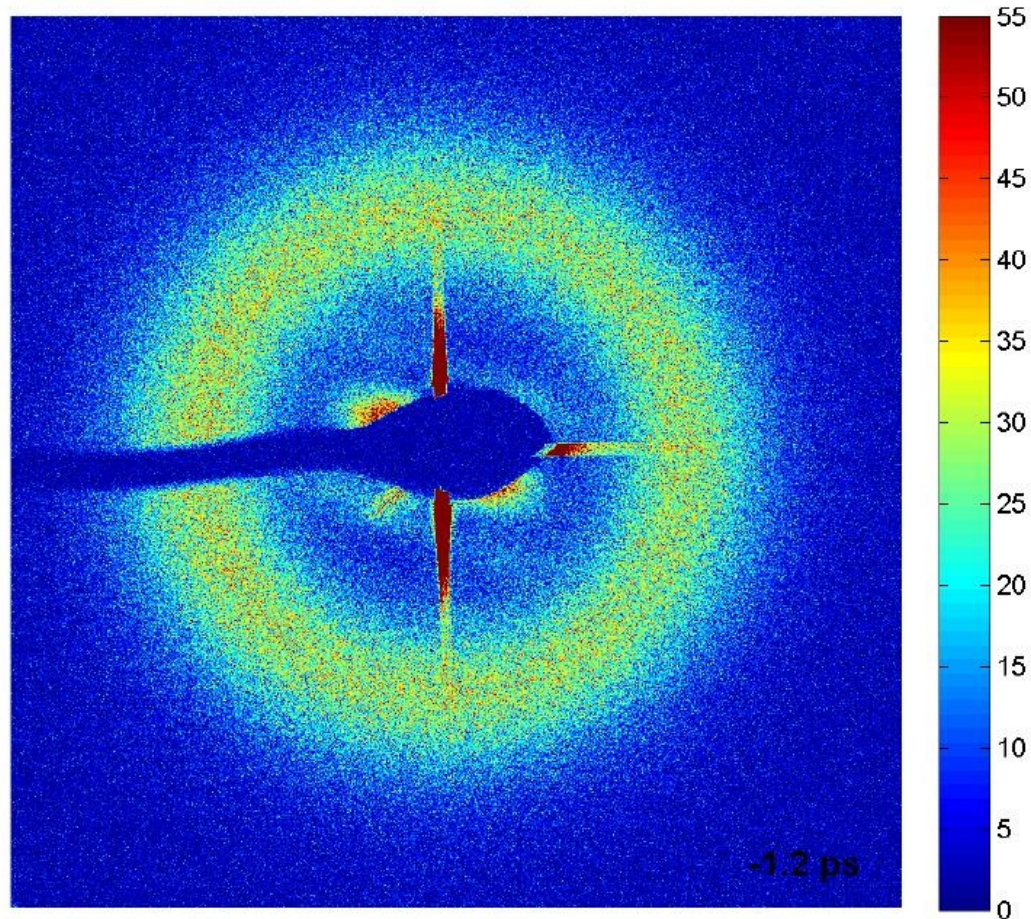
- spin lattice excitations: domains, vortices spin waves & their non-linear interactions
- optical pump – x ray scattering probe

