

# Methoden moderner Röntgenphysik I

## Coherence based techniques III

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# Outline

**18.12. 2008**

**Introduction to Coherence**

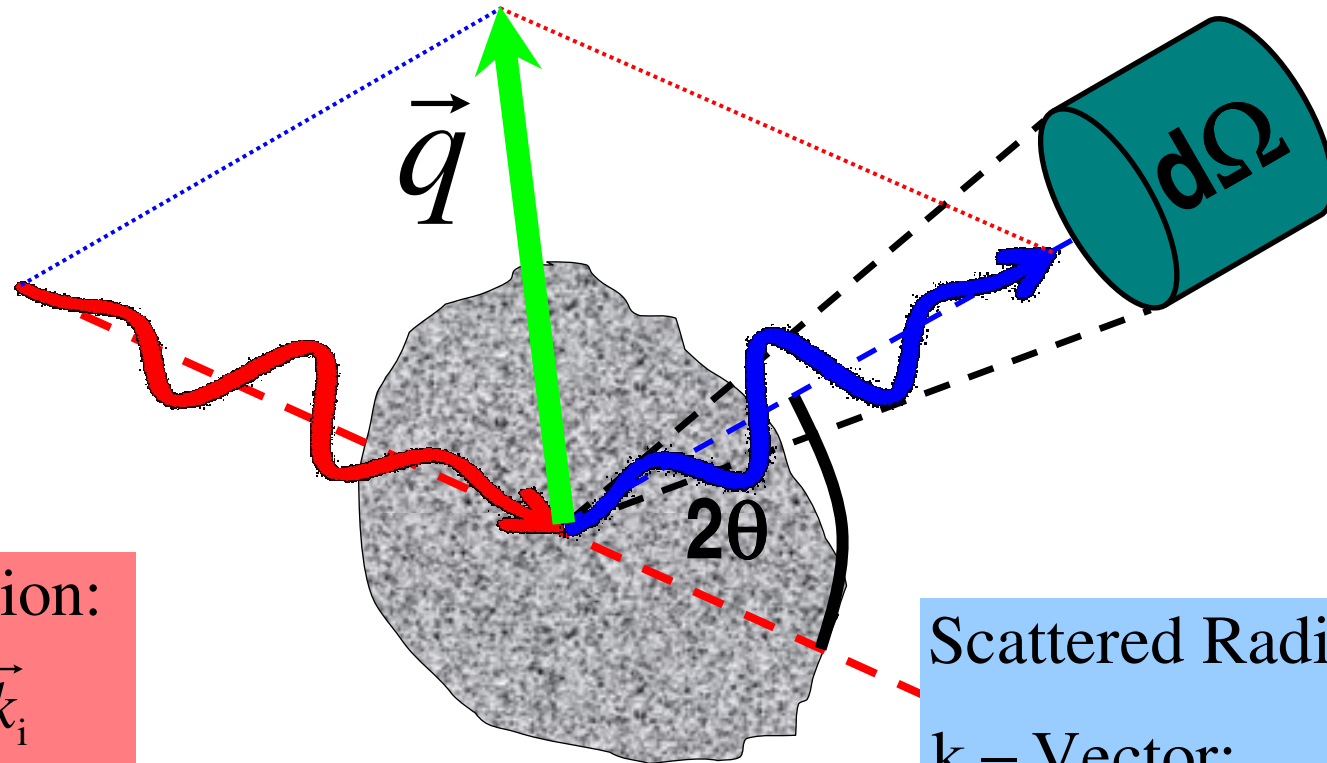
**8.01. 2009**

**Structure determination techniques**

**15.01.2009**

**Correlation Spectroscopy**

# Scattering Geometry



Incident Radiation:

k – Vector:  $\vec{k}_i$

$$|\vec{k}_i| = 2\pi / \lambda$$

Energy:  $E_i$

Polarization:  $\vec{p}_i$

Scattered Radiation:

k – Vector:  $\vec{k}_f$

Energy:  $E_f$

Polarization:  $\vec{p}_f$

Wavevector Transfer:

$$\vec{q} = \vec{k}_f - \vec{k}_i$$

Energy Transfer:

$$\Delta E = E_f - E_i = \hbar\omega$$

Polarization:

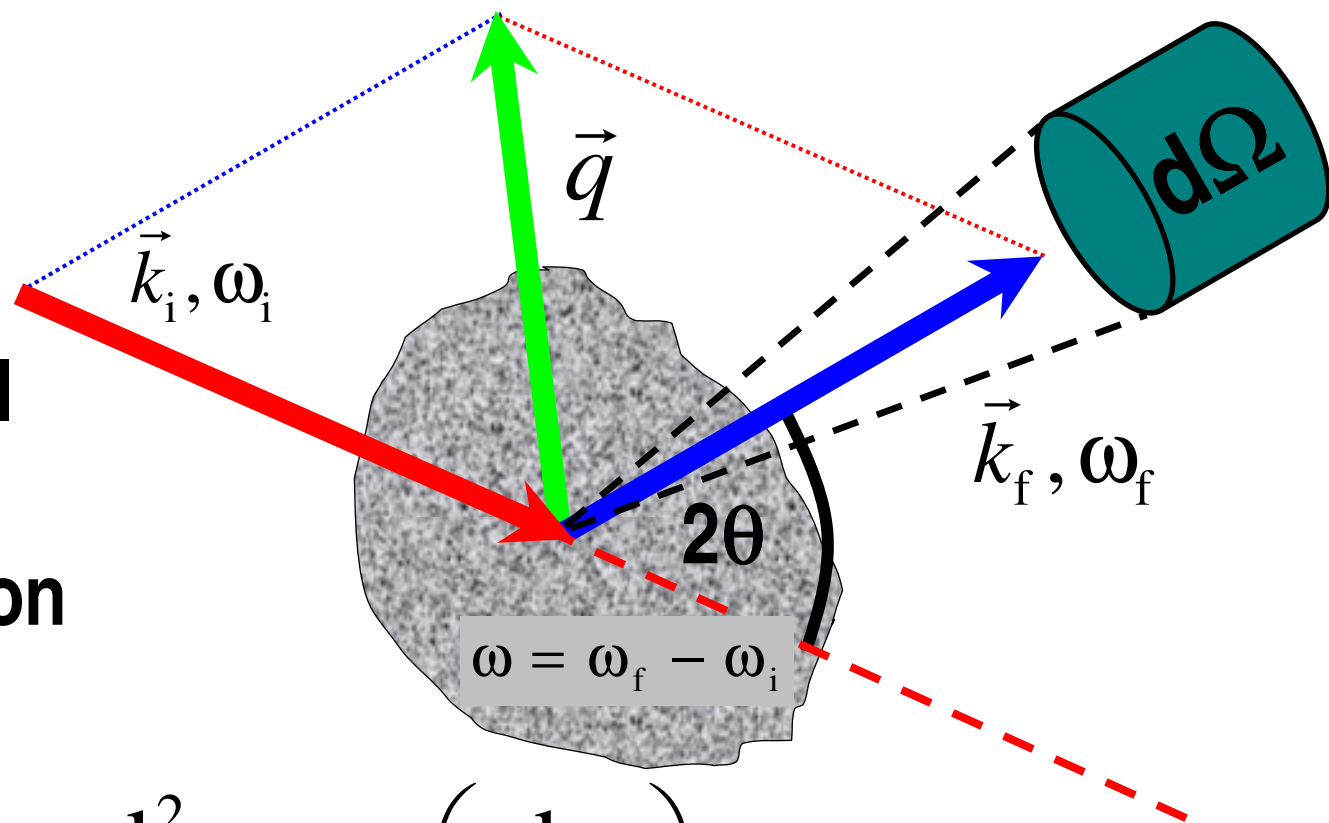
$$\vec{p}_i \rightarrow \vec{p}_f$$

For X-Rays:

$$\Delta E \ll E_f, E_i$$

$$\Rightarrow |\vec{q}| = 2k_i \sin(2\theta/2)$$

# Double-Differential Scattering Cross Section



$$\frac{d^2\sigma}{d\Omega d\omega_f} = \left( \frac{d\sigma}{d\Omega} \right)_0 S(\vec{q}, \omega)$$

**Intrinsic Cross Section  
Coupling Beam  $\Leftrightarrow$  Sample**

**Properties of  
the Sample  
without Beam**

$$\frac{d\sigma}{d\Omega} = \int d\omega \frac{d^2\sigma}{d\Omega d\omega} = \left( \frac{d\sigma}{d\Omega} \right)_0 S(\vec{q})$$

**Intrinsic Cross Section  
Coupling Beam  $\Leftrightarrow$  Sample**

**Properties of  
the Sample  
without Beam**

# Time and frequency dependent correlation functions

$$\rho(r, t) = e^{iHt/\hbar} \rho(r, 0) e^{-iHt/\hbar}$$

Time dependence in the Heisenberg representation

H - Hamilton operator of the system

$$\rho(r, t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-i\omega t} \rho(r, \omega)$$

time and frequency dependent electron densities are conjugate quantities connected via Fourier Transform

$$\rho(r, \omega) = \int_{-\infty}^{\infty} e^{i\omega t} \rho(r, t)$$

$$C(r, t, r', t') = \langle \rho(r, t) \rho(r', t') \rangle$$

time dependent correlation function

# Time and frequency dependent correlation functions

system spatially homogenous and temporal stationary

$$C(r_1, t_1, r_2, t_2) = \langle \rho(r_1, t_1) \rho(r_2, t_2) \rangle = \langle \rho(r_1 - r_2, t_1 - t_2) \rho(0, 0) \rangle = C(r, t)$$

$$r = r_1 - r_2, \text{ and } t = t_1 - t_2$$

Intermediate scattering function

$$S(q, t) = \iint dr_1 dr_2 \langle \rho(r_1, 0) \rho(r_2, t) \rangle e^{iq(r_1 - r_2)} = \langle \rho(-q, 0) \rho(q, t) \rangle = \tilde{C}(q, t)$$

Dyanmic scattering function

$$S(q, \omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\omega e^{-i\omega t} S(q, t) = \tilde{C}(q, \omega)$$

Static scattering function

$$S(q, t = 0) = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\omega S(q, \omega)$$

How to calculate the dynamic correlation functions?

### 1. Linear response theory

$$H = \int dr F(r, t) \rho(r, t) \quad \begin{array}{l} H \text{ interaction Hamiltonian} \\ F \text{ Force} \end{array}$$

$$\rho(r, t) = \iint dr' dt' \chi(r, r', t - t') F(r', t')$$

$$\rho(r, \omega) = \int dr' \chi(r, r', \omega) F(r', \omega)$$

response function  $\chi$

### Fluctuation Dissipation Theorem

$$S(q, \omega) = \frac{2k_B T}{\omega} \text{Im} \chi(q, \omega)$$

valid for thermal equilibrium



# Fluctuation Dissipation Theorem

Example : the driven, damped harmonic oscillator

equation of motion  $\ddot{x} + \omega_0^2 x + \gamma \dot{x} = f / m$

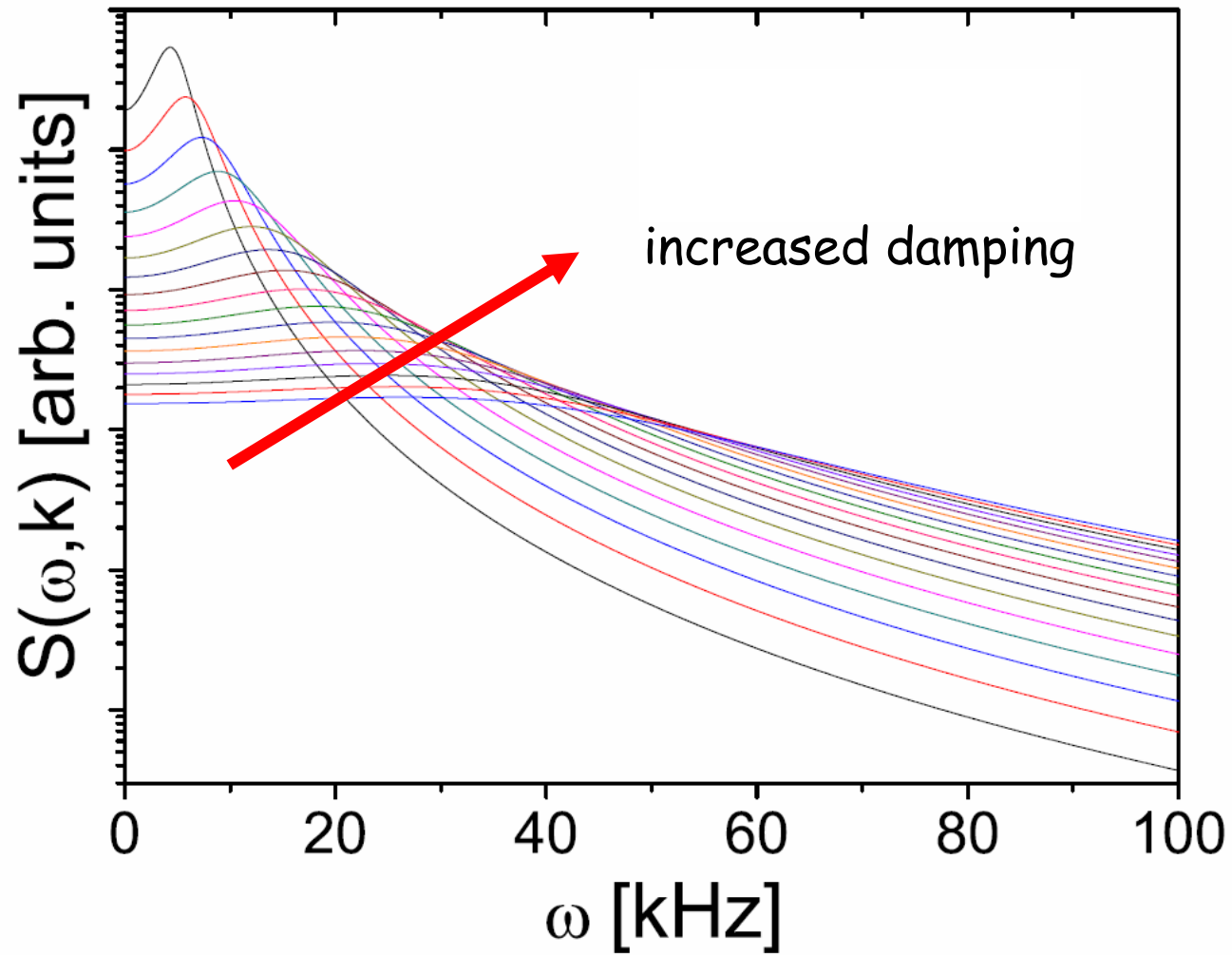
frequency dependent response  $\chi(\omega) = \frac{x(\omega)}{f(\omega)} = \frac{1}{m} \frac{1}{-\omega^2 + \omega_0^2 - i\omega\gamma}$

$$\text{Im } \chi(\omega) = \frac{1}{2m\omega_1} \left[ \frac{\gamma/2}{(\omega - \omega_1)^2 + (\gamma/2)^2} - \frac{\gamma/2}{(\omega + \omega_1)^2 + (\gamma/2)^2} \right]$$

$$\omega_1 = \sqrt{\omega_0^2 - \gamma^2/4}, \quad \tau^{-1} = \frac{1}{2}\gamma[1 - (1 - 4\omega_0^2\gamma^{-2})^{1/2}]$$

- (a)  $\omega_1$  is real  $\rightarrow$  Lorentzian lines, centered at  $\omega_1$  with HWHM  $\gamma/2$
- (b)  $\omega_1$  is imag  $\rightarrow$  Lorentzian line, centered at 0 with HWHM  $\tau$

# $S(q, \omega)$ of the damped harmonic oscillator



How to calculate the dynamic correlation functions?

