

# Methoden moderner Röntgenphysik I

## Coherence based techniques

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18. December 2008

# Outline

**18.12. 2008**

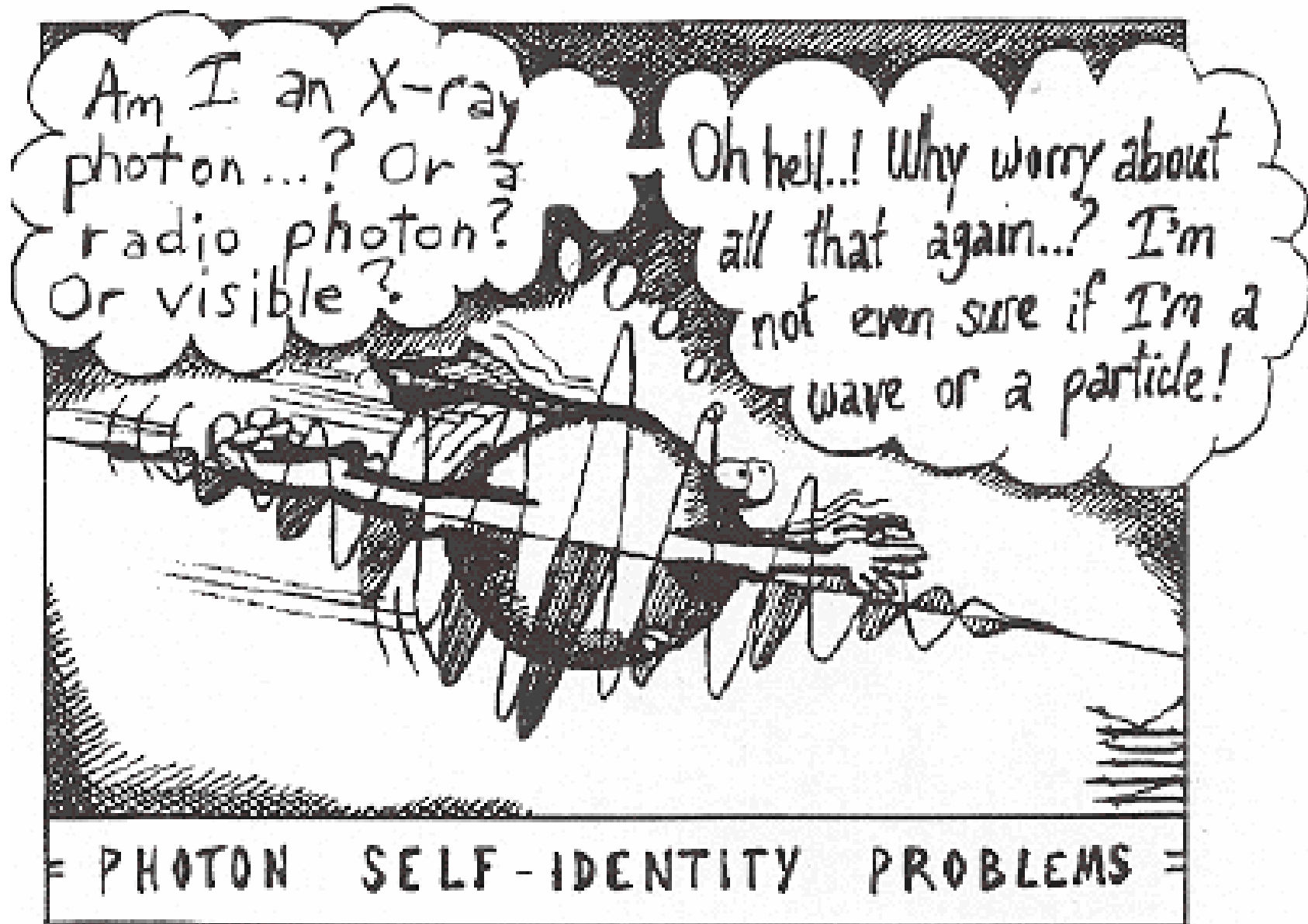
- **Introduction to Correlations**
- **Coherence Lengths**
- **van Zittert-Cernike Theorem**
- **Examples**
- **Speckle Pattern**

**8.01. 2009**

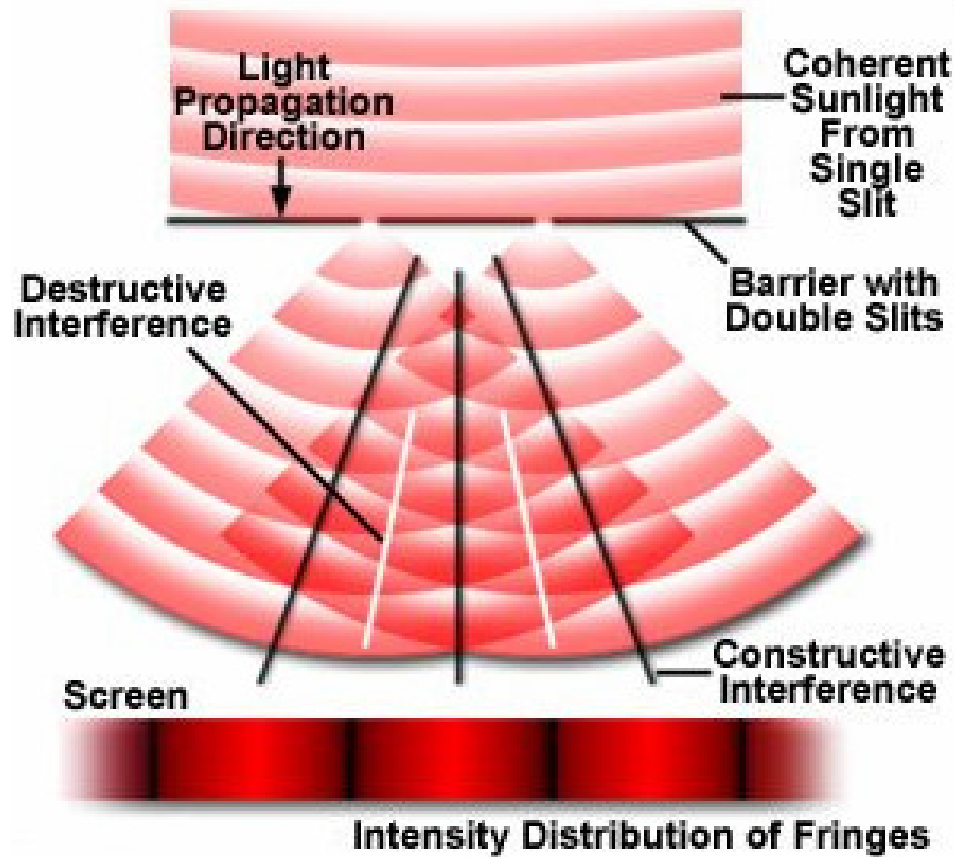
**Coherent Diffractive Imaging**

**15.01.2009**

**Correlation Spectroscopy**



# Young's Double Slit Experiment

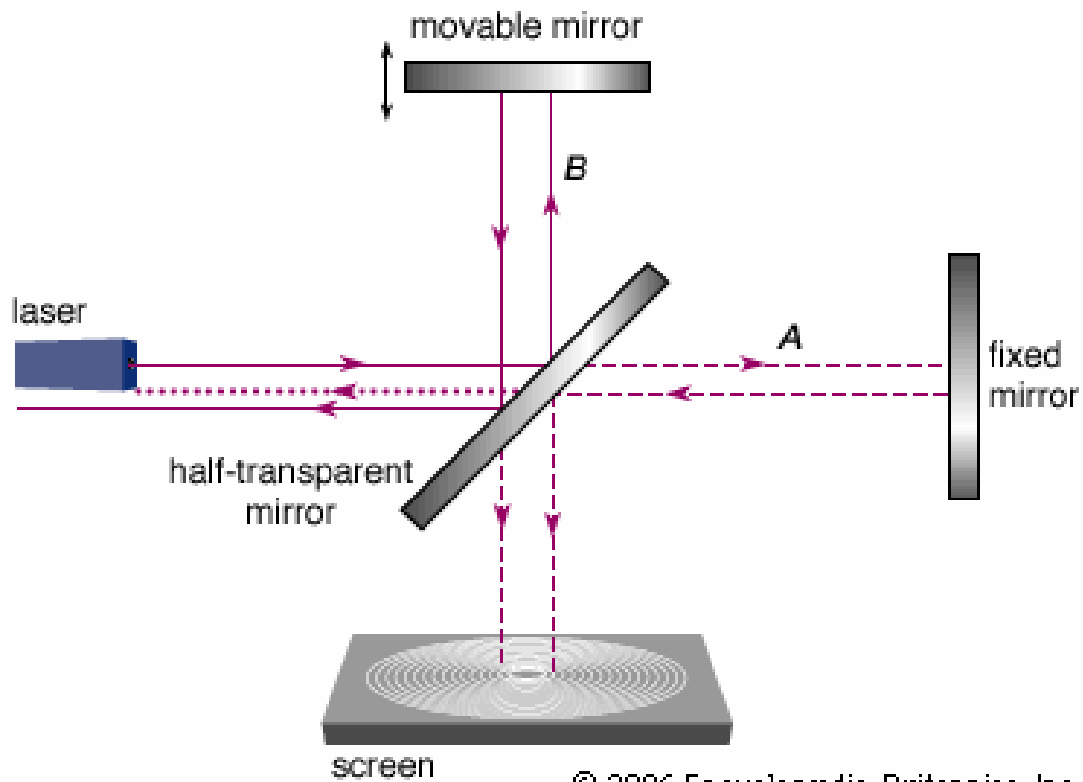


*Thomas Young*

Thomas Young, 1773-1829

Light is a wave

# Michelson Interferometer



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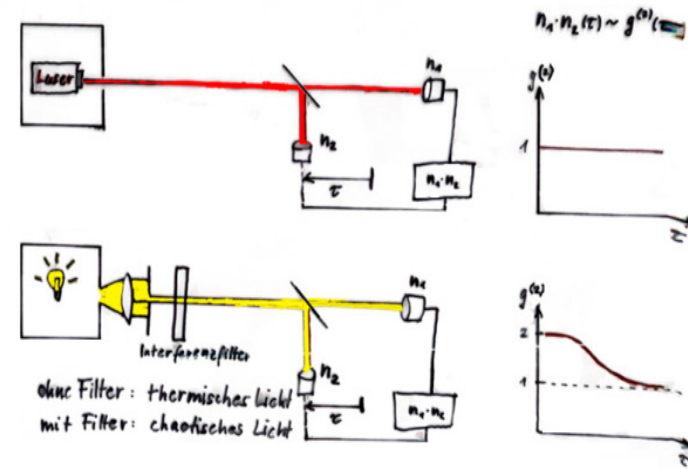
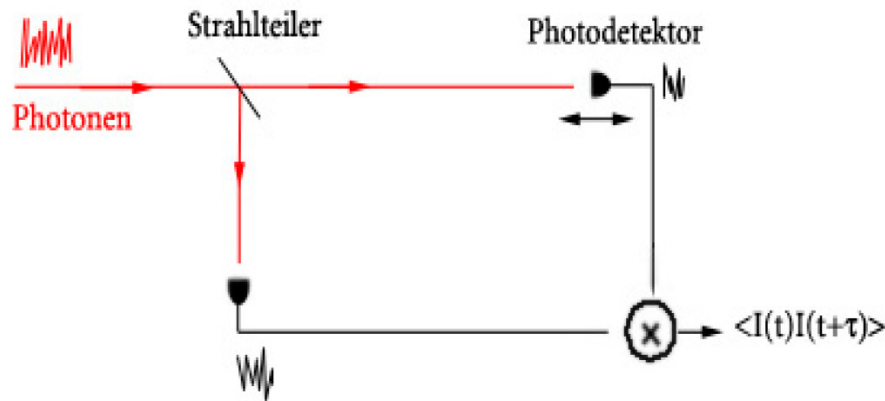
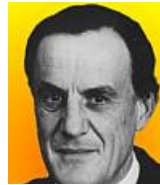
Albert Abraham Michelson  
1852-1931