E. Beaurepaire, et al.,
PRL **76**, p4250 (1996)

B. Koopmans et al. *Nature Materials* **9**, 259 (2010)

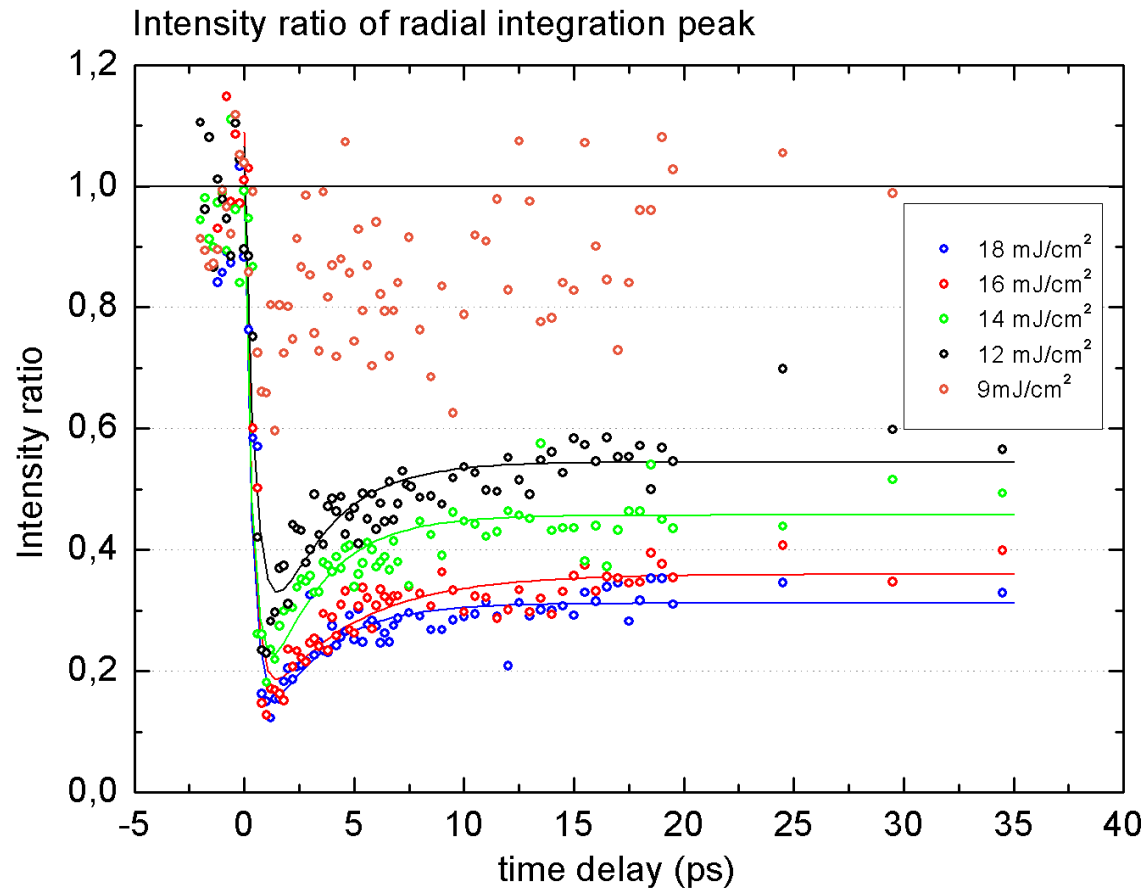
C. Boeglin et al. *Nature* **465**, 458 (2010)