## 2. Differential equations of correlation functions

Example : Fick's law of diffusion

$$\frac{\partial}{\partial t} C(R, t) = D \nabla^2 C(R, t) \quad (1)$$

Spatial Fourier transform of Eq. (1)

$$\frac{\partial}{\partial t} S(q, t) = -Dq^2 S(q, t) \quad (2)$$

Intermediate Scattering function for free diffusion

$$S(q, t) = \exp(-Dq^2 t) \quad (3)$$

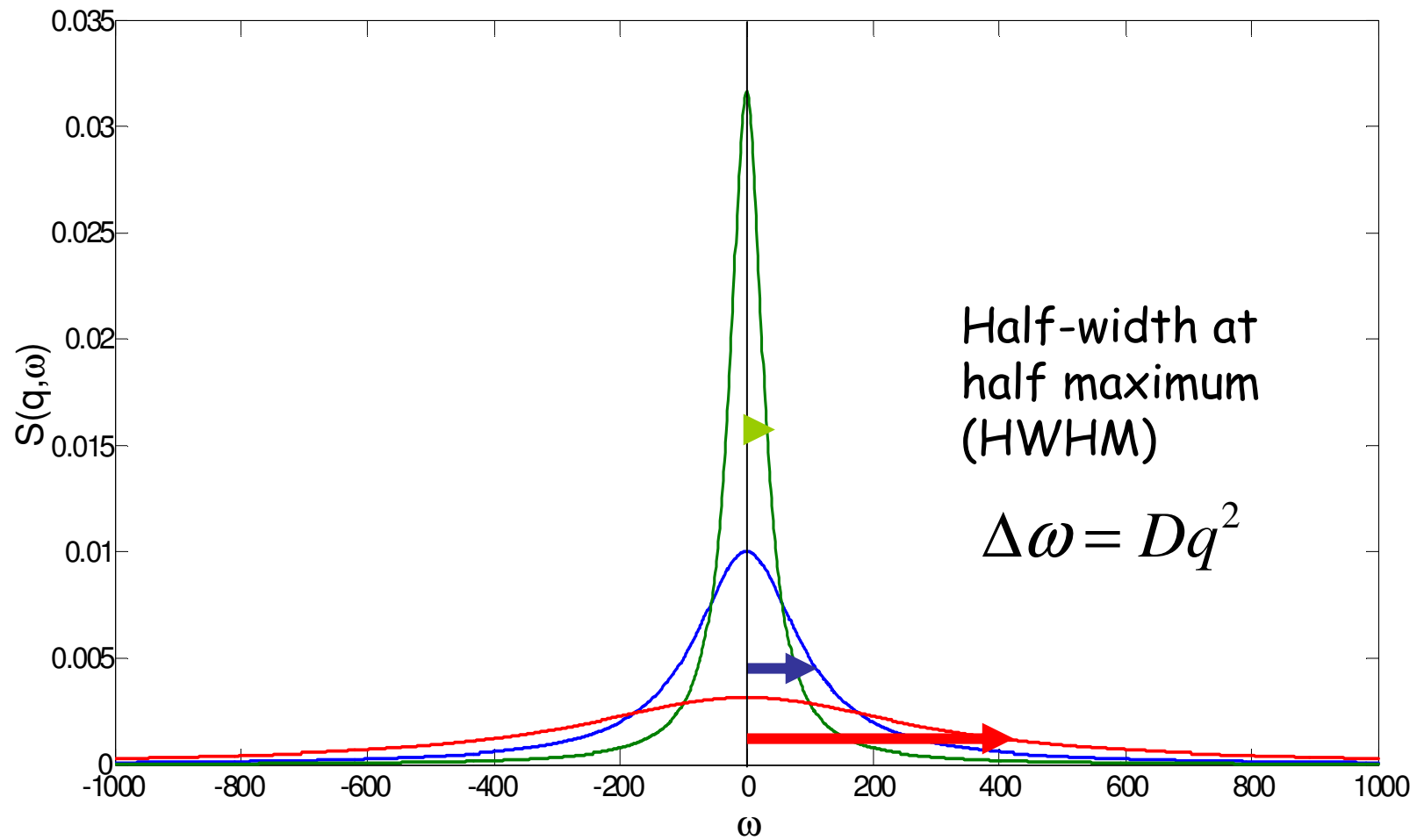
Dynamic Scattering function for free diffusion

$$D = \frac{k_B T}{6\pi\eta R}$$

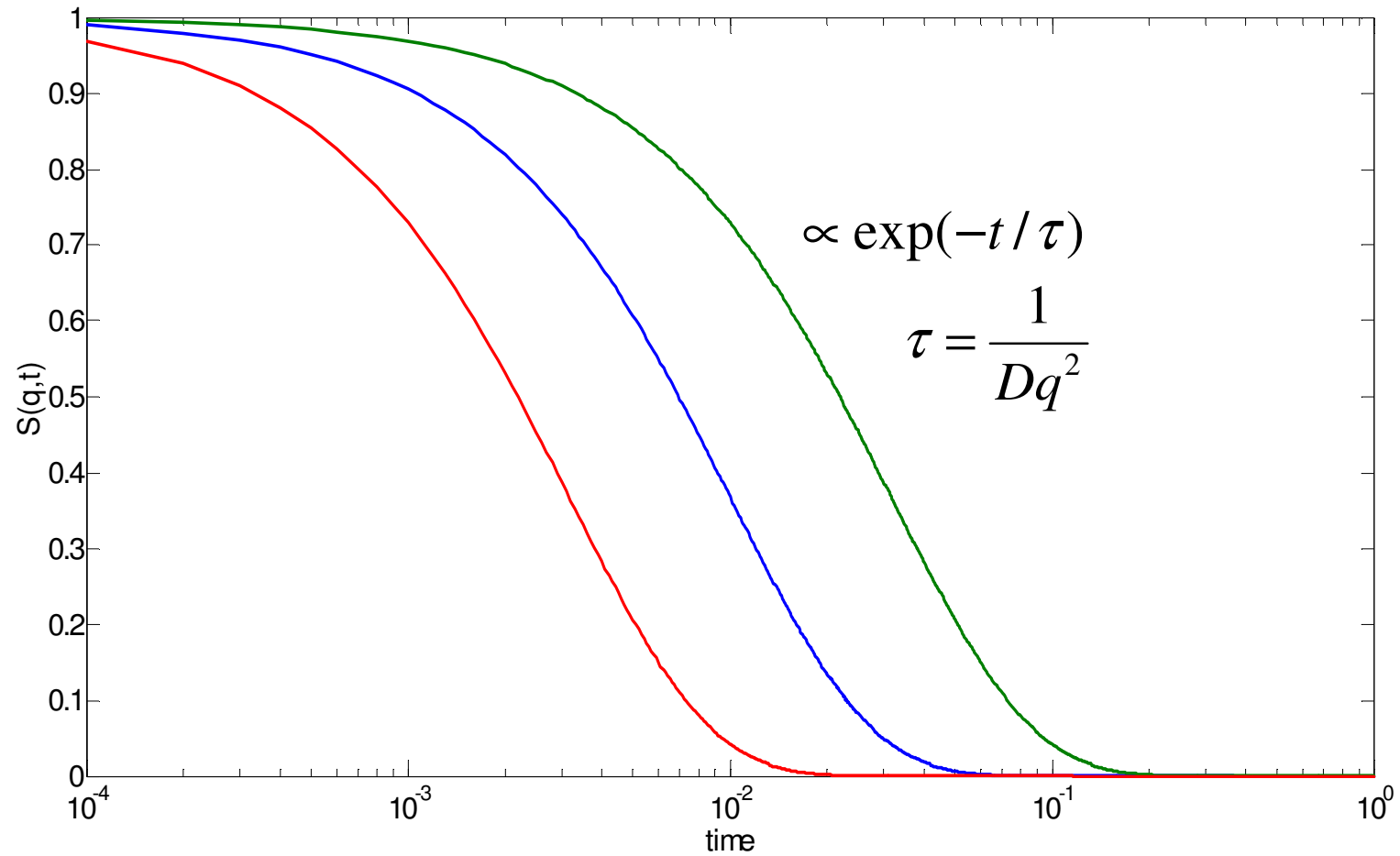
$$S(q, \omega) = \frac{Dq^2}{\omega^2 + (Dq^2)^2} \quad (4)$$

Einstein relation

# Dynamic Scattering function for free diffusion

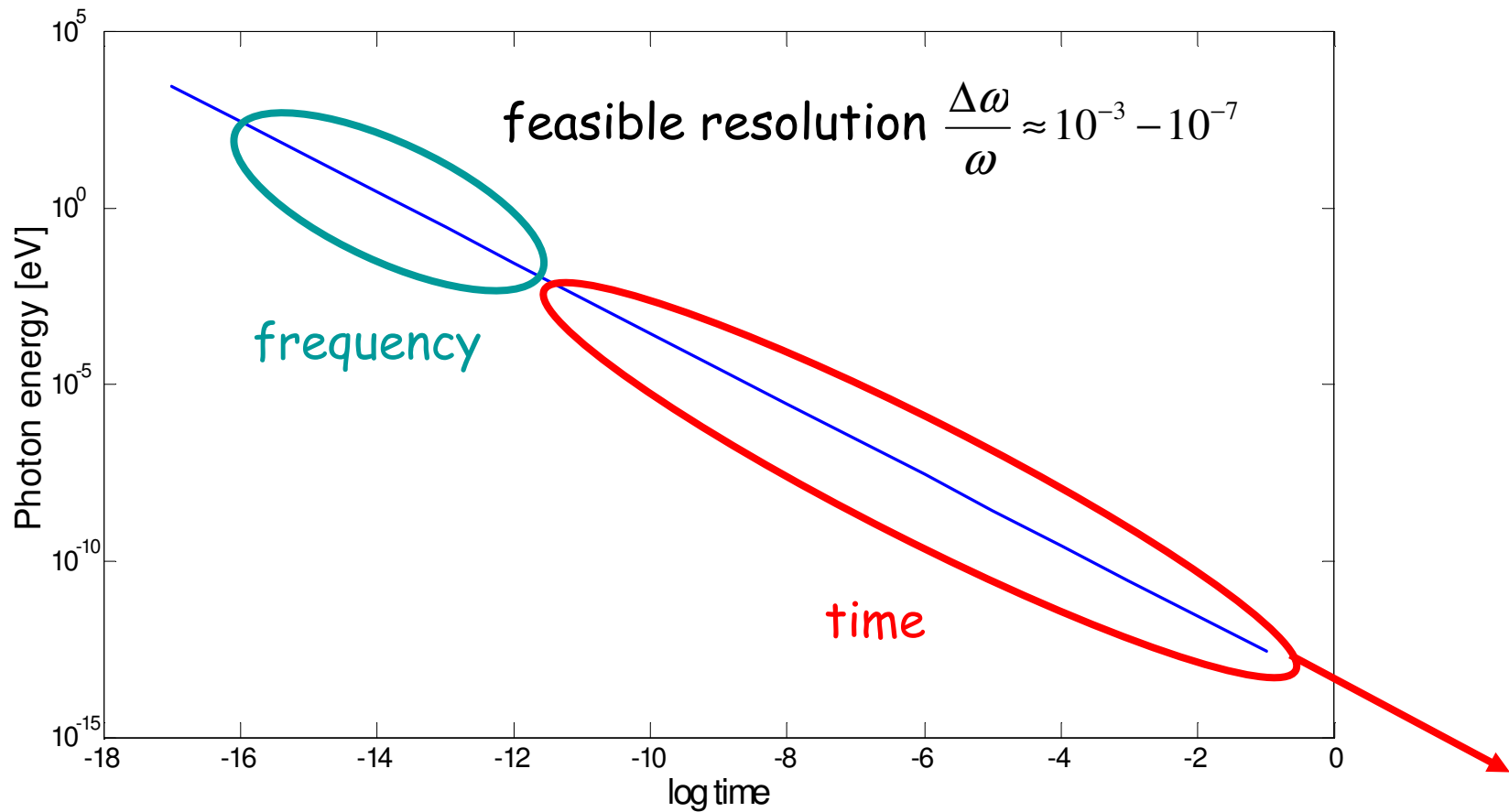


# Intermediate Scattering function for free diffusion

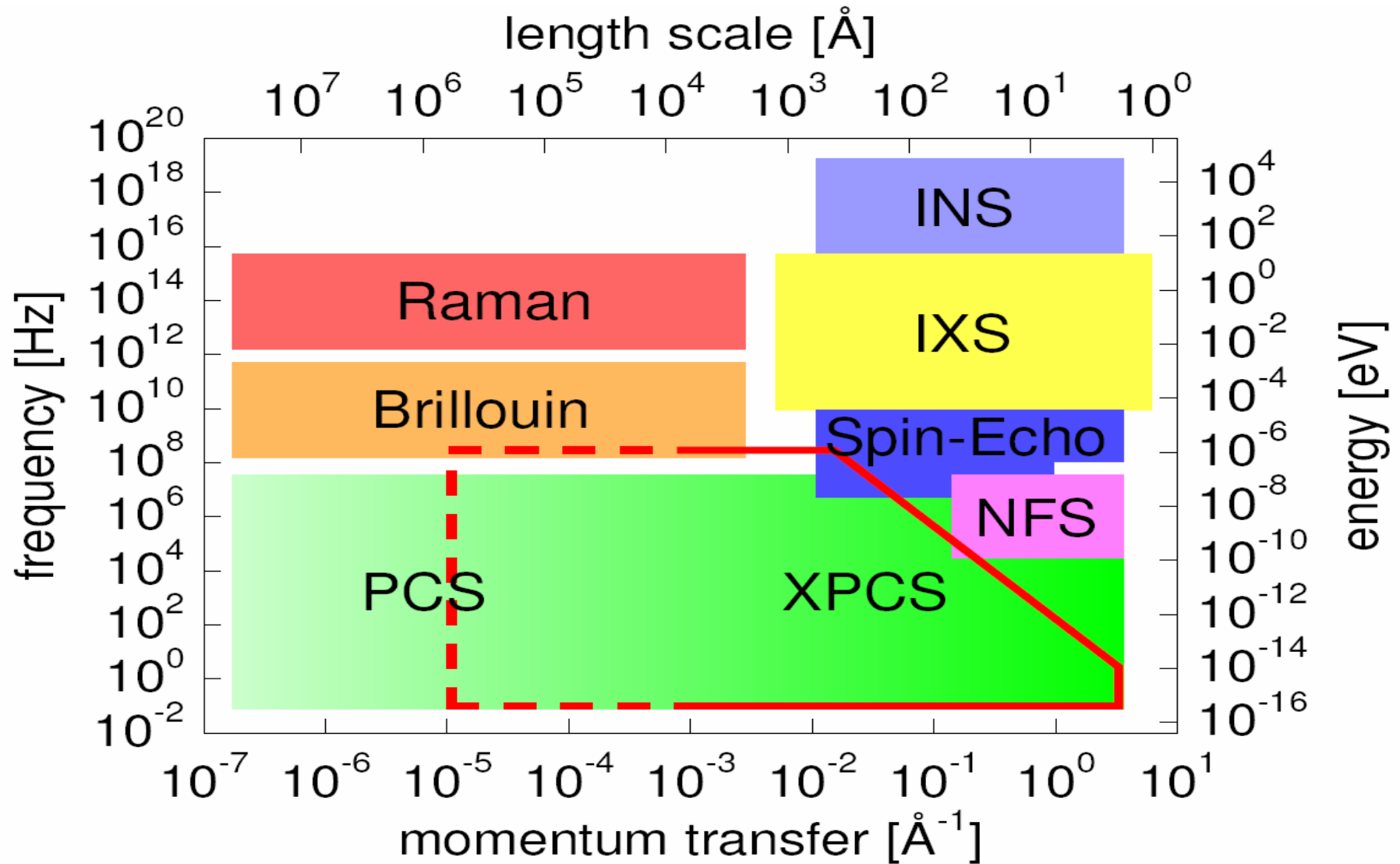


# Measuring in frequency or time domain ? equilibrium processes

typical frequency of X-rays     $h\nu = 8\text{keV}$      $h = 4.14 \cdot 10^{-15} \text{ eVs}$