no Äther  
-> Einstein

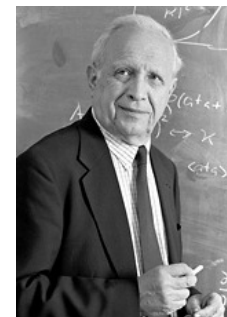
# Intensity Correlations

## 2nd order correlations

Robert Hanbury Brown  
Richard Q. Twiss

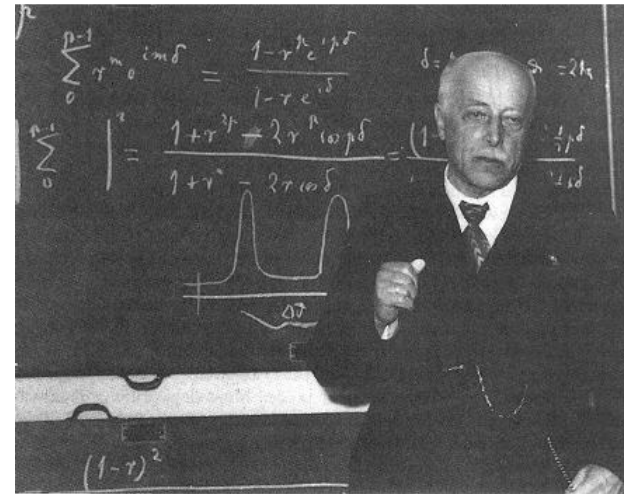
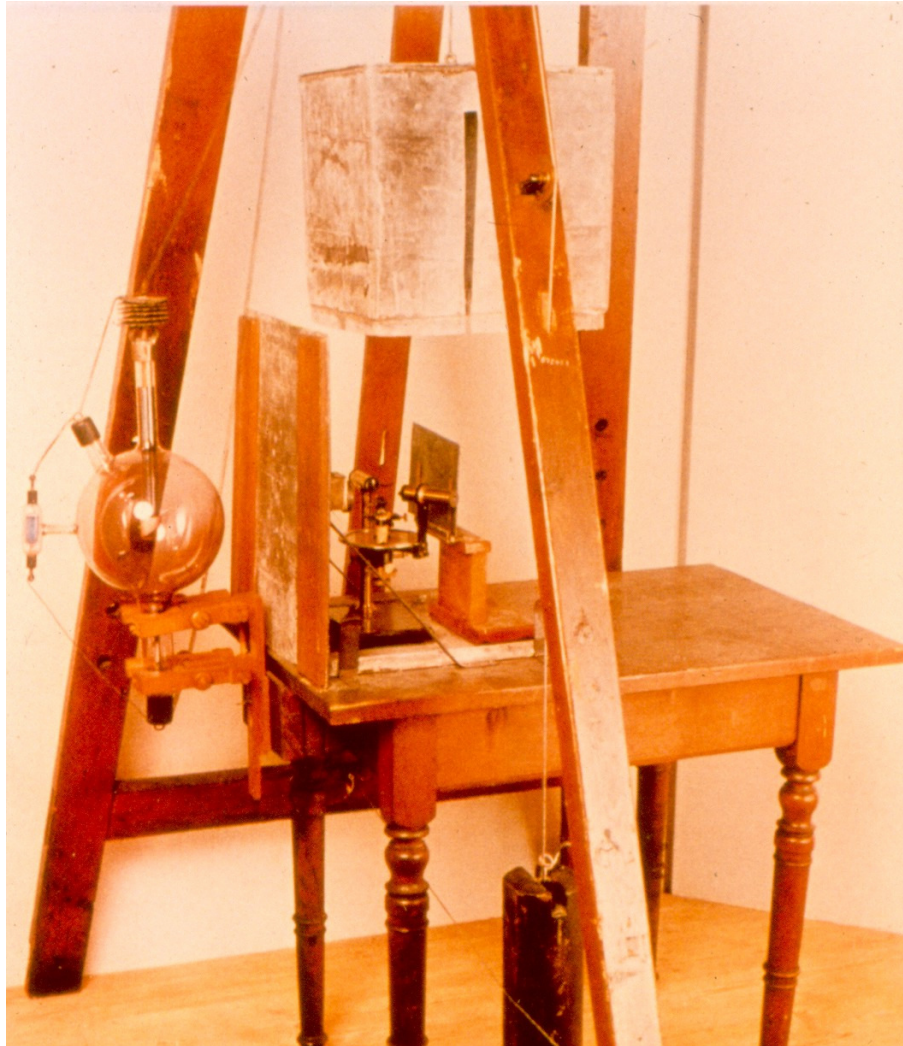


Roy Glauber- Nobel price in Physics 2005  
Theory of quantum coherence and optics

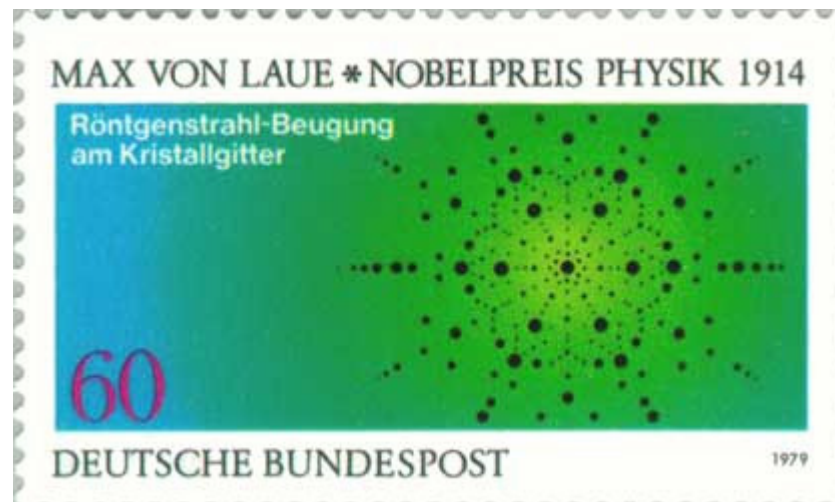


# Interference from a crystal lattice

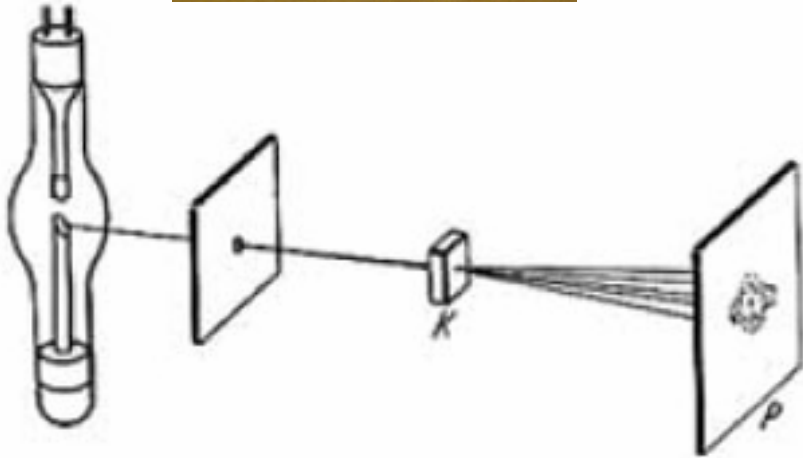
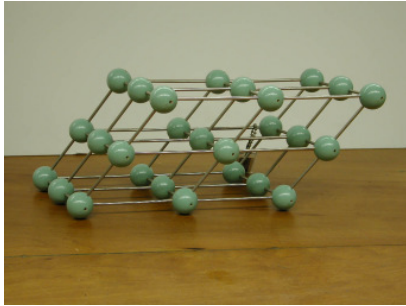
Laue, Friedrich and Knipping 1912



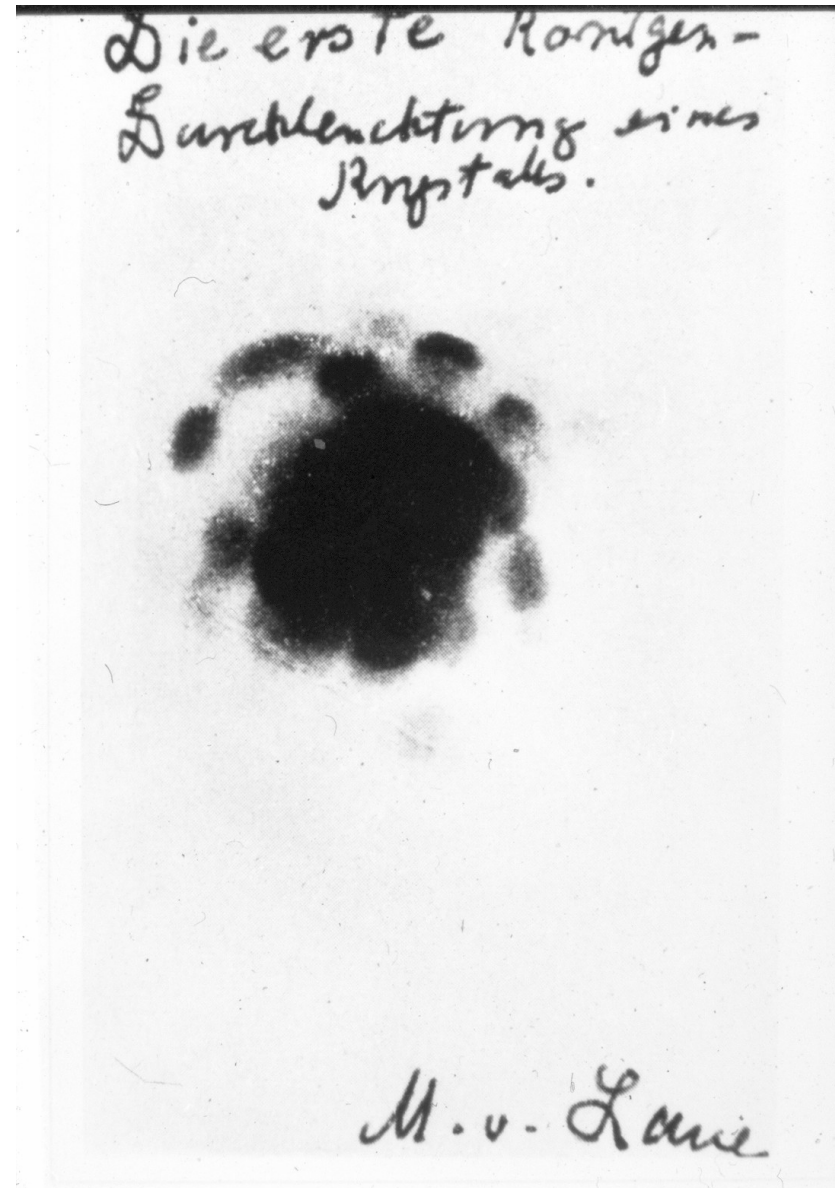
LAUE hielt seine Idee keineswegs geheim, sondern diskutierte sie sowohl mit SOMMERFELD als auch in der berühmten Runde junger Physiker, die sich täglich im Cafe Lutz traf. Aber weder SOMMERFELD noch die Runde waren überzeugt: Die Temperaturbewegung der Atome würde wohl die vermuteten Interferenzen hoffnungslos zerstören. Auch WIEN und MIE, mit denen sich SOMMERFELD zum Schilaulen traf, verwarfen die Idee. Nur LAUE insistierte weiter, und so erklärte sich schließlich WALTHER FRIEDRICH, der als Schüler RÖNTGENs SOMMERFELDs experimenteller Assistent wurde, zu dem Experiment bereit. Allerdings stimmte SOMMERFELD nicht zu. Erst als ein weiterer Doktorand RÖNTGENs, PAUL KNIPPING, seine Mithilfe anbot, wagte FRIEDRICH den Aufbau einer ersten, noch primitiven Apparatur. Das erste Ergebnis war negativ. Aber KNIPPING bestand auf einer anderen Aufstellung der Photoplatte, und nun zeigte die Durchstrahlungsaufnahme des Kupfersulfats einen Kranz abgebeugter Spektren um den direkten Strahl herum. LAUE schreibt dazu: "Tief in Gedanken ging ich durch die Leopoldstraße nach Haus, als mir FRIEDRICH die Aufnahme gezeigt hatte. Und schon nahe meiner Wohnung... kam mir der Gedanke für die mathematische Theorie der Erscheinung". Die dreifache Anwendung der einfachen Beugungsbedingung des eindimensionalen Gitters erklärte die neue Entdeckung.