Magnetic SAXS data – time dependant



Time-delay scan for 18 mJ/cm² pumping power, $\Delta t = -1.2$ ps...34.5 ps

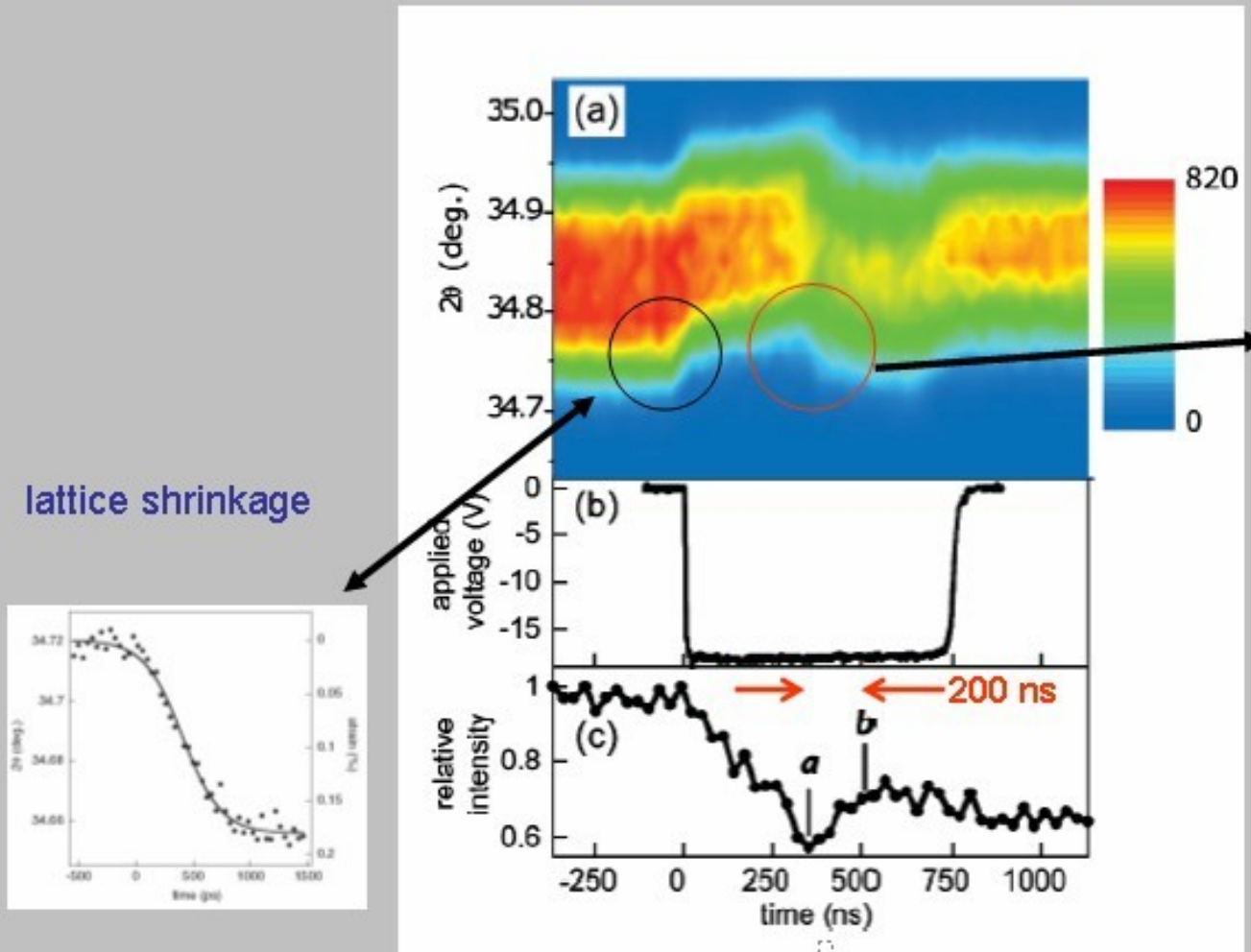
Ultra-fast demagnetization from SAXS data



High pump fluence, characteristic demagnetization time below 1 ps

Polarization Switching in Ferroelectrics

Domain wall motion during polarization switching in ferroelectric $\text{Pb}(\text{Zr},\text{Ti})\text{O}_3$



PZT (002) reflection
(0.1 μm , 10 keV, APS)

polarization switching:

Nucleation of domains with reversed polarity, followed by domain growth with domain wall velocities of here about 40m/s.

Study fluctuations and domain growth via XPCS.

A. Grigoriev, D.-H. Do, D. M. Kim, C.-B. Eom, B. Adams, E.M. Dufresne and P. G. Evans, Phys. Rev. Lett. 96 (2006)187601

XPCS at a XFEL source: An example

Ultrafast dynamics at surfaces and interfaces in:

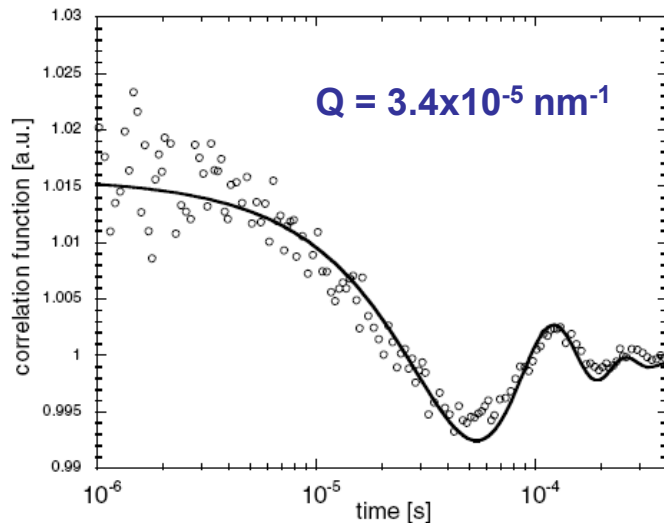
- o liquids,
- o membranes,
- o ...