# Measuring in frequency or time domain ?



# Time domain

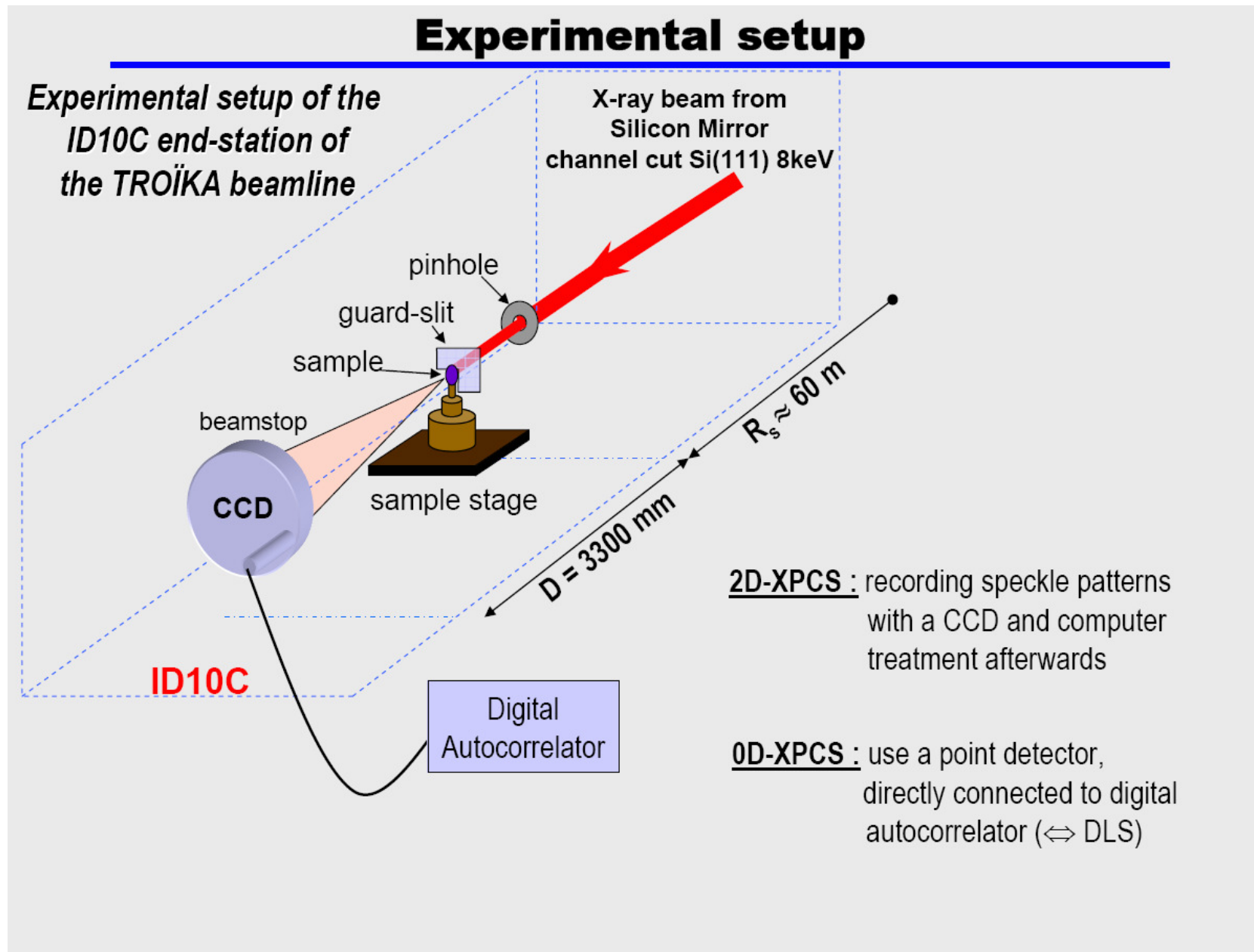
- + cover a very large time window 1000 seconds  $\rightarrow$   $1e-12$  seconds (XFEL)
- + sensitive to non-equilibrium processes
- photon hungry



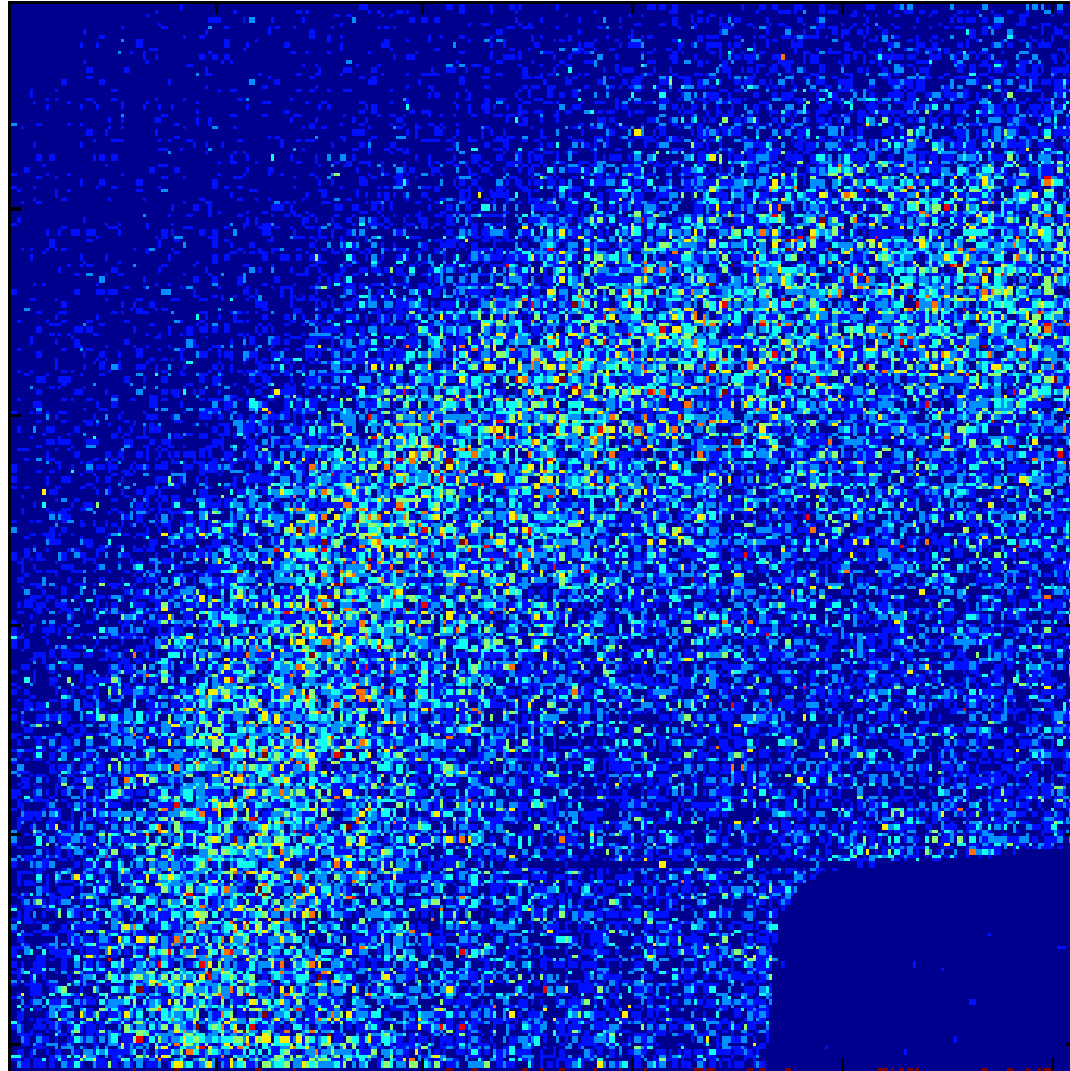
**XPCS is a  
„Photon-  
Hungry“  
Method ....**



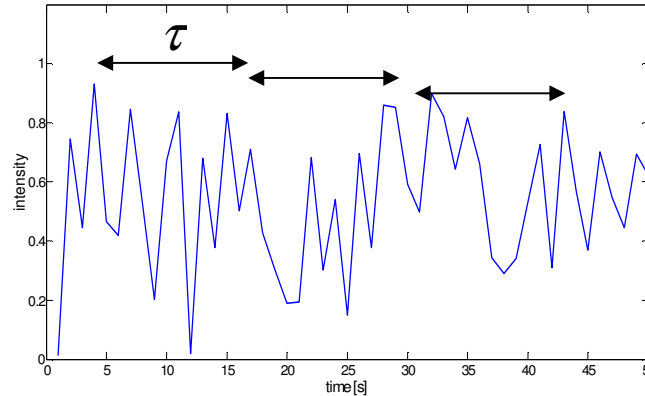
# X-ray Photon Correlation Spectroscopy



# Fluctuating Speckle Pattern



# Intensity Autocorrelation Function



$$\langle I(q,0)I(q,\tau) \rangle = \frac{1}{T} \int_0^T I(q,t)I(q,t+\tau)$$

remember  $I(q,t) = \langle \rho(q,t)\rho(-q,t) \rangle$

$$\langle I(q,t_1)I(q,t_2) \rangle = \langle \rho(q,t_1)\rho(-q,t_1)\rho(q,t_2)\rho(-q,t_2) \rangle$$

## Gaussian momentum theorem and Siegert relation

$$\begin{aligned} \langle \rho(q, t_1) \rho(-q, t_1) \rho(q, t_2) \rho(-q, t_2) \rangle &= \langle \rho(q, t_1) \rho(-q, t_1) \rangle \langle \rho(q, t_2) \rho(-q, t_2) \rangle \\ &+ \underbrace{\langle \rho(q, t_1) \rho(q, t_2) \rangle \langle \rho(-q, t_1) \rho(-q, t_2) \rangle + \langle \rho(q, t_1) \rho(-q, t_2) \rangle \langle \rho(q, t_1) \rho(-q, t_2) \rangle}_{= 0} \end{aligned}$$

$$\begin{aligned} \langle \rho(q, t_1) \rho(-q, t_1) \rho(q, t_2) \rho(-q, t_2) \rangle &= \langle \rho(q, t_1) \rho(-q, t_1) \rangle \langle \rho(q, t_2) \rho(-q, t_2) \rangle \\ &+ \langle \rho(q, t_1) \rho(-q, t_2) \rangle \langle \rho(q, t_1) \rho(-q, t_2) \rangle \end{aligned}$$

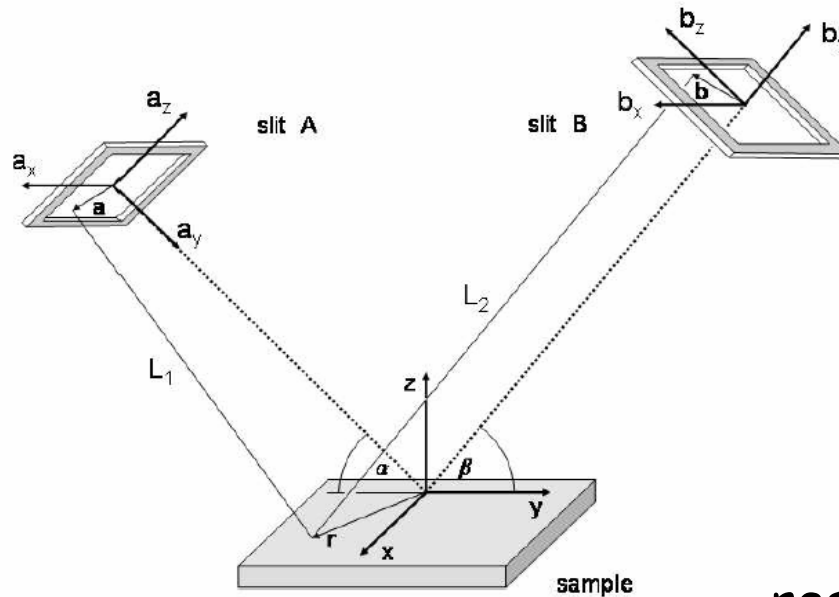
$$= \langle \rho(q, 0) \rho(-q, 0) \rangle \langle \rho(q, 0) \rho(-q, 0) \rangle + \langle \rho(q, 0) \rho(-q, \tau) \rangle \langle \rho(q, 0) \rho(-q, \tau) \rangle$$



$$\langle I(q, 0) I(q, \tau) \rangle = \langle S(q, 0) \rangle^2 + \left| \langle S(q, \tau) \rangle \right|^2$$

phase information lost, if S is complex (QM)  $k_B T \ll \hbar \omega$

# Effects of partial coherence



resolution functions  $F, H$

$$g_2(q, \tau) = \frac{\langle I(q, 0) I(q, \tau) \rangle}{\langle I(q, 0) \rangle^2} = 1 + \frac{\iint dq' dq'' \tilde{C}(q', \tau) \tilde{C}(q'', \tau) F(q, q', q'')}{\left( \int dq' \tilde{C}(q', 0) H(q, q') \right)^2}$$

approximation

$$g_2(q, \tau) = \frac{\langle I(q, 0) I(q, \tau) \rangle}{\langle I(q, 0) \rangle^2} \approx 1 + \beta^2 \frac{|I(q, \tau)|^2}{\langle I(q, 0) \rangle^2}$$

comprising effects of partial coherence

## Heterodyne mixing in XPCS

analogue to Holography - built in a reference source

$$\rho(q,t) = \rho_0 + \rho(q,t)$$

$$\langle I(q,0)I(q,\tau) \rangle \sim 2I_s I_r \langle S(q,\tau) \rangle + I_s^2 |\langle S(q,\tau) \rangle|^2$$

field correlation function

$$I_s = \langle \rho(q,t)\rho(-q,t) \rangle, I_r = \langle \rho_0 \rho_0^* \rangle$$

By choosing a strong reference signal the intensity autocorrelation function is dominated by  $S(q,t)$  - on the expense of signal to noise ratio.