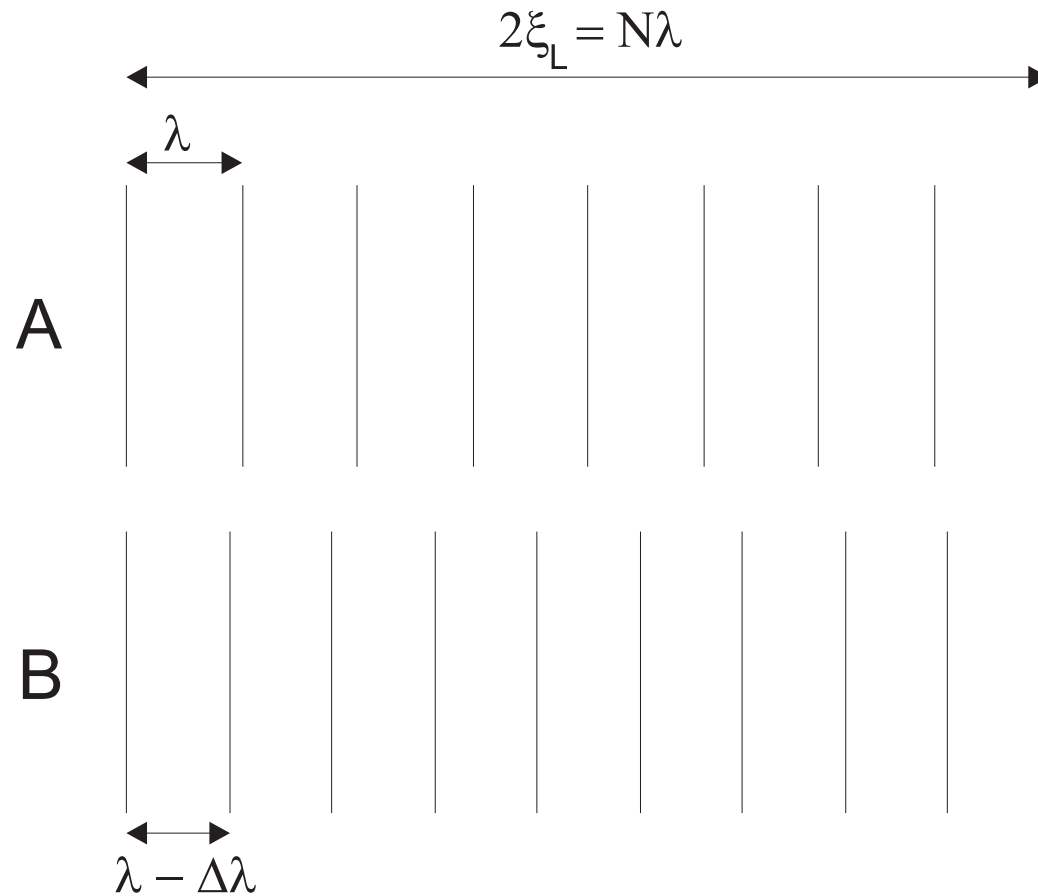


Kupfersulfat  
 $\text{CuSO}_4 \cdot 5 \text{H}_2\text{O}$



exposure time: many hours

# Longitudinal coherence



$$N\lambda = (N + 1)(\lambda - \Delta\lambda)$$

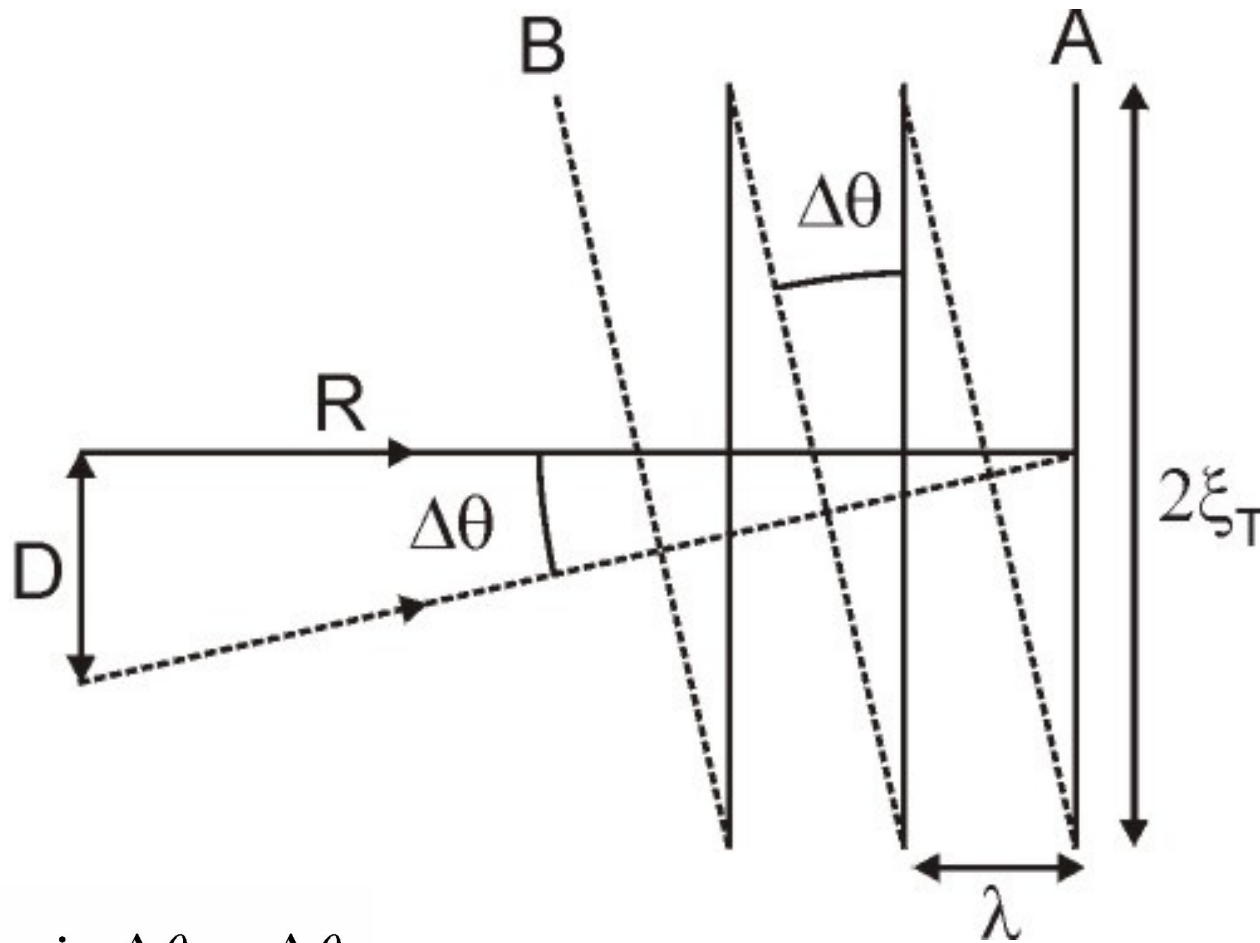
$$(N + 1)\Delta\lambda = \lambda$$

$$N \approx \frac{\lambda}{\Delta\lambda}$$

$$\xi_L = \frac{N\lambda}{2} = \frac{1}{2} \frac{\lambda^2}{\Delta\lambda}$$

longitudinal coherence depends  
on bandwidth

# Transverse coherence



$$\frac{\lambda}{2\xi_T} = \sin \Delta\theta \approx \Delta\theta$$

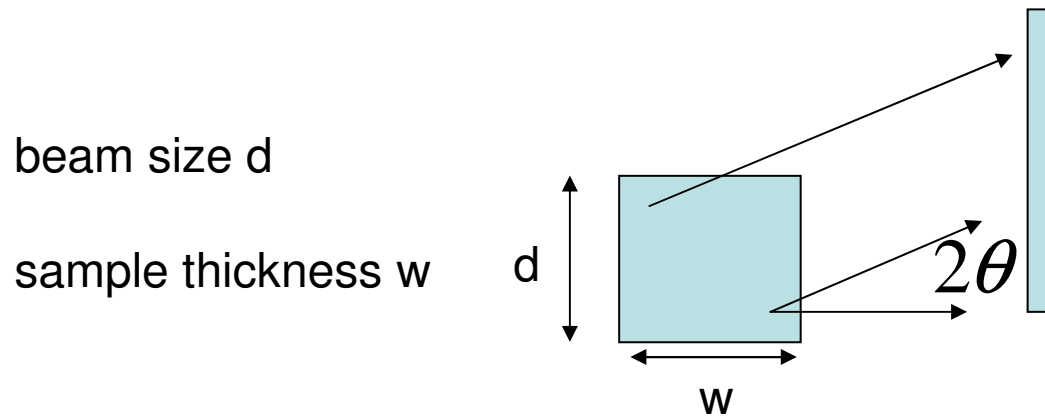
$$\frac{D}{R} = \tan \Delta\theta \approx \Delta\theta$$