XFEL:

Onset of non-classical behaviour ($Q > 1 \text{ nm}^{-1}$, beyond continuum hydrodynamics)

Study capillary wave dynamics at high Q in model systems ($\lambda=1\text{\AA}$, $Q=1 \text{ nm}^{-1}$)

Today: $Q_{\text{max}} \approx 10^{-4} \text{ nm}^{-1}$ (water)



τ [s] countrate (FEL)

Water	$\approx 25 \text{ ps}$	20
Mercury	$\approx 0.5 \text{ ps}$	0.3 *
Glycerol	$\approx 1 \text{ }\mu\text{s}$	3000

* expect important deviations due to layering at the surface

G. Grübel et al., TDR XFEL, DESY (2006)

G. Grübel et al., NIM B, submitted

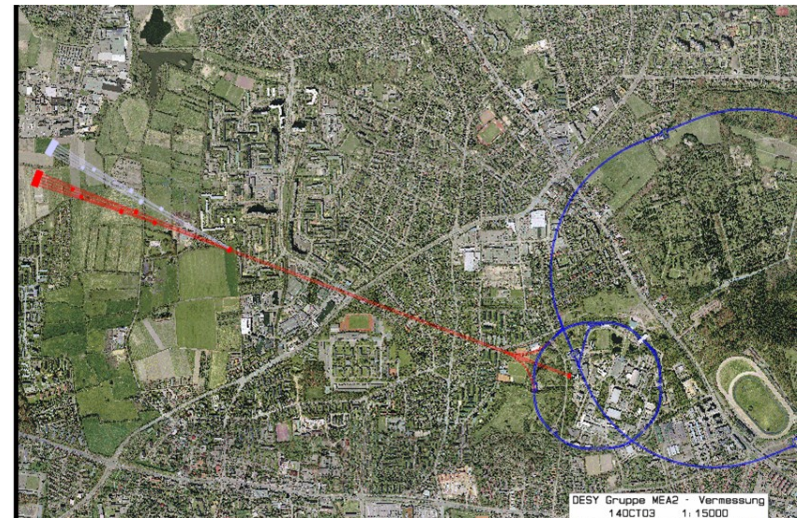
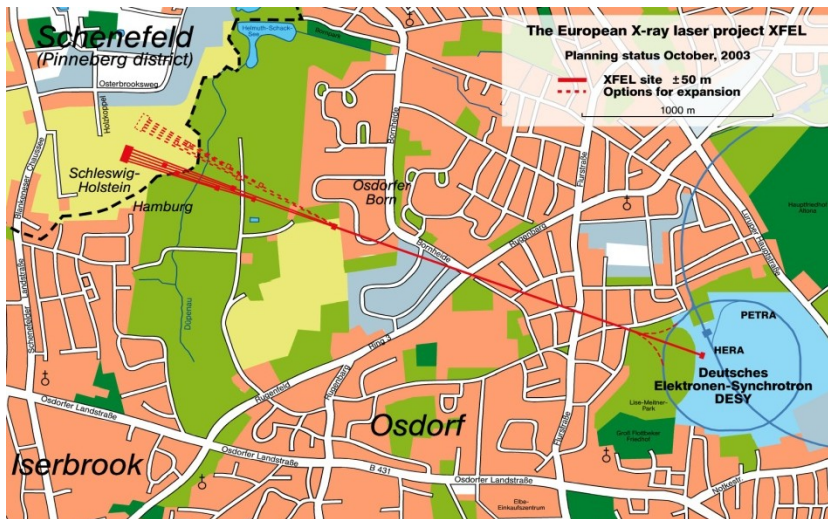
C. Gutt et al., PRL 91 (2003)76104

The XFEL

www.xfel.eu



← 3.4km →



- # Acknowledgements

GRK 1355

L. Stadler, C. Gutt, A. Robert