# Examples

Surface fluctuations

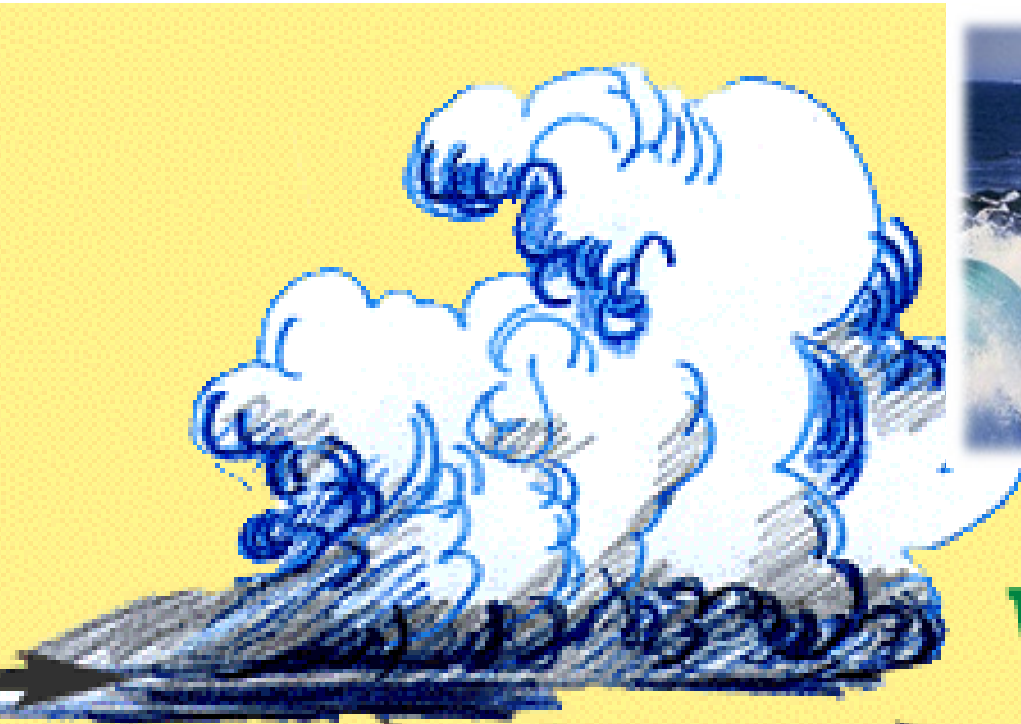
Heterodyne mixing

Transition propagating - overdamped behavior

Bulk fluctuations

Colloidal diffusion

Domain dynamics in Chromium

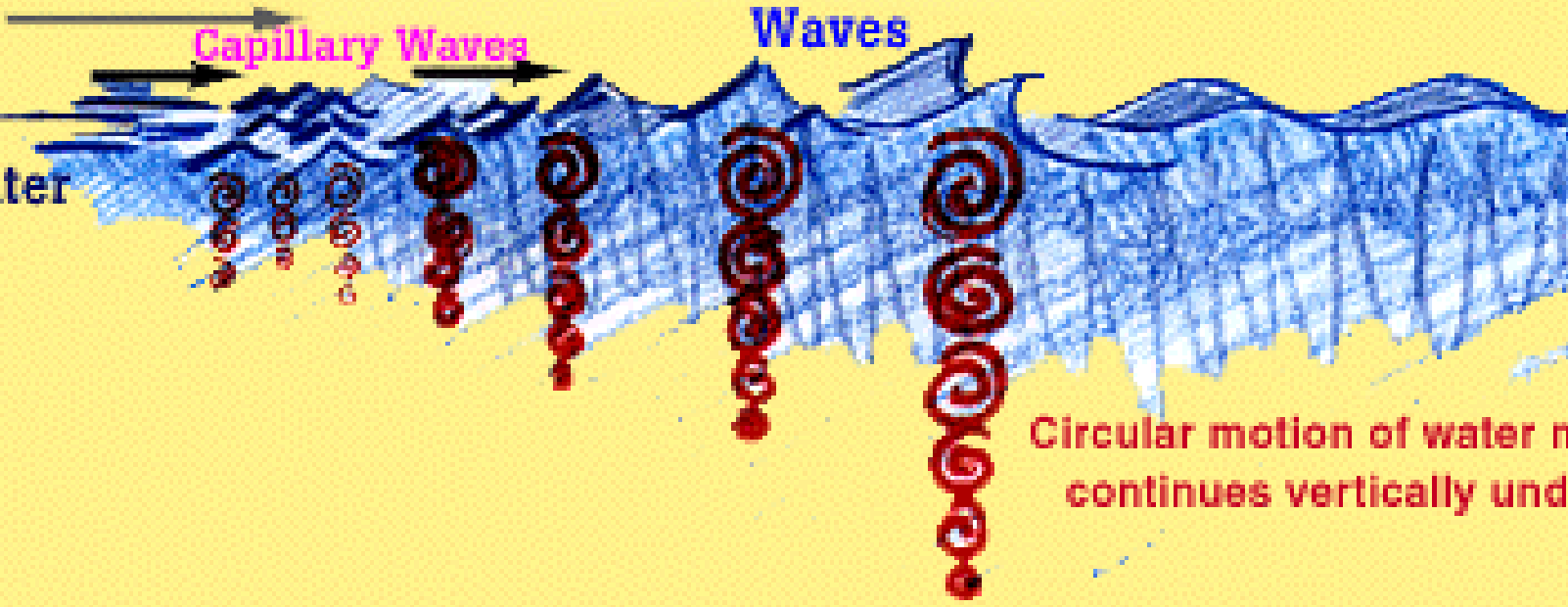


**WIND**

Smooth Water Surface

Capillary Waves

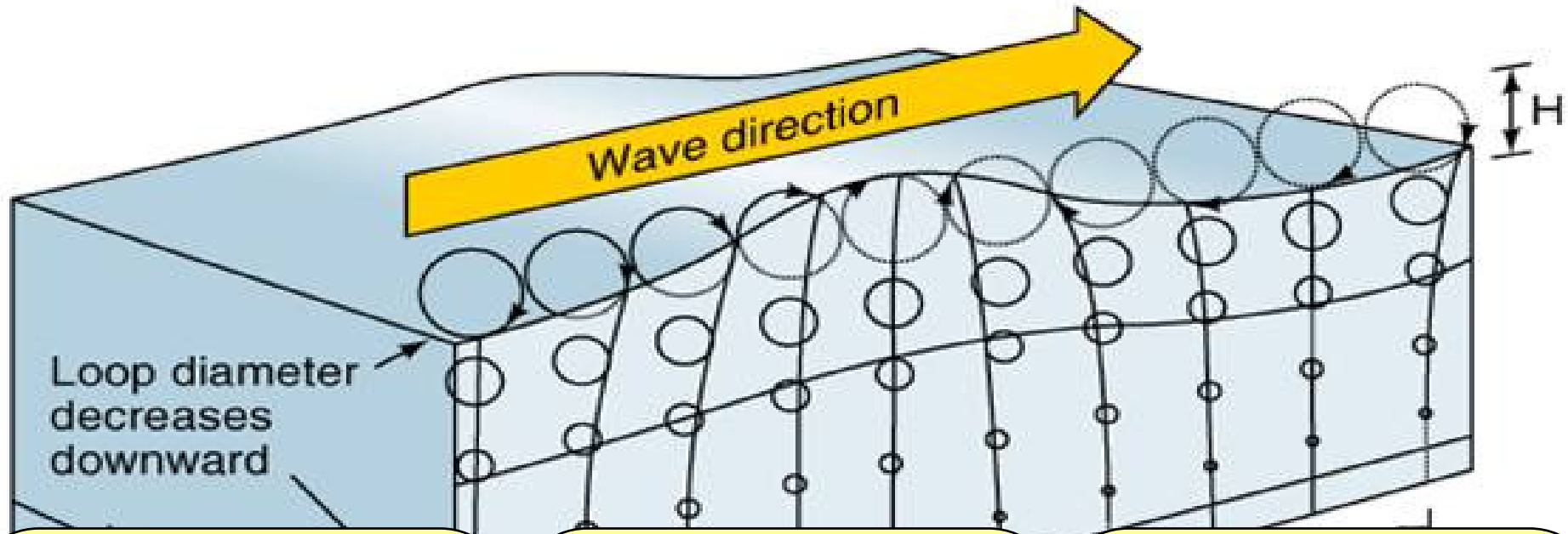
Waves



Circular motion of water particles continues vertically under



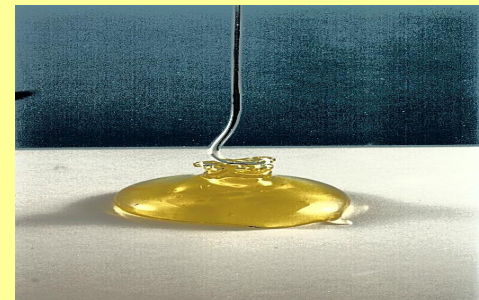
# Fluctuating Interfaces: Capillary Waves



**Thermal energy**

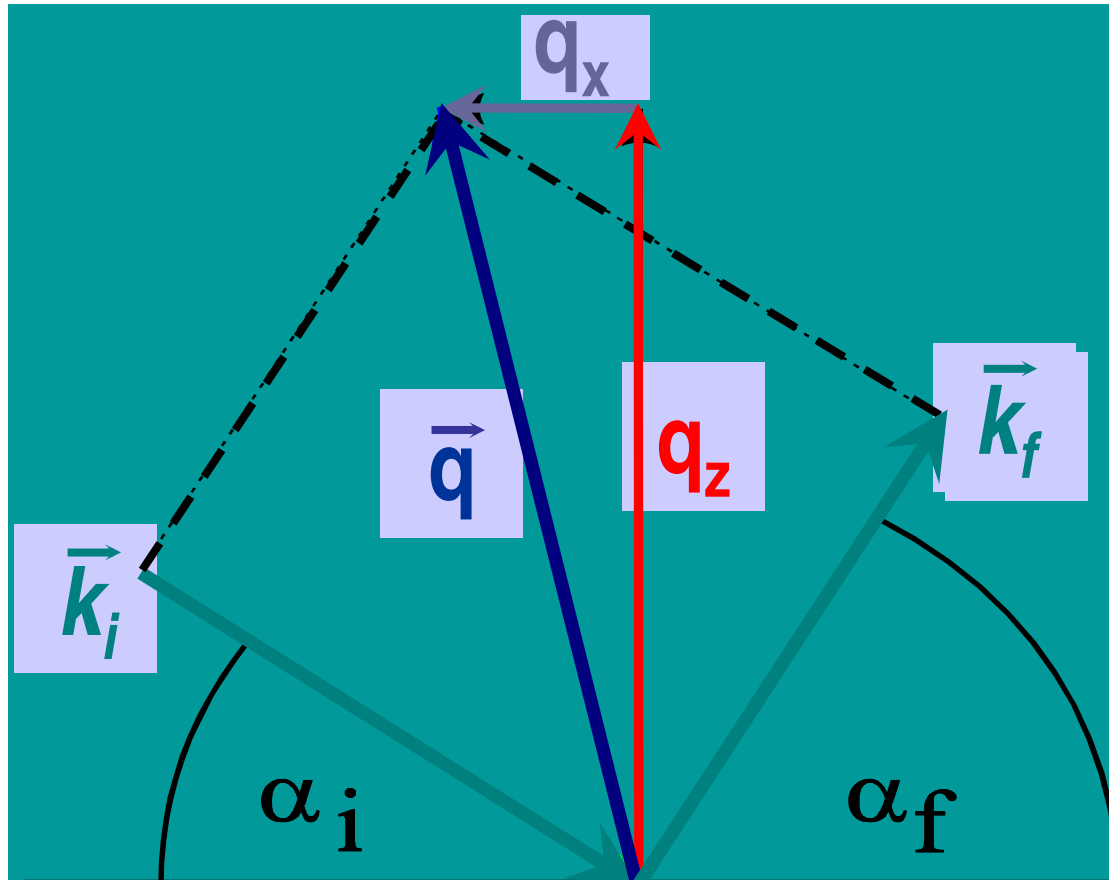


**Surface tension**



**Viscosity**

# Scattering Geometry & Notation



**Reflectivity:**

$$q_x = q_y = 0$$

$$q_z = (4\pi/\lambda)\sin\alpha_i$$

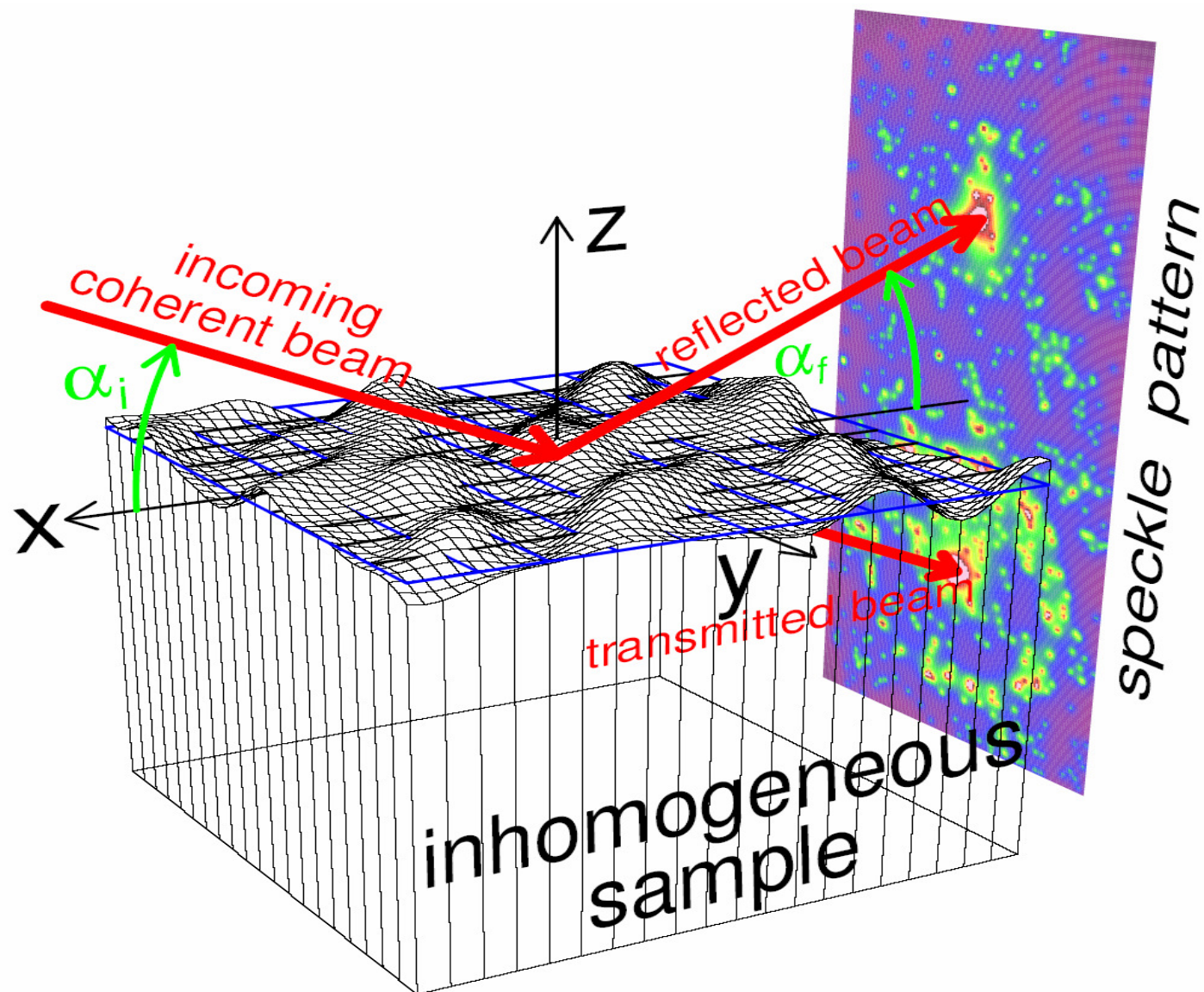
**Diffuse Scattering:**

$$q_x, q_y \neq 0$$

$$q_z = (4\pi/\lambda)\sin(\alpha_i + \alpha_f)/2$$

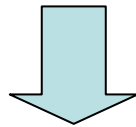
**Wave-Vector:  $\vec{q} = \vec{k}_f - \vec{k}_i$**