$$\xi_T \approx \frac{\lambda R}{2 D}$$

transverse coherence depends  
on distance and source size

# Path length difference and coherence lengths

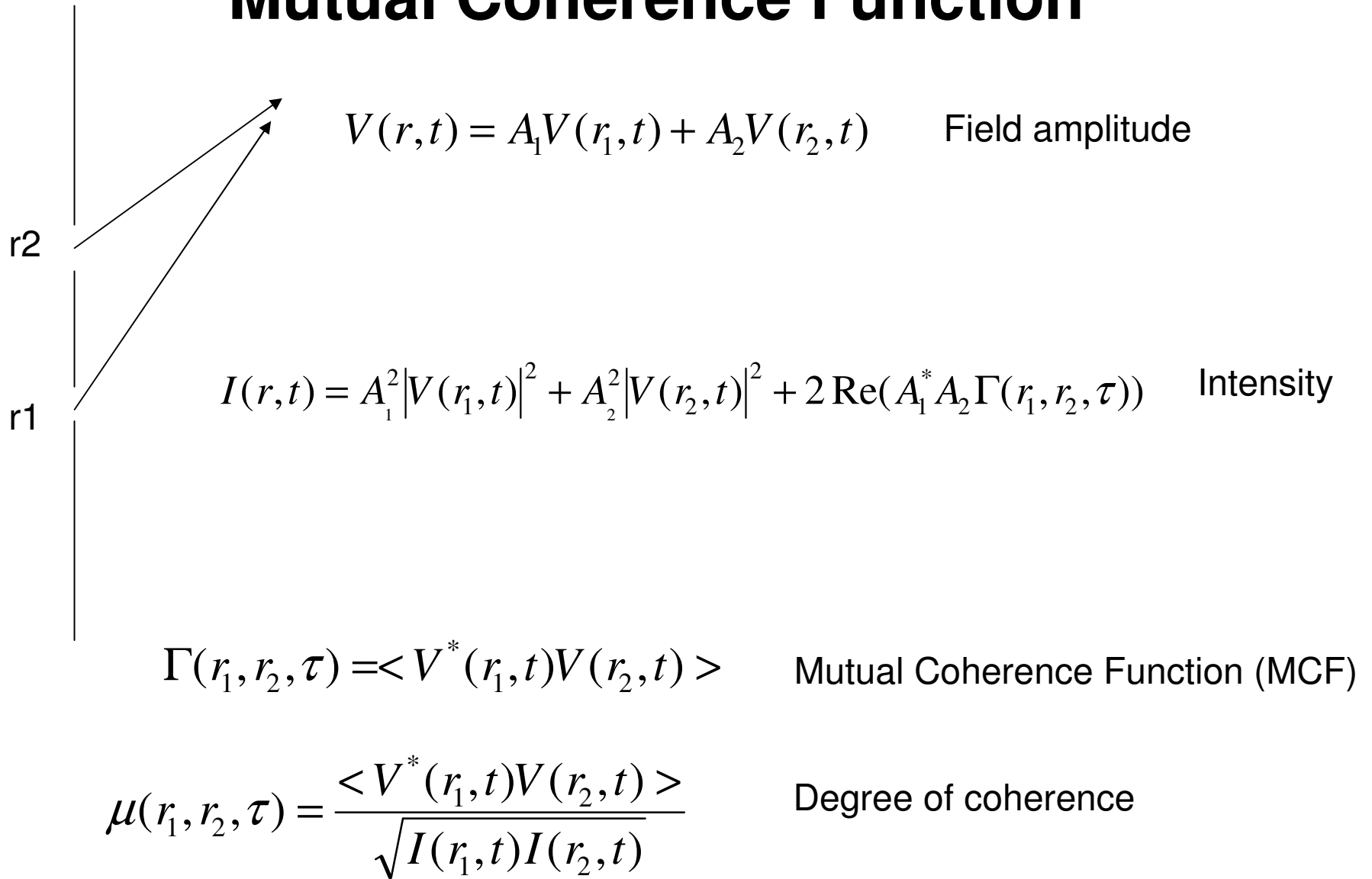
Interference is only possible when the path length differences (PLD) involved are of the same order as the coherence length.



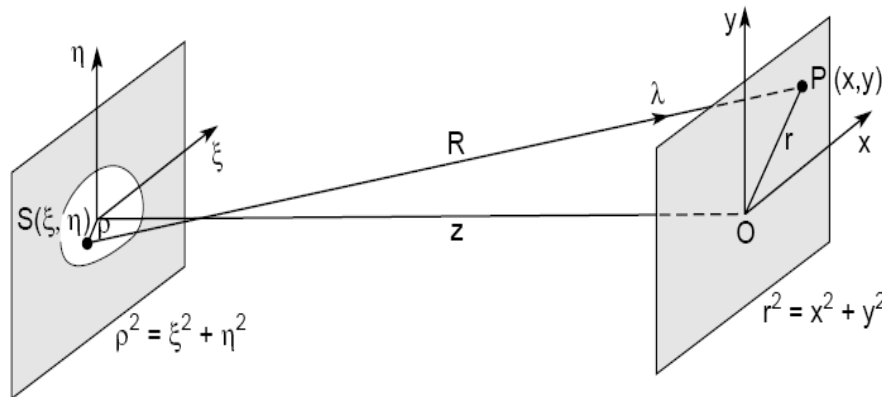
$$\text{PLD} \approx 2w \sin^2(\theta) + d \sin(2\theta)$$

for x-rays typically 1 micron

# Mutual Coherence Function



# van Cittert - Zernike Theorem



$$p = \frac{X}{R}, \quad q = \frac{Y}{R}, \quad \psi = \frac{k(X^2 + Y^2)}{R}$$

$$\theta = \frac{r}{z}$$

complex degree of coherence

$$\mu(0, P) = \frac{e^{-i\psi} \iint_S I(\xi, \eta) e^{ik(p\xi + q\eta)} d\xi d\eta}{\iint_S I(\xi, \eta) e^{ik(p\xi + q\eta)} d\xi d\eta}$$

Fourier Transform of the source intensity distribution!

Axial symmetry

$$\mu(0, P) = \frac{e^{i\psi} \int_0^{\infty} I(\rho) J_0(k\rho\theta) \rho d\rho}{\int_0^{\infty} I(\rho) \rho d\rho}$$

## Applications van Cittert - Zernike Theorem (1)

point source

$$I(\rho) = \frac{I_0 \delta(\rho)}{2\pi\rho}$$

$$\mu(0, P) = \frac{e^{-i\psi} \int_0^{\infty} \delta(\rho) J_0(k\rho\theta) d\rho}{\int_0^{\infty} \delta(\rho) d\rho} = e^{-i\psi} J_0(0)$$

$$|\mu(0, P)| = 1$$

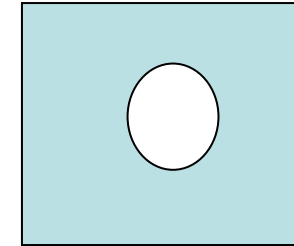
For a point source the normalized absolute degree of coherence is 1, i.e. it is a fully coherent source!

# Applications van Cittert - Zernike Theorem (2)

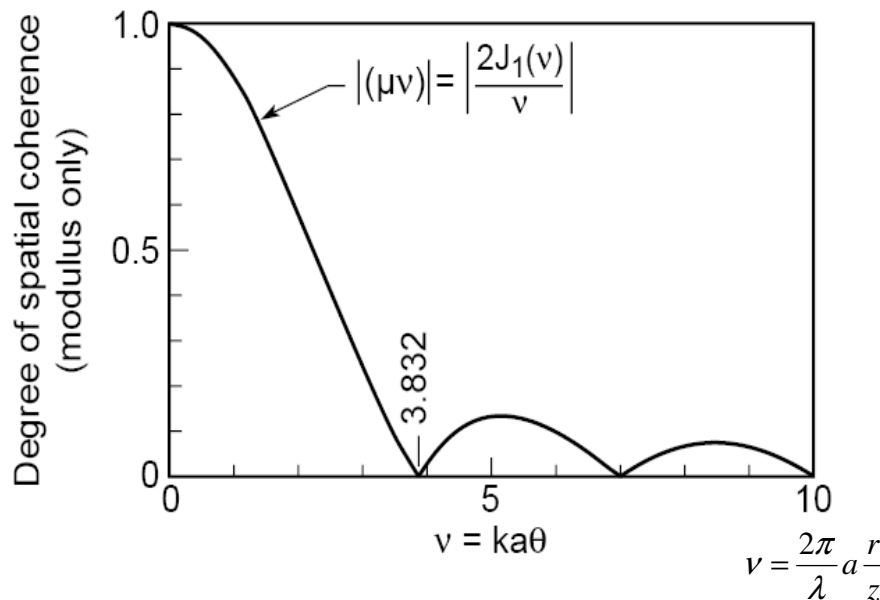
pinhole diameter  $2a$

$$I(\rho) = 1 \text{ for } \rho \leq a$$

$$I(\rho) = 0 \text{ for } \rho > a$$



$$\mu(0, P) = \frac{e^{-i\psi} \int_0^a J_0(k\rho\theta) \rho d\rho}{\int_0^a \rho d\rho} = \frac{8e^{-i\psi}}{d^2} \int_0^a J_0(k\rho\theta) \rho d\rho = e^{-i\psi} \frac{2J_1(ka\theta)}{ka\theta}$$