$$\alpha_i, \alpha_f < 5^\circ$$

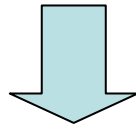


# Electron Density: Surface Scattering

$$\tilde{\rho}(q, t) = \frac{i \rho_0}{q_z} \iint \exp \left\{ -i [q_x x + q_y y + q_z z(x, y, t)] \right\} dx dy$$



$$|q_z z(x, y, t)| \ll 1$$



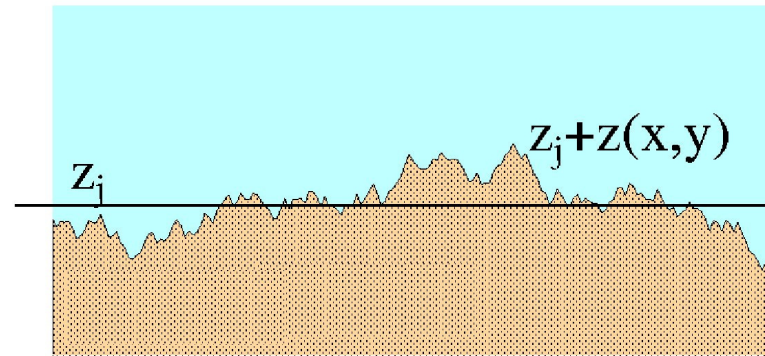
$$\tilde{\rho}(q, t) \approx \delta(q_x) \delta(q_y) + \rho_0 \tilde{z}(q_x, q_y, t)$$



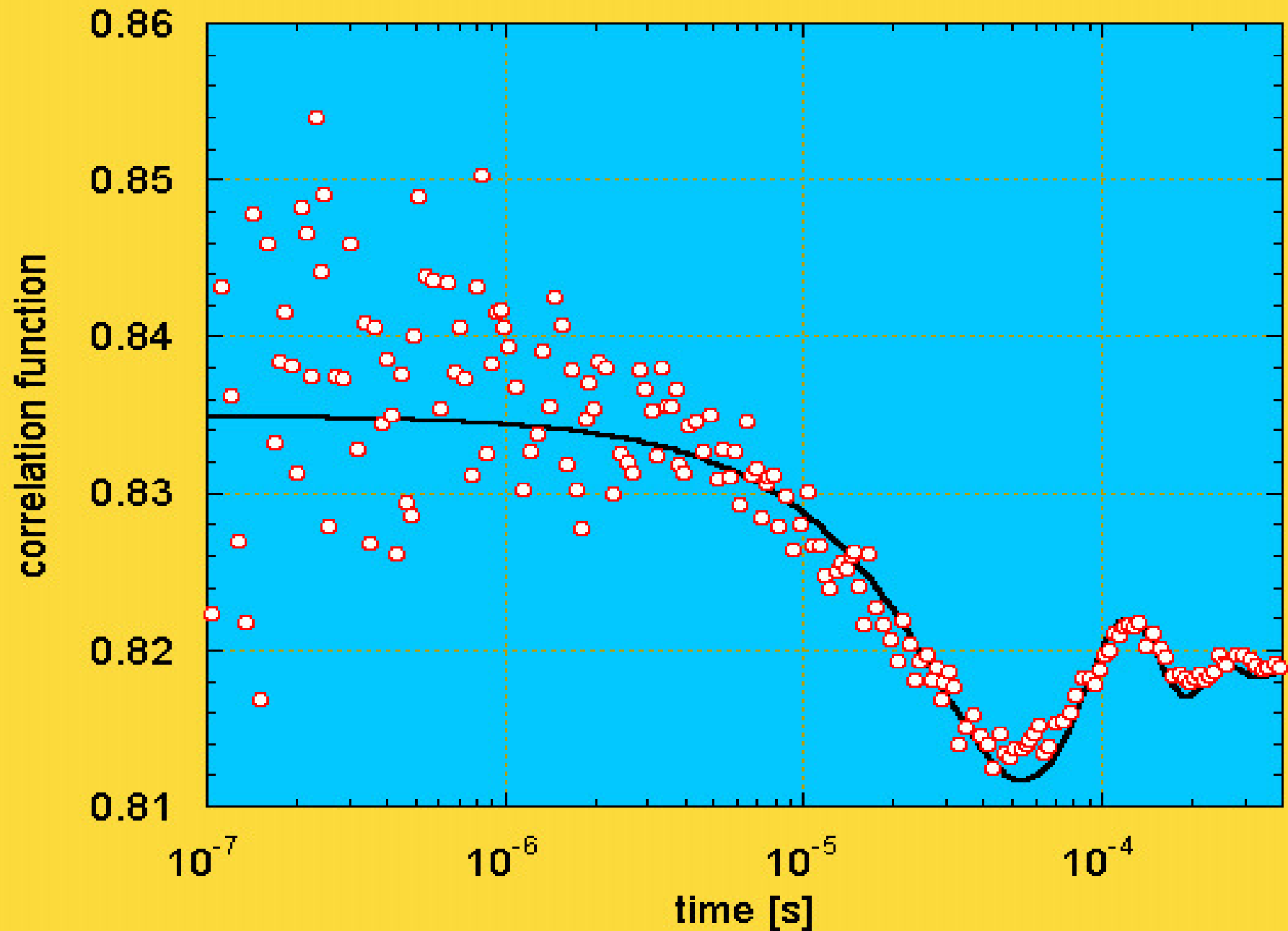
reference source for heterodyning



fluctuating signal



# Capillary waves on liquid water

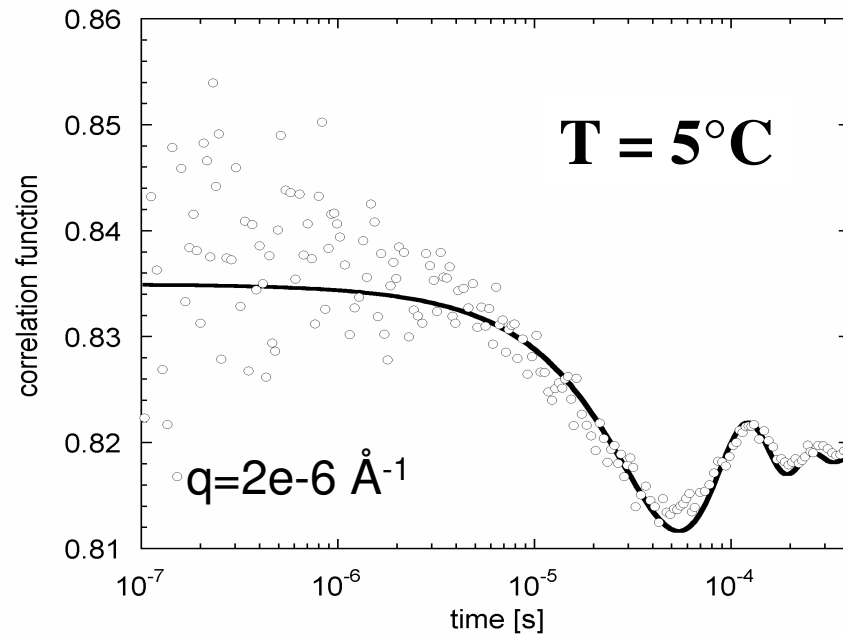


# Correlation functions of surface fluctuations

$$\omega_s = \sqrt{\frac{\gamma}{\rho}} q^3$$

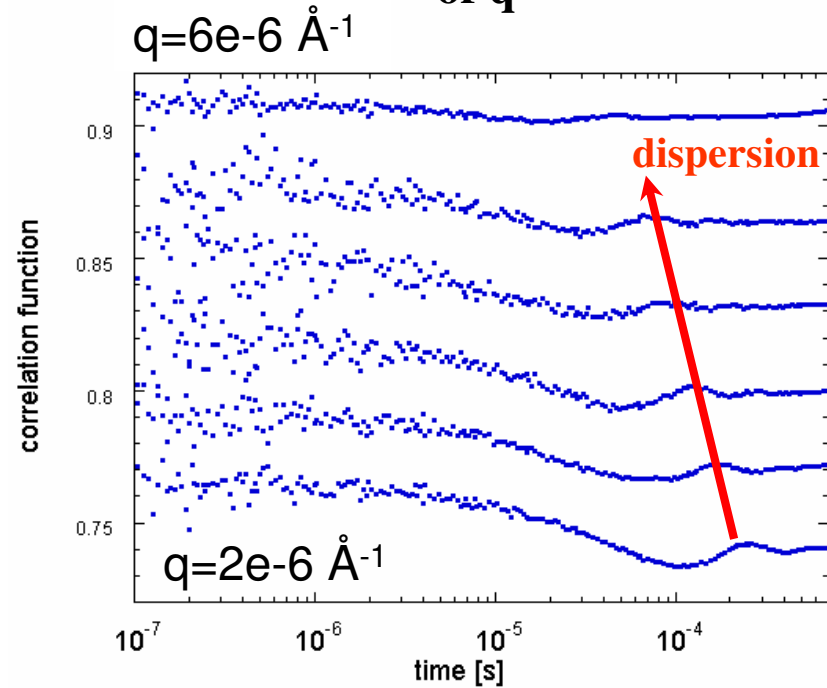
$$\langle I(q,0)I(q,\tau) \rangle \sim 2I_s I_r \langle S(q,\tau) \rangle + I_s^2 \left| \langle S(q,\tau) \rangle \right|^2$$

$$\cos(\omega\tau)e^{-\Gamma\tau} \quad \cos^2(\omega\tau)e^{-2\Gamma\tau}$$



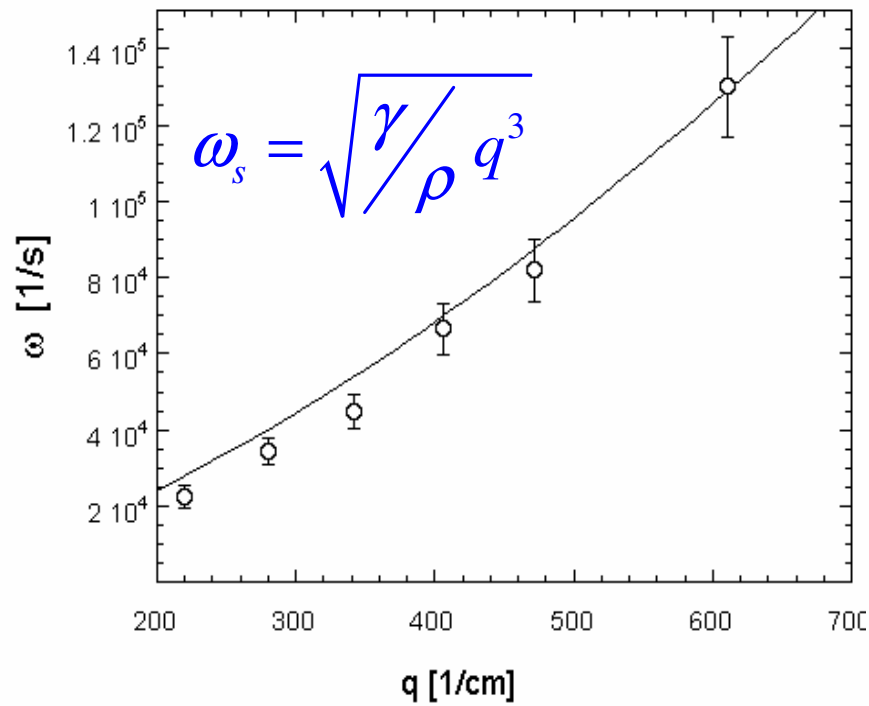
damped cosine behavior  
 → heterodyne mixing

correlation functions of a water surface at T=5°C as a function of q

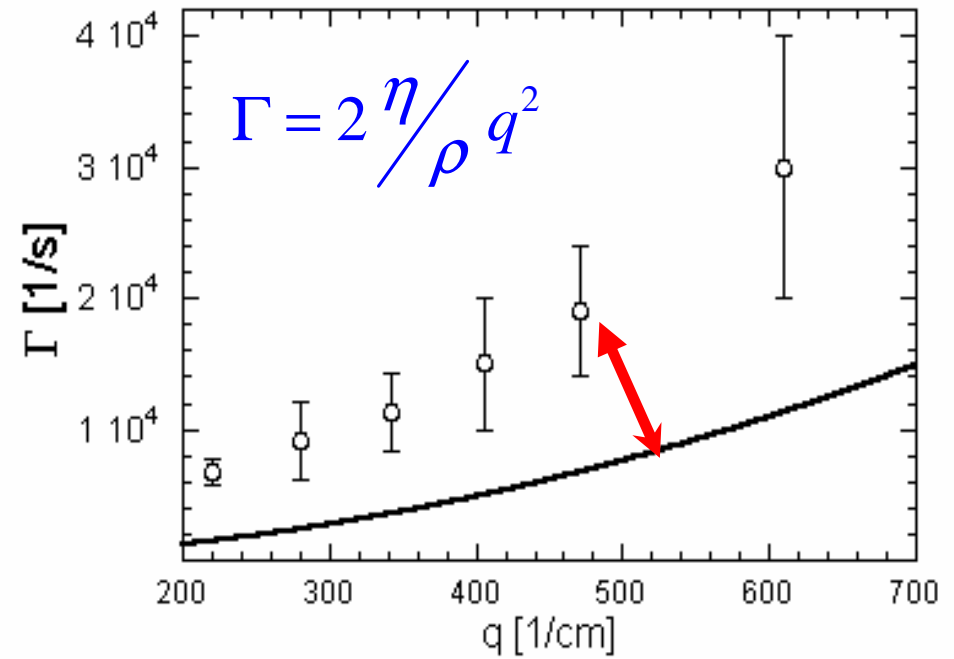


# Correlation functions of surface fluctuations

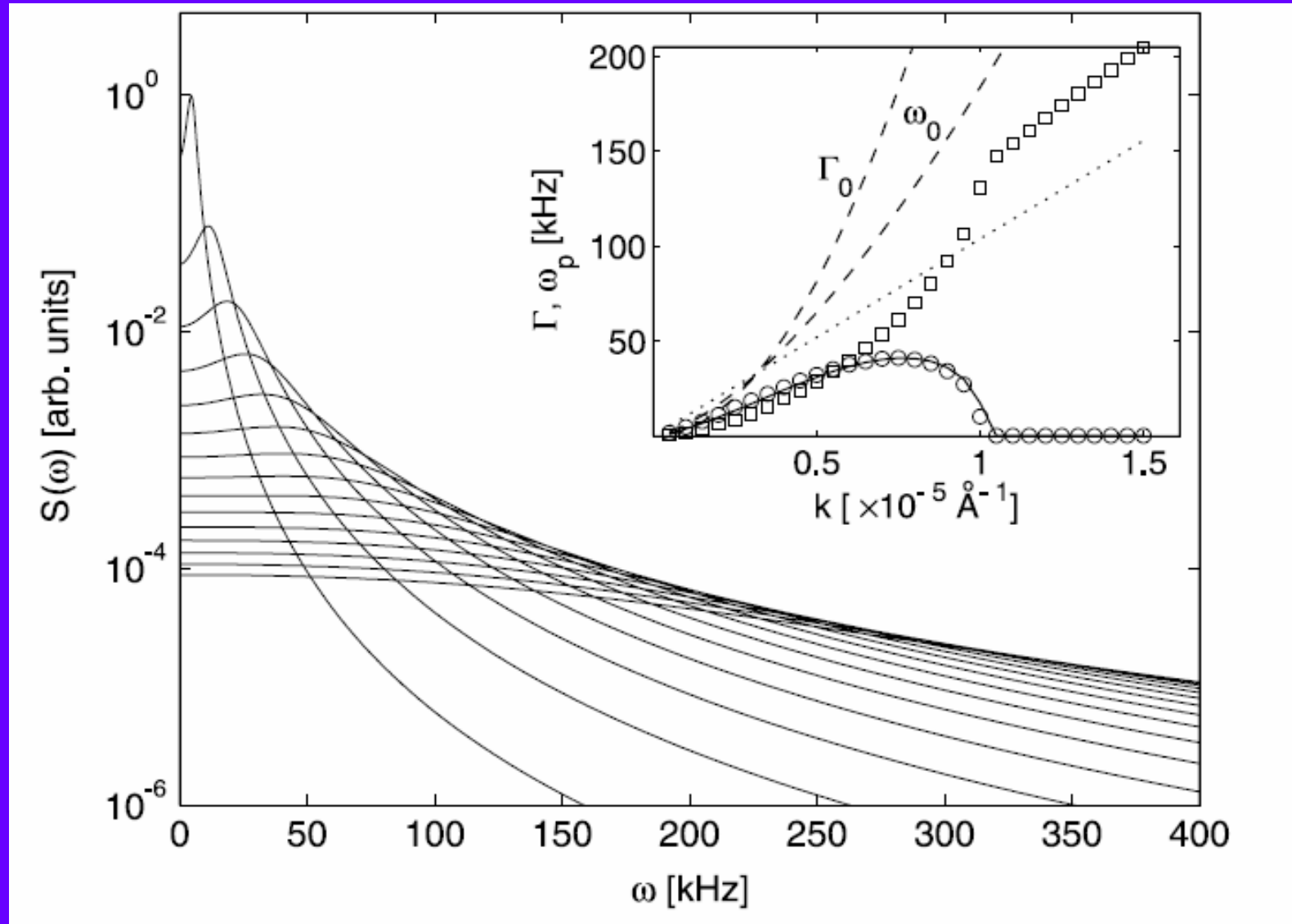
## Dispersion



## Damping

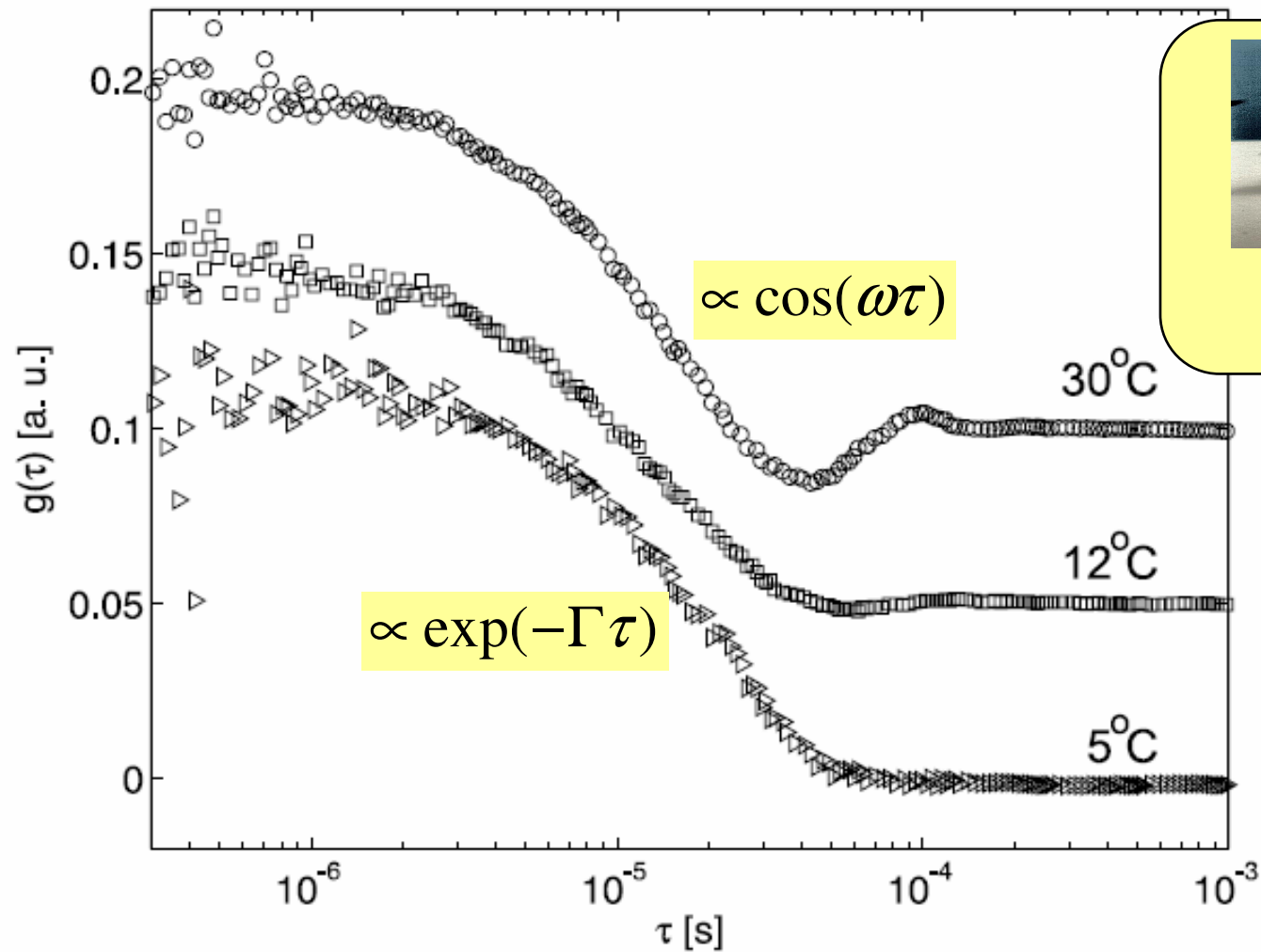


# Spectral features of damped capillary waves





# Mixtures of liquid water and glycerol tuning the viscosity (damping) as a function of temperature



Viscosity ( $T$ )

# Measurement

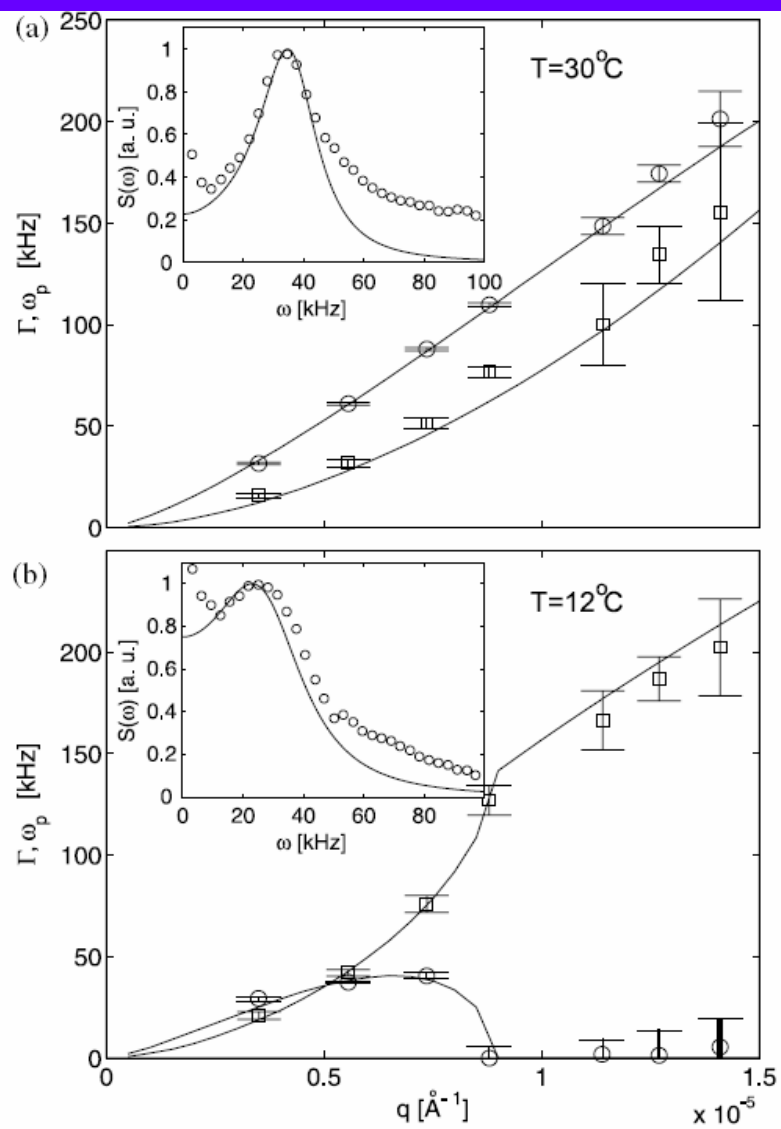
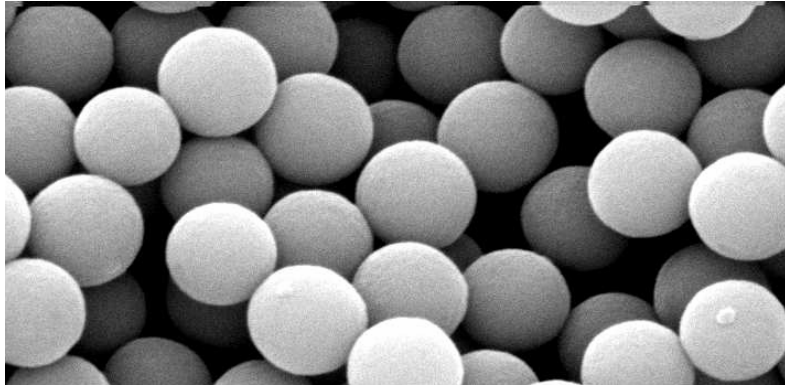


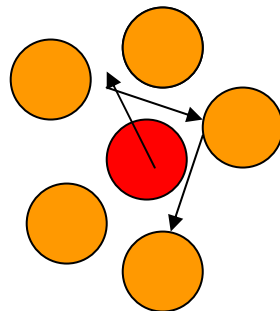
FIG. 4. Experimental data (from Fig. 2) that are well de-

# Diffusion in Colloidal Suspensions

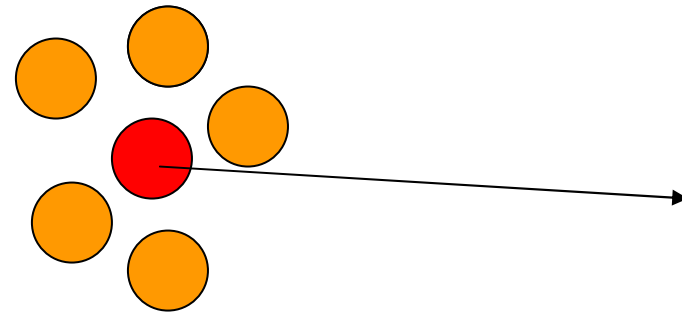


Free diffusion is valid for large length scales. What is happening when the probed length scales becomes comparable to particle size ?