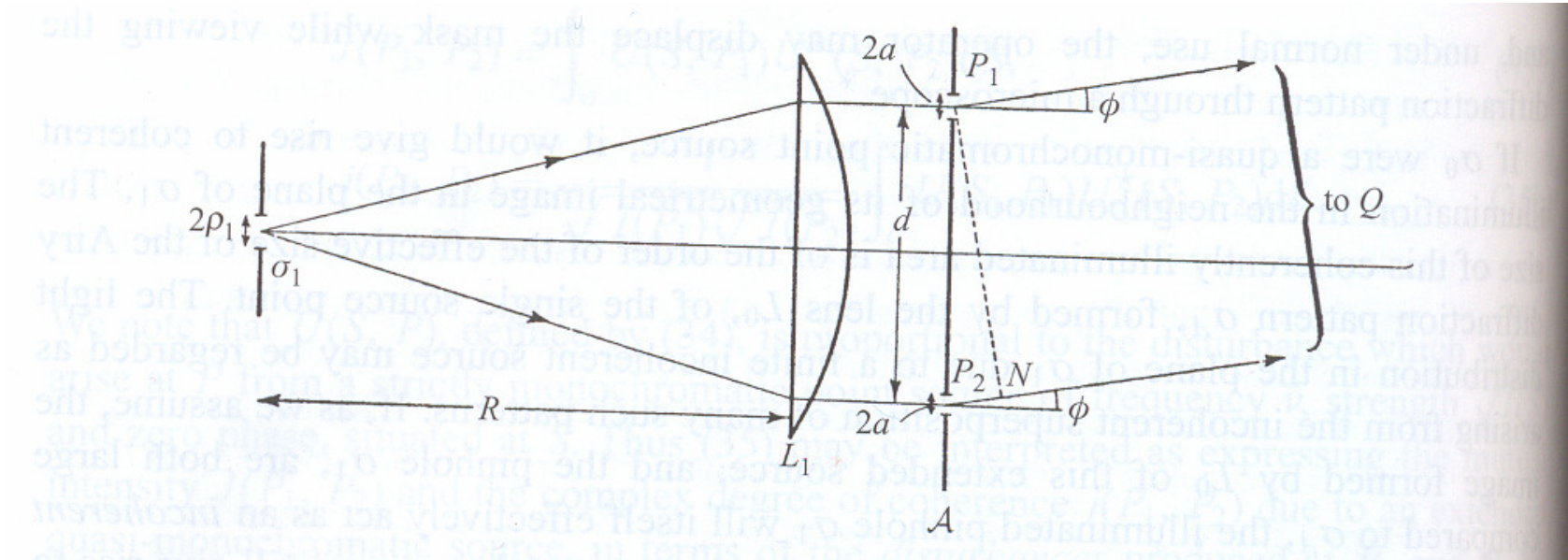
$$|\mu(0, P)| = 0.88$$

$$r = 0.16 \frac{z\lambda}{a}$$

Sun:  $r=19$  microns  
 x-rays at synchrotron  
 $r=10-100$  microns



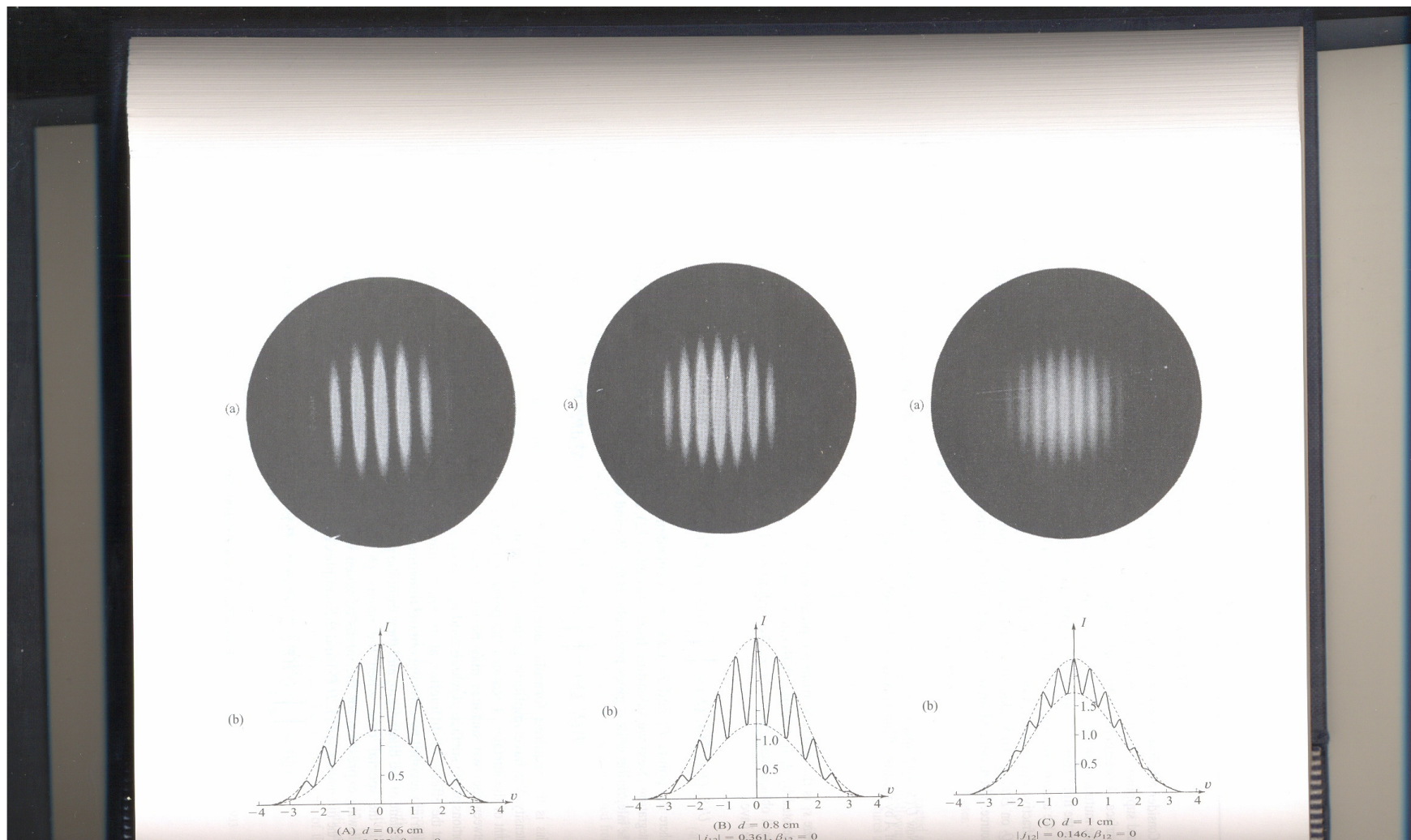
# Revisiting Young's Double-Slit experiment



$$I(\phi, d) = 2 \left( \frac{2J_1(u)}{u} \right)^2 \left( 1 + \left| \frac{2J_1(v)}{v} \right| \cos(\beta(v) - Cuv) \right)$$

$$u = ka \sin(\phi), \quad v = \frac{k\rho_1 d}{R}, \quad C = k \frac{R}{\rho_1 a} \quad \begin{array}{l} \beta = 0 \text{ when } J_1(v) > 0 \\ \beta = \pi \text{ else} \end{array}$$

# Measuring the Degree of Coherence (1)



Born and Wolf, Optics

# Measuring the Degree of Coherence (1)

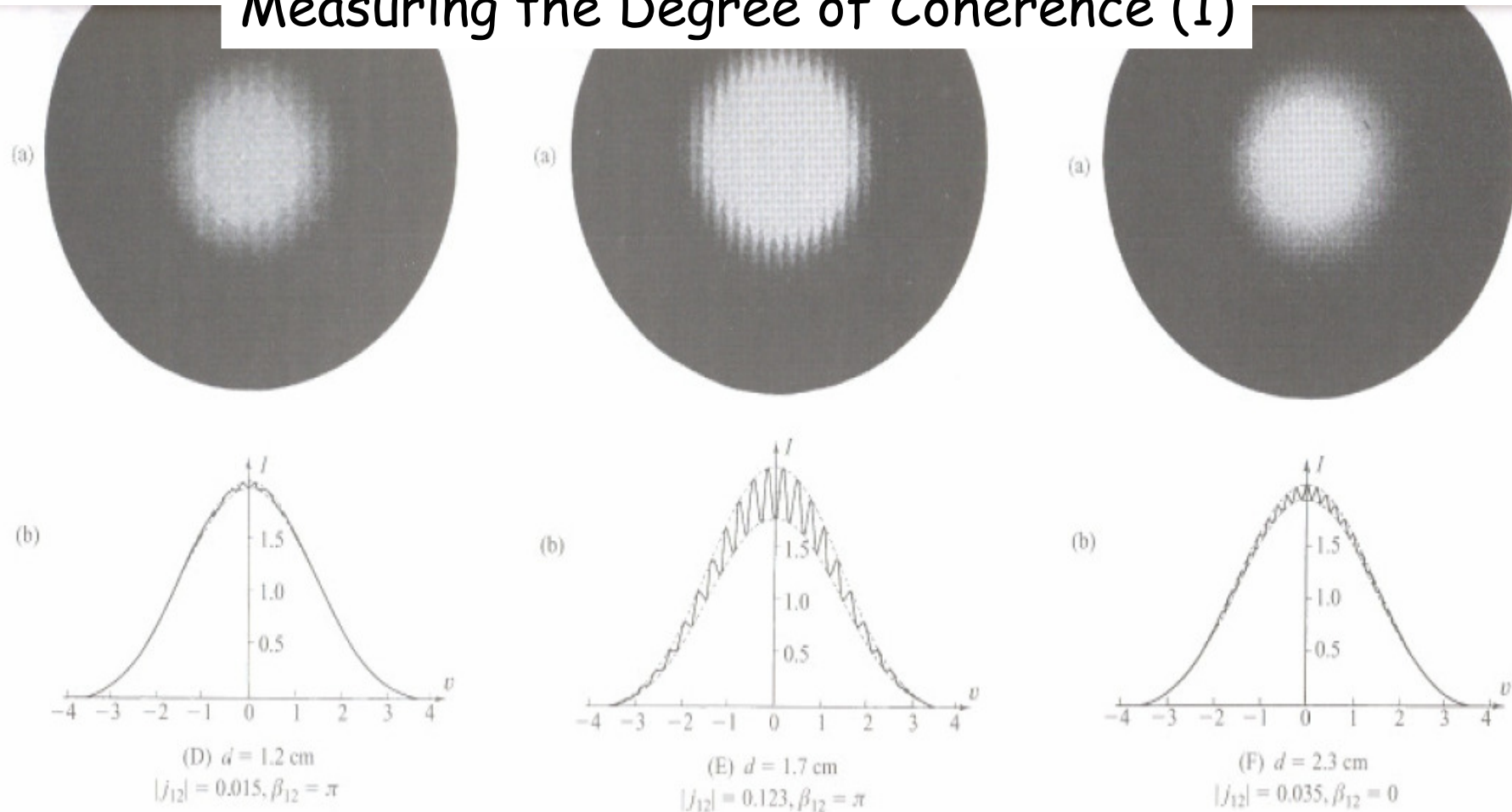
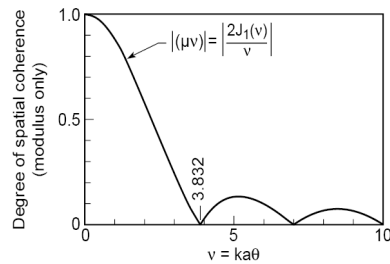


Fig. 10.6 Two-beam interference with partially coherent light. (a) Observed patterns. (b) theoretical intensity curves. Focal length of lenses  $f_{\Delta}$ .