caging



free diffusion



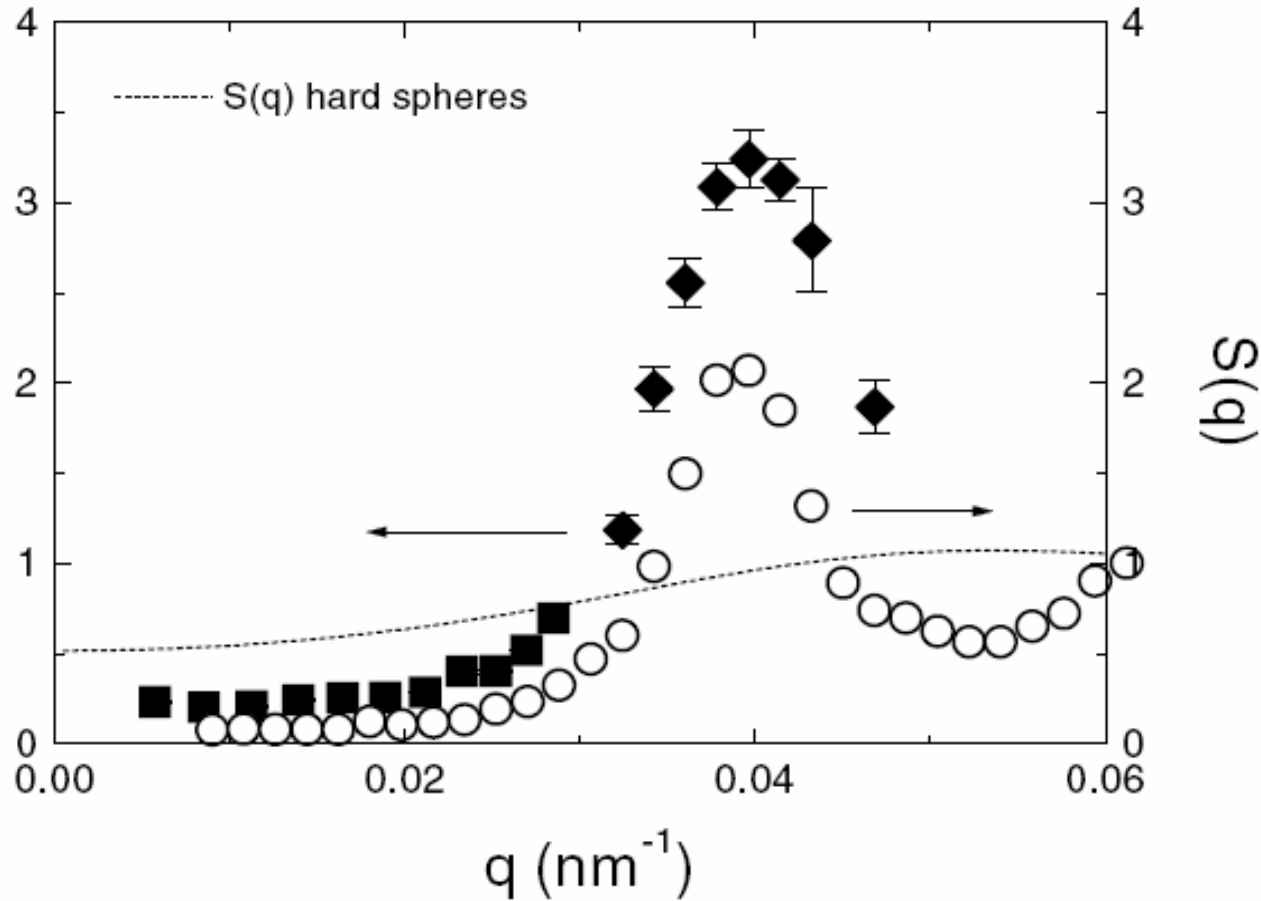
# Diffusion in Colloidal Suspensions

$$D(q) = \frac{\Gamma(q)}{q^2}$$

de Gennes narrowing  
microscopic structure of the liquid  
becomes visible

$$D_0 = \frac{k_B T}{6\pi\eta R}$$

$D_0 / D(q)$



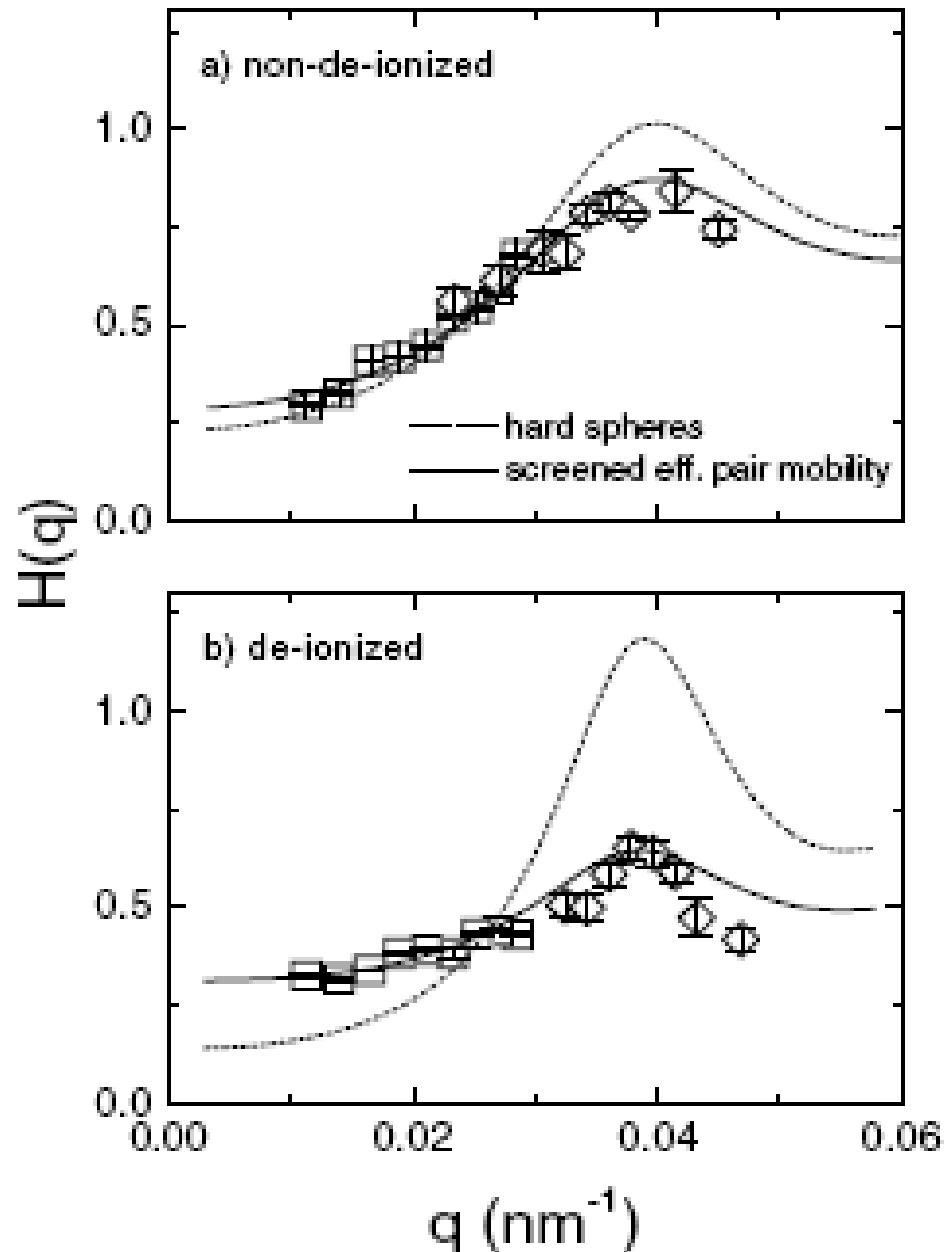
# Hydrodynamic function $H(q)$

indirect hydrodynamic interactions mediated by the solvent medium

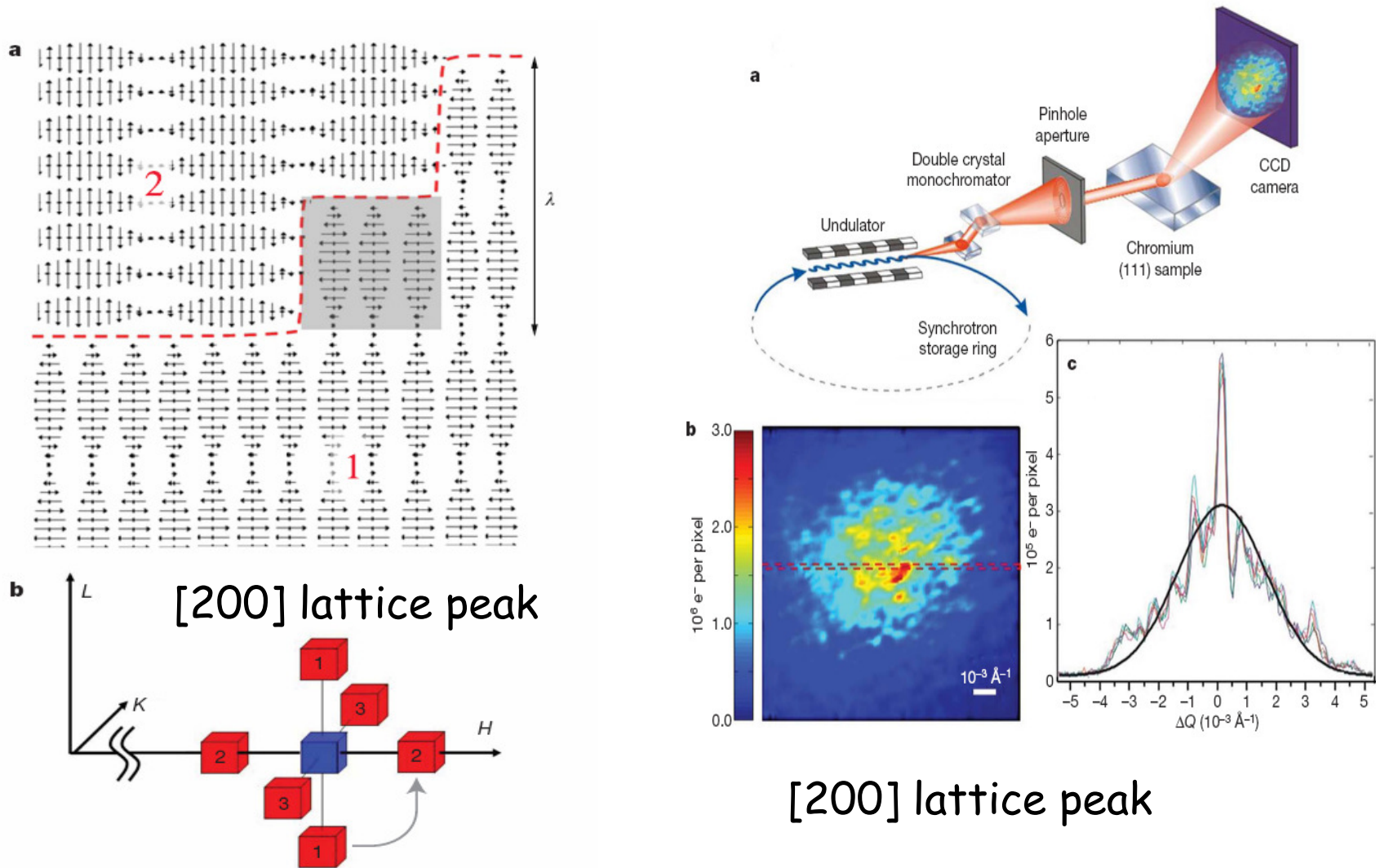
$$H(q) = \frac{D(q)}{D_0} S(q)$$

$$H(q) < 1$$

indirect hydrodynamic interactions slow down the dynamics

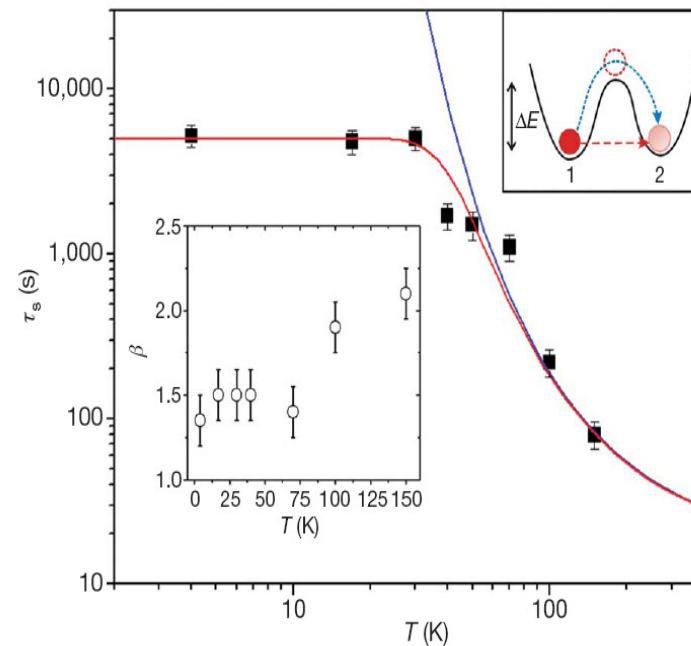
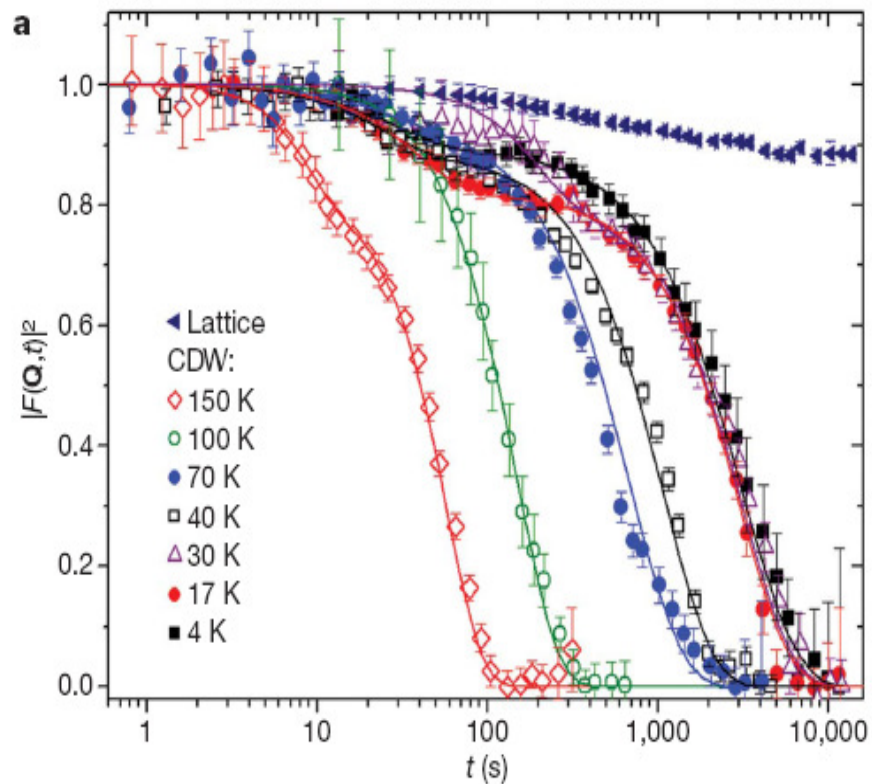


# Antiferromagnetic Domain fluctuations in Chromium

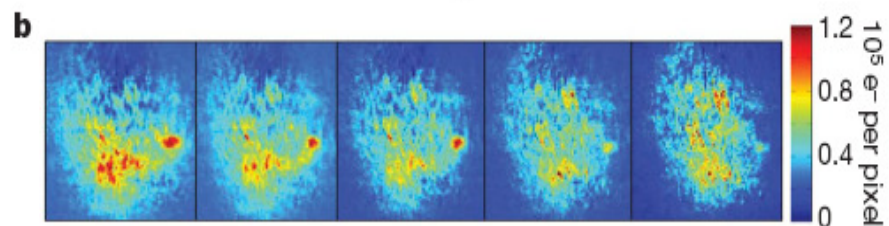


Charge density wave satellite peaks (CDW)

# Intensity Autocorrelation Function

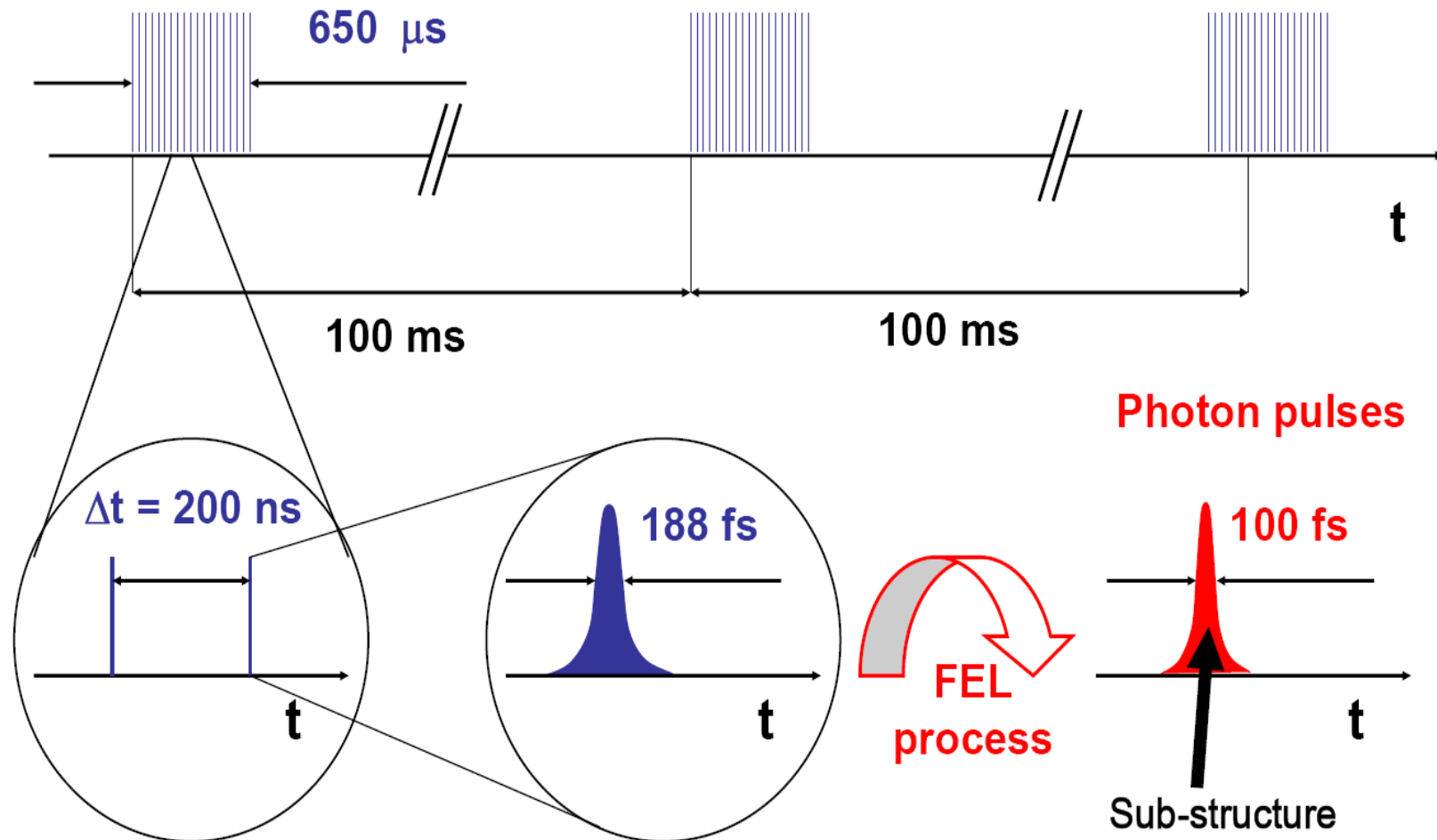


non Arrhenius behavior



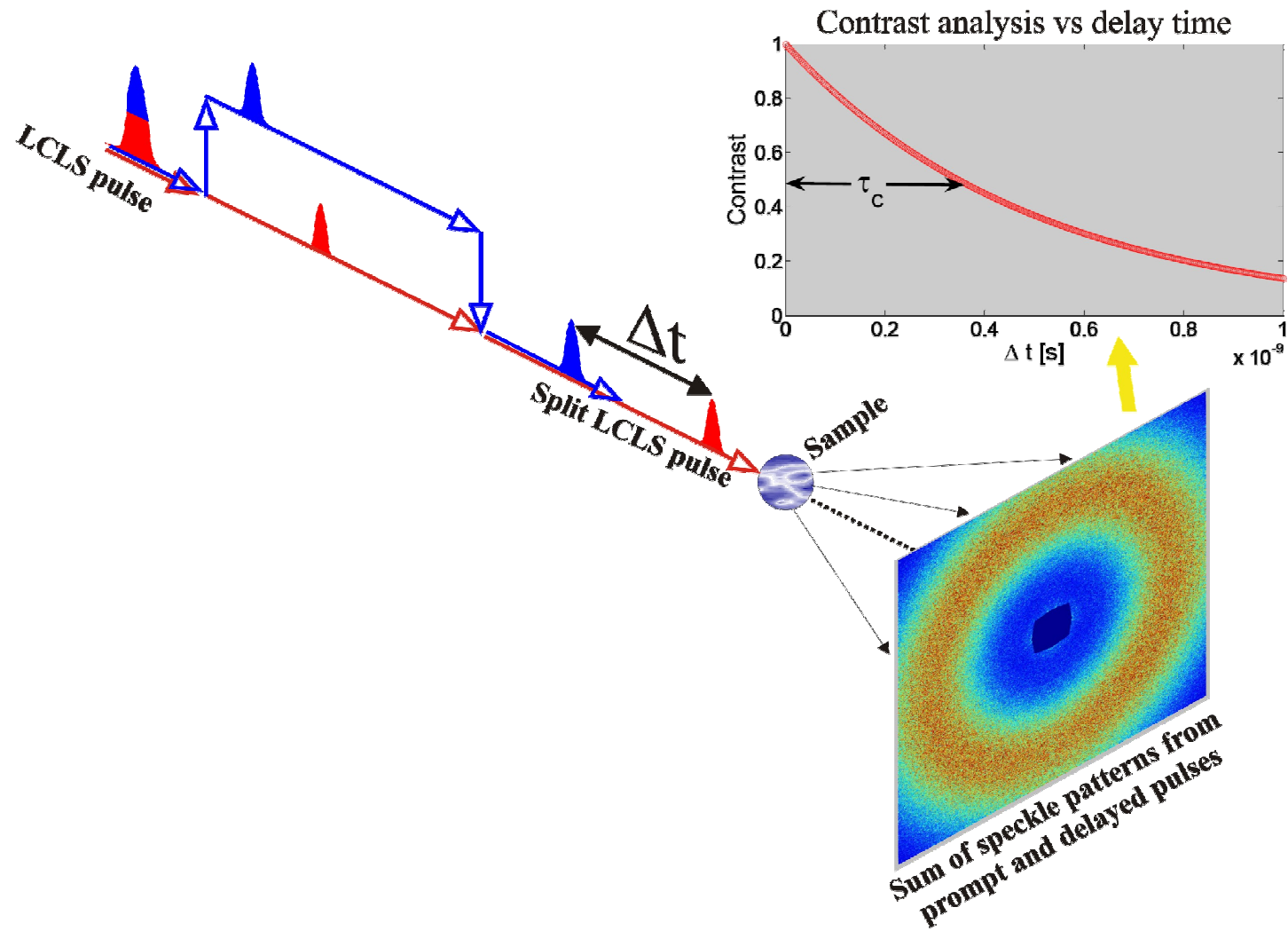
CDW satellite peak

## Electron bunch trains (with up to 3250 bunches à 1 nC)





# XPCS at XFEL



## Sequential technique

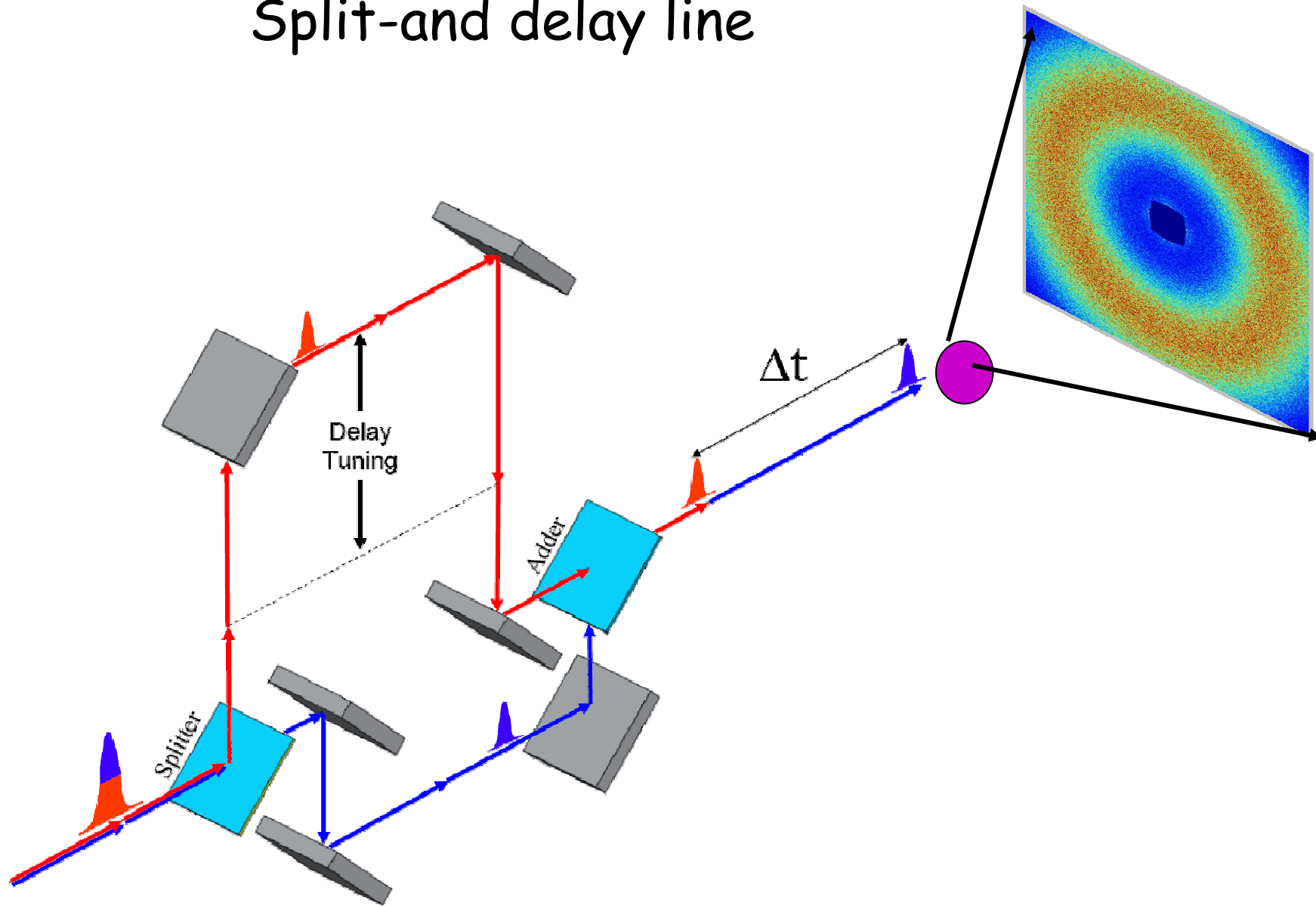
$$g_2(\tau) = \frac{\langle I(t)I(t+\tau) \rangle}{\langle I(t) \rangle^2} = 1 + \beta |f(\tau)|^2$$

## Split-Pulse technique

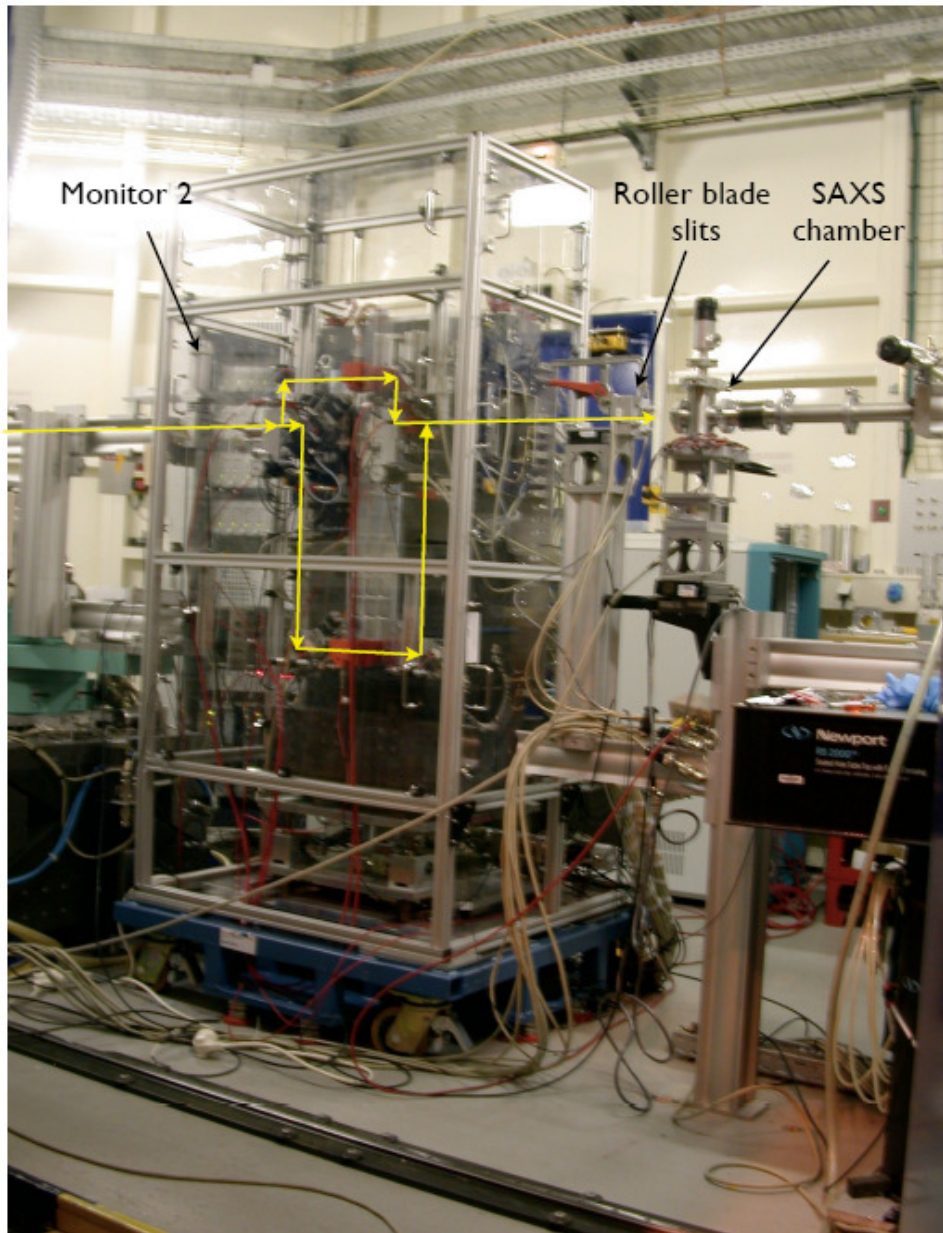
$$S(\tau) = I(t) + I(t+\tau)$$

$$c_2(\tau) = \frac{\langle S(\tau)^2 \rangle - \langle S(\tau) \rangle^2}{\langle S(\tau) \rangle^2} = \frac{\beta}{2} \left( 1 + |f(\tau)|^2 \right)$$

# Split-and delay line



# X-ray Split and Delay Unit

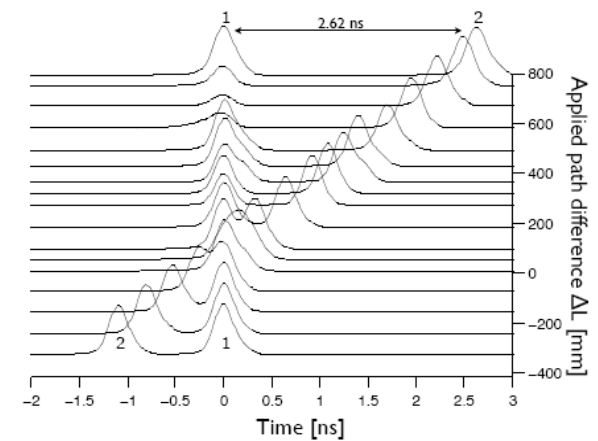
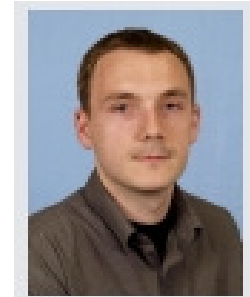


Development of an X-ray delay unit for  
correlation spectroscopy and pump - probe  
experiments

Dissertation zur Erlangung des Doktorgrades  
des Fachbereichs Physik der Universität  
Hamburg

vorgelegt von

Wojciech Roseker  
aus Bydgoszcz, Polen  
Hamburg  
2008



# First XFEL light in Stanford 2009



to come XFEL Japan, European XFEL DESY, PSI, Korea ?, China ?

The End