Born and Wolf, Optics

# Coherence

Measure of coherence : Degeneracy parameter  
 - number of photons per mode

$\Delta_c$

$$\Delta_c := \text{Br} \cdot \tau_0 \cdot \delta\Omega \cdot \frac{\Delta\nu}{\nu} = \text{Br} \cdot \frac{\lambda^3}{\pi c}$$

**Br.: Brilliance** [photons/(s · mm<sup>2</sup> · mrad<sup>2</sup> · 10<sup>-3</sup> BW)]

Quelle	Energie	Leistung	Br.	$\Delta_c$
Quecksilber-Lampe	4.9 eV	1 W	$2 \cdot 10^{11}$	$3 \cdot 10^{-3}$
ESRF high $\beta$ Undulator	6.4 keV	einige kW	$2 \cdot 10^{20}$	$2 \cdot 10^{-3}$
HeNe Laser	1.96 eV	1 mW	$8 \cdot 10^{19}$	$2 \cdot 10^7$
ILL neutrons	25 meV		$5 \cdot 10^2$	$3 \cdot 10^{-15}$

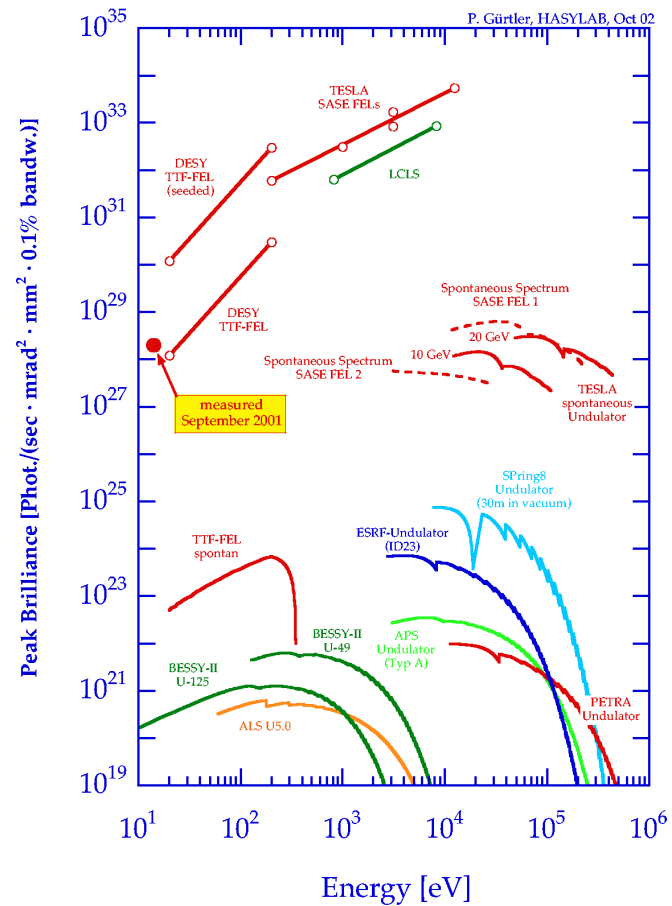
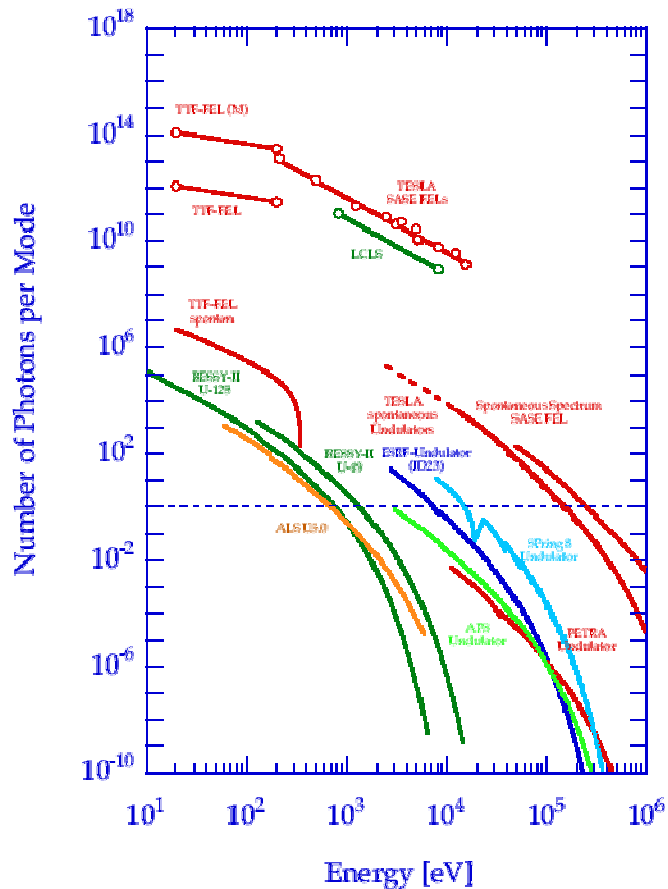
# Brilliance of X-rays Sources

Number of photons per mode

$$N_C = B \lambda^3 / (16\pi c)$$

Coherent Flux:  $F_0 = B \lambda^2 (\Delta\lambda/\lambda)$

(ESRF: ID10A  $F_0 \sim 10^{10}$  ph/s)



# Coherent X-rays

third generation synchrotron sources

( ESRF, APS, Spring-8,... )

(and soon : Diamond, Soleil, SLS, Petra-III)



Transverse coherence lengths  $\xi_T$

$$\xi_T = \frac{\lambda}{2} \cdot \frac{R}{s} = 155(V) \times 3.8(H) \mu\text{m}^2$$

Troika I Beamline parameters

$\lambda = 1.55\text{\AA}$ ,  $R=46\text{m}$ ,  
source size :  $23(V) \times 928(H)\mu\text{m}$

Longitudinal coherence length  $\xi_L$

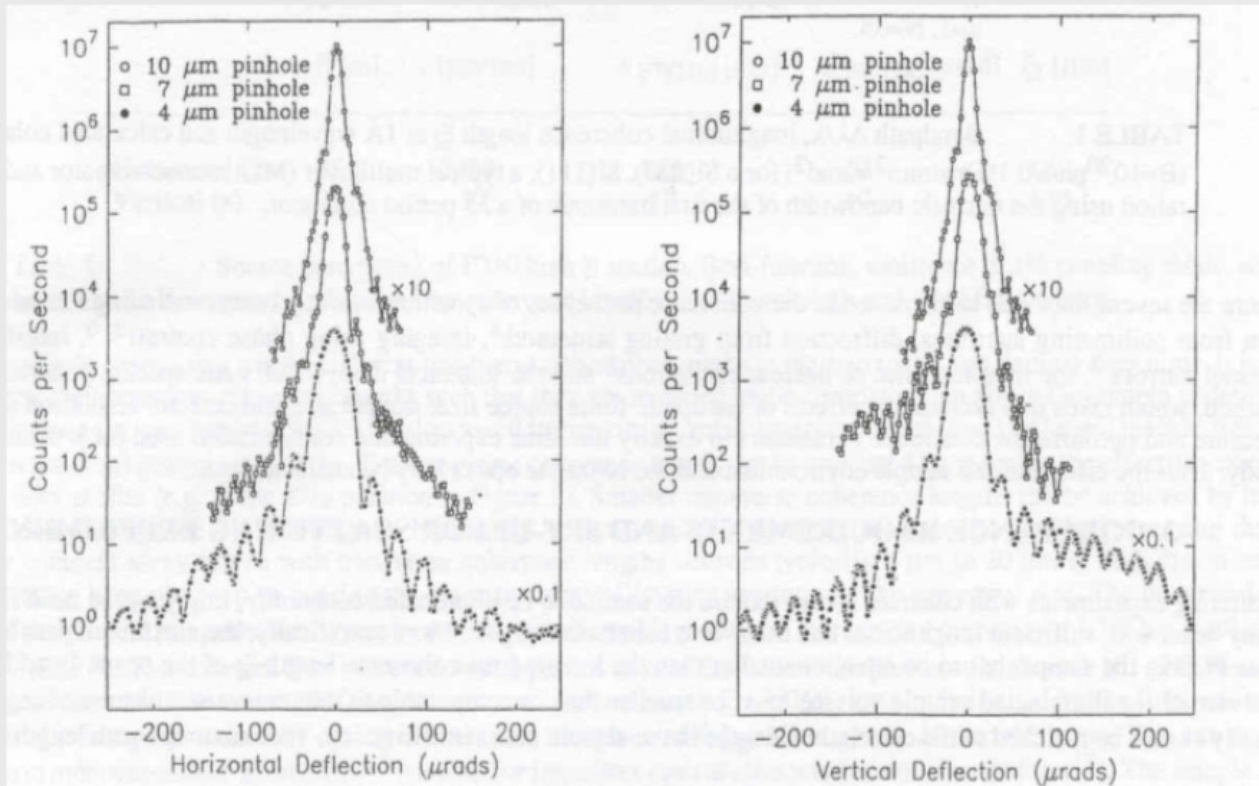
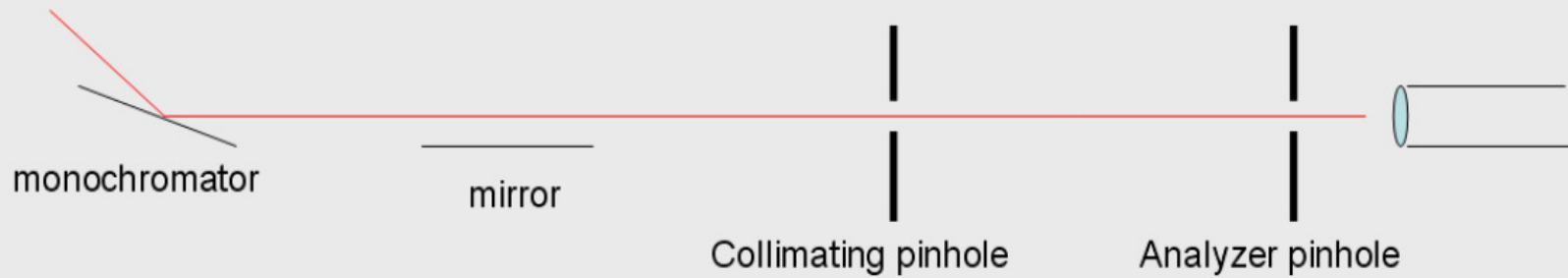
$$\xi_L = \lambda \cdot \left( \frac{\lambda}{\Delta\lambda} \right) \approx 1\mu\text{m}$$

$\lambda = 1.55\text{\AA}$ , Si(111)  $\Delta\lambda/\lambda = 1.4 \cdot 10^{-4}$

Selectioning coherent  
part of the beam

beam size  $\approx \xi_T \times \xi_T$

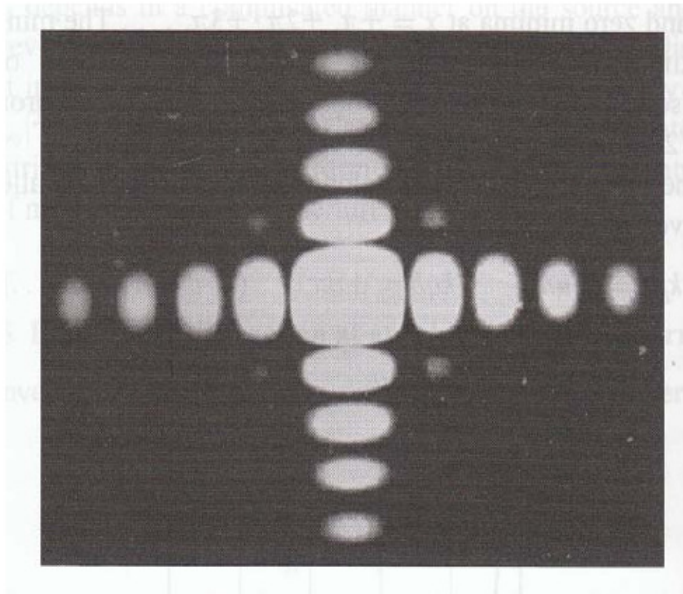
# Fraunhofer Diffraction ( $\lambda=0.1\text{nm}$ )



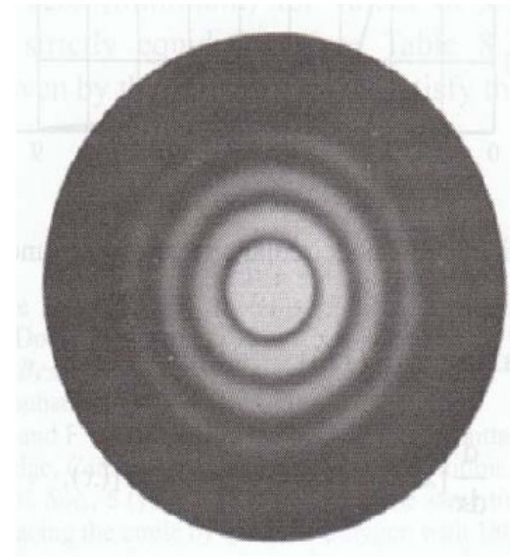
ID10 (ESRF)  
E=11.8 keV;  
 $\lambda=1.05\text{\AA}$

# Fraunhofer Diffraction

coherent illumination of apertures



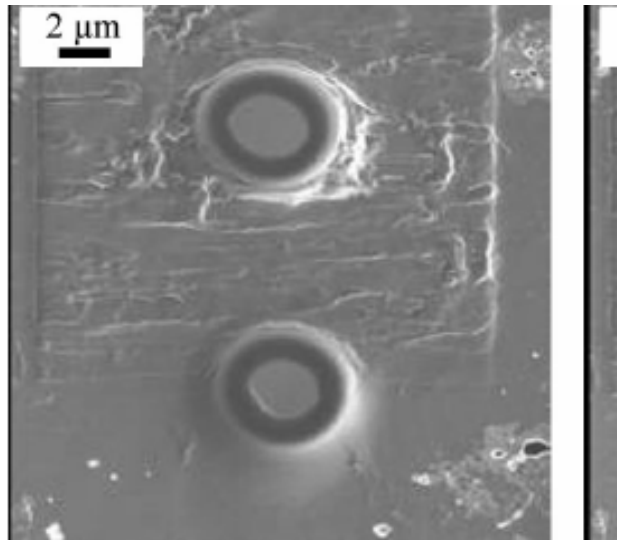
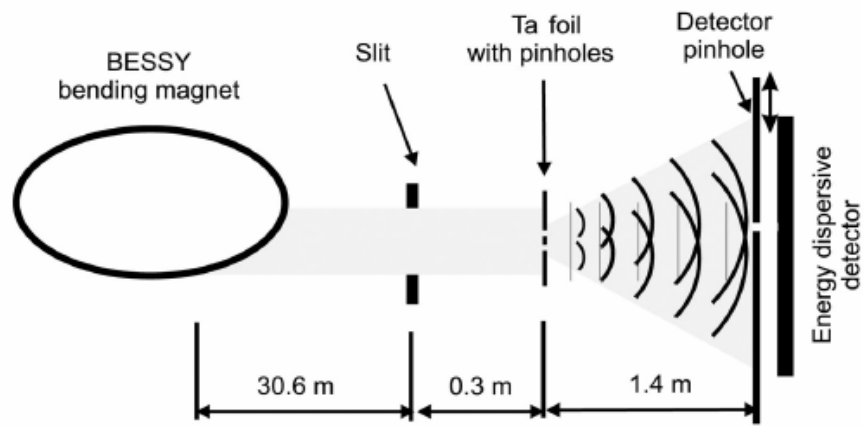
Fraunhofer diffraction of a rectangular aperture  $8 \times 7 \text{ mm}^2$ , taken with mercury light  $\lambda=579\text{nm}$  (from Born&Wolf, chap. 8)



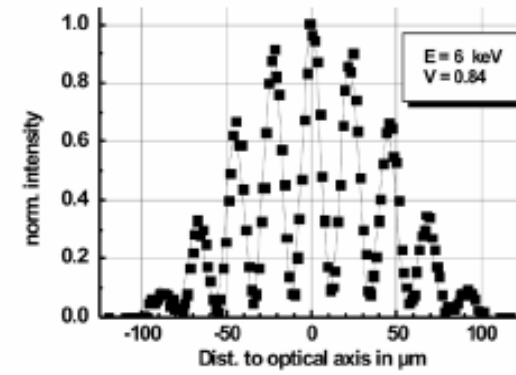
Fraunhofer diffraction of a circular aperture, taken with mercury light  $\lambda=579\text{nm}$  (from Born&Wolf, chap. 8)



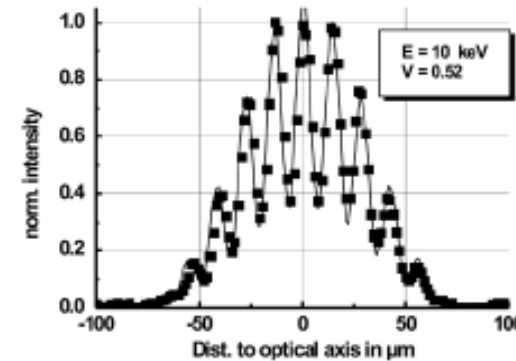
# Young's experiment with X-rays



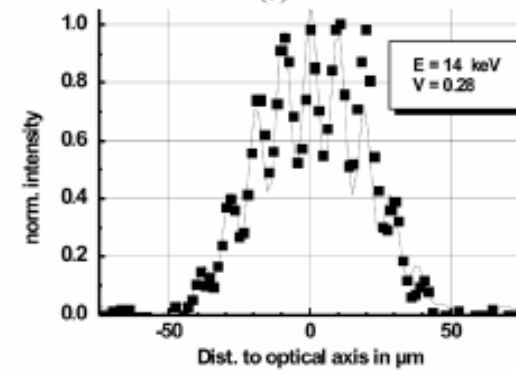
(a)



(a)

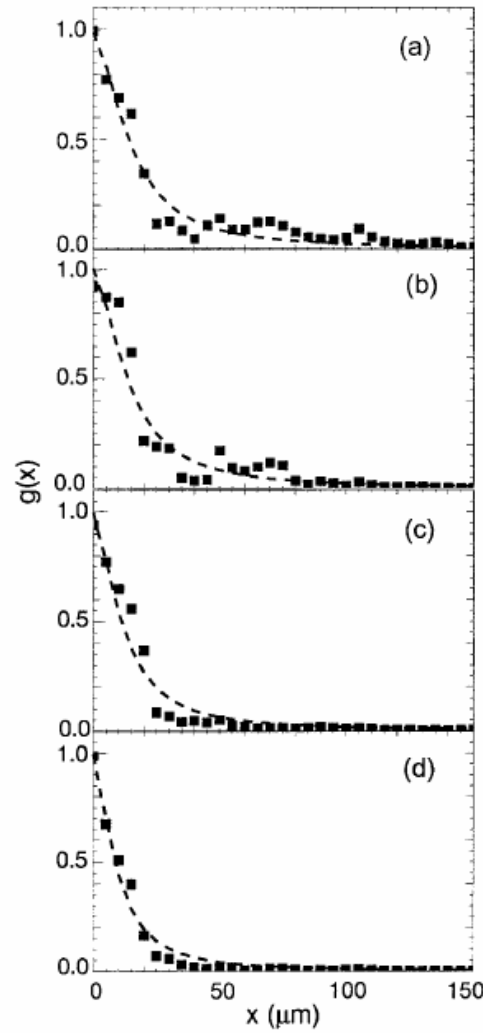


(b)



(c)

# Measuring the absolute degree of coherence for x-rays



10  $\mu\text{m}$  pinhole

50  $\mu\text{m}$  pinhole

90  $\mu\text{m}$  pinhole

100  $\mu\text{m}$  pinhole

# Speckle Pattern

"Everything interferes with everything"

# Coherent scattering

perfect coherence

$$I(q) = \left| \int_{-A}^A \rho(r) \exp(iqr) dr \right|^2$$

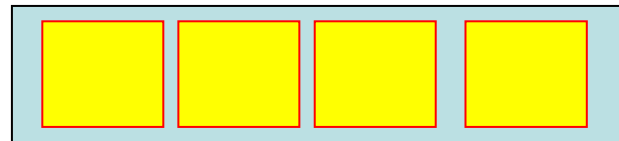
- absolute phase is lost
- relative phases are still in  $I(q)$

A sample size

partial coherence

$$I(q) = \sum_i \left| \int_{-L(i)}^{L(i)} \rho(r) \exp(iqr) dr \right|^2$$

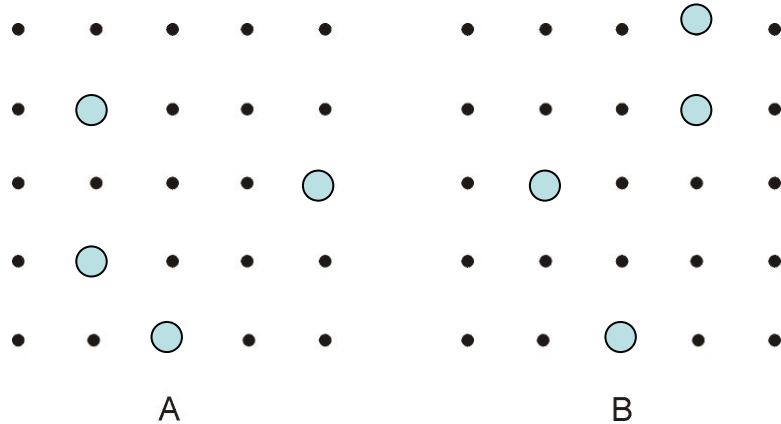
- absolute phase is lost
- relative phases are lost to a large extent



L(1)

L(2)

## Simulation of coherence scattering from a lattice:



- formfactor  $f_1$
- formfactor  $f_2 = 1.5 * f_1$

simulation parameter:

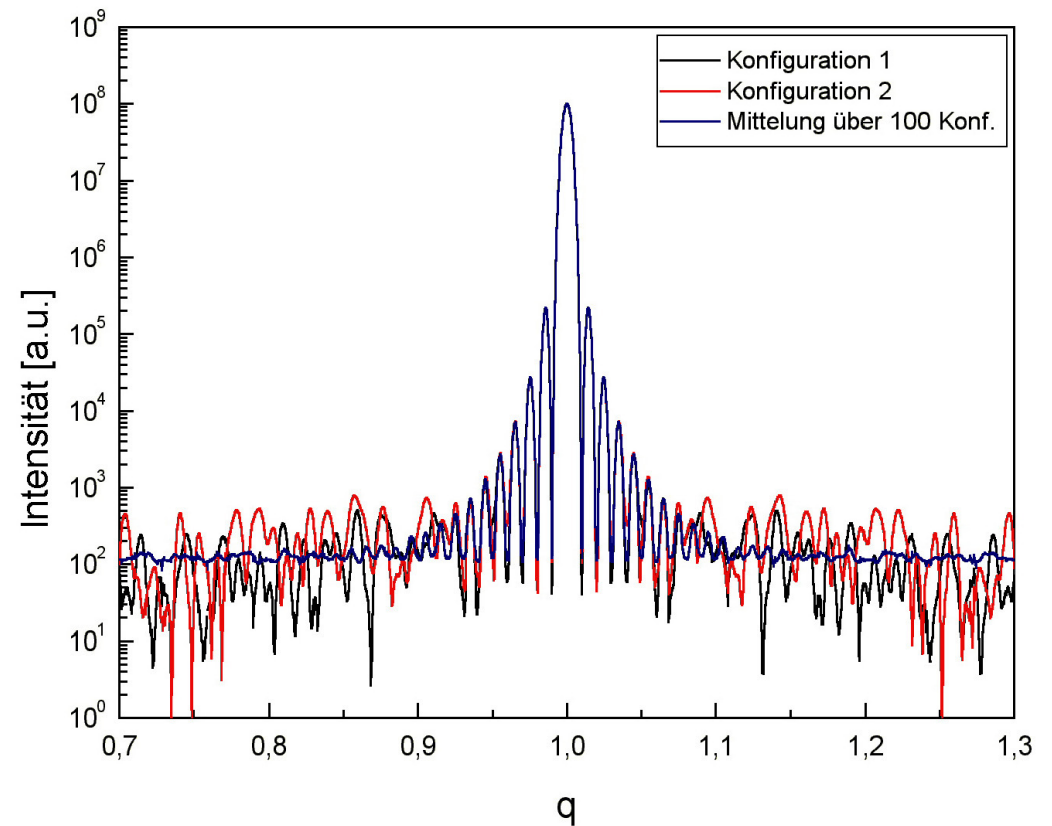
lattice 100x100

99 % : ●

1 % : ○

incoherence scattering:

averaging over 100 configurations

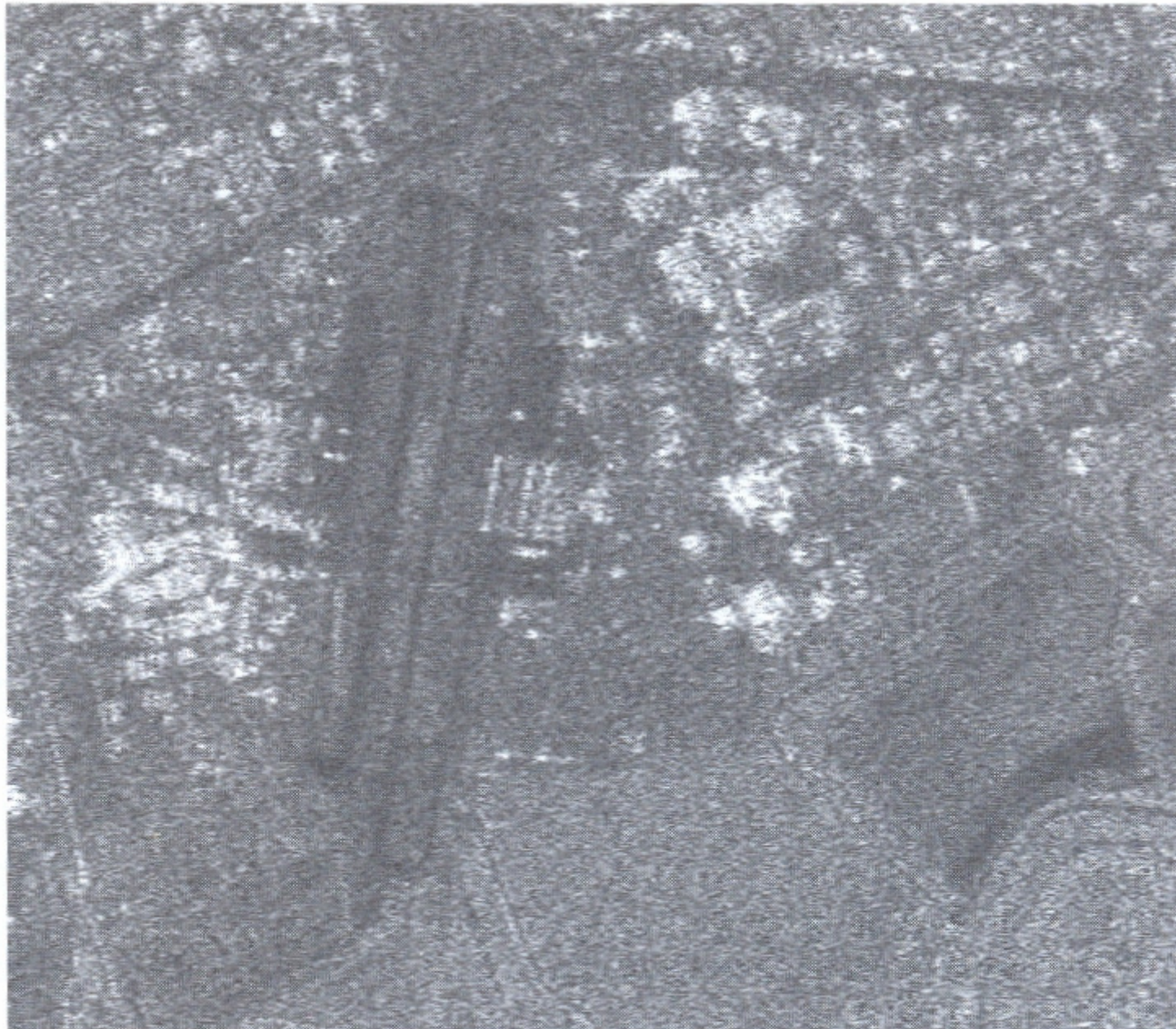


Moffet field, San Francisco, CA as shown on Google Earth

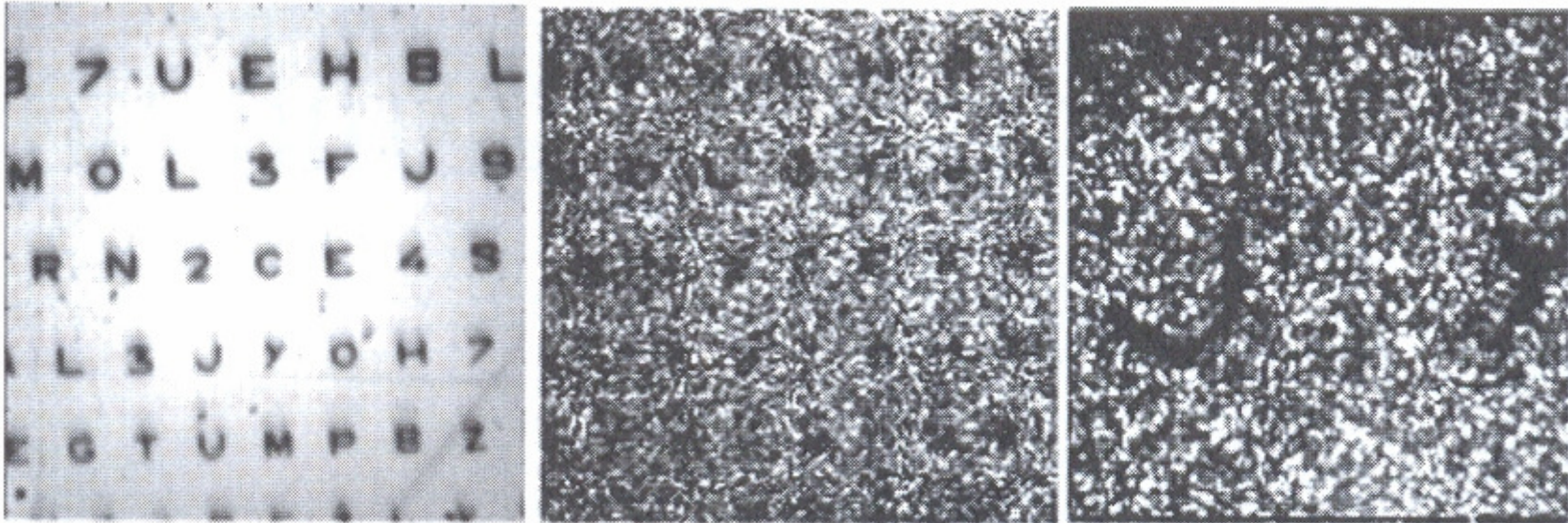


Moffet field, San Francisco, CA

## Radar Speckles



# Optical Speckles



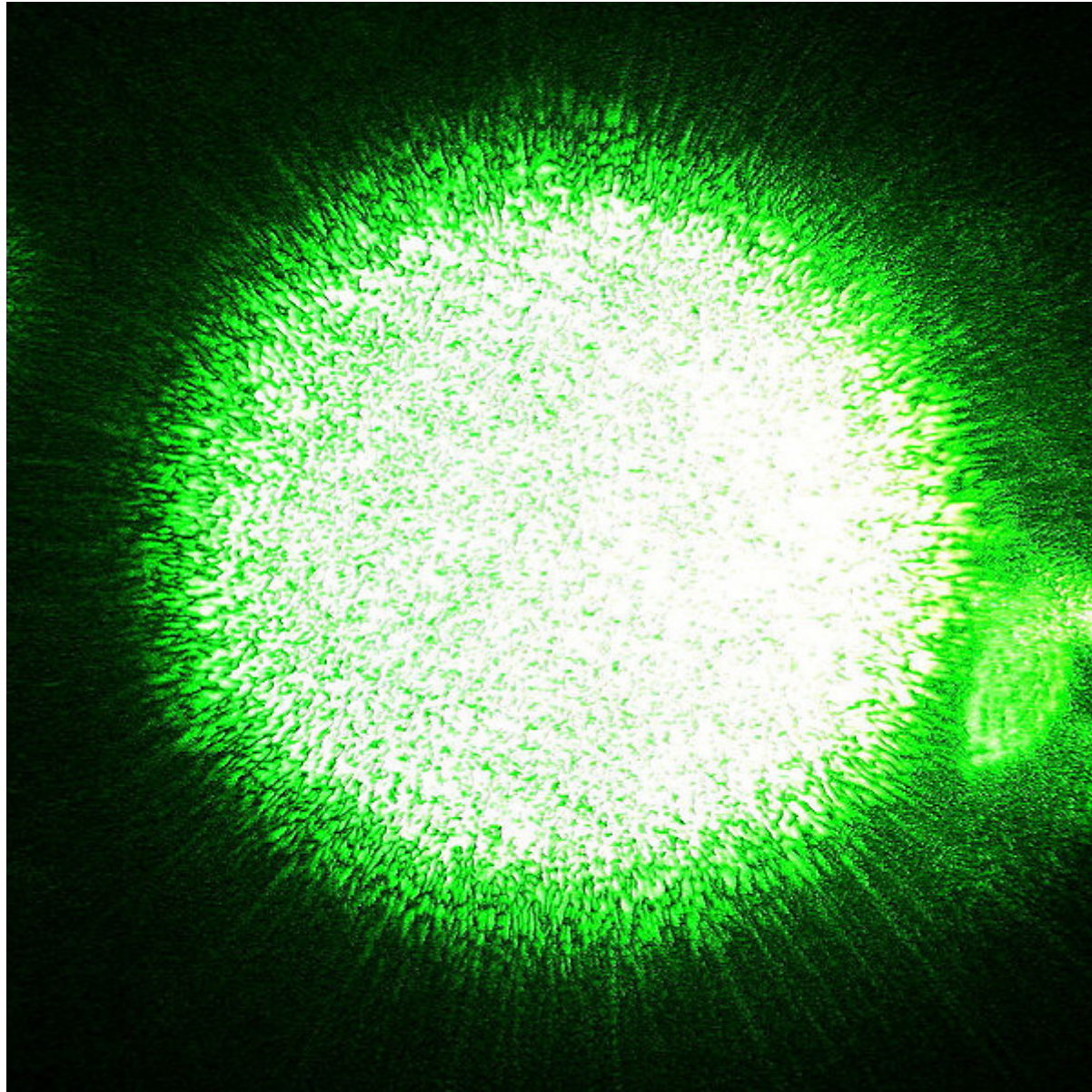
Incoherent light

Coherent light

Close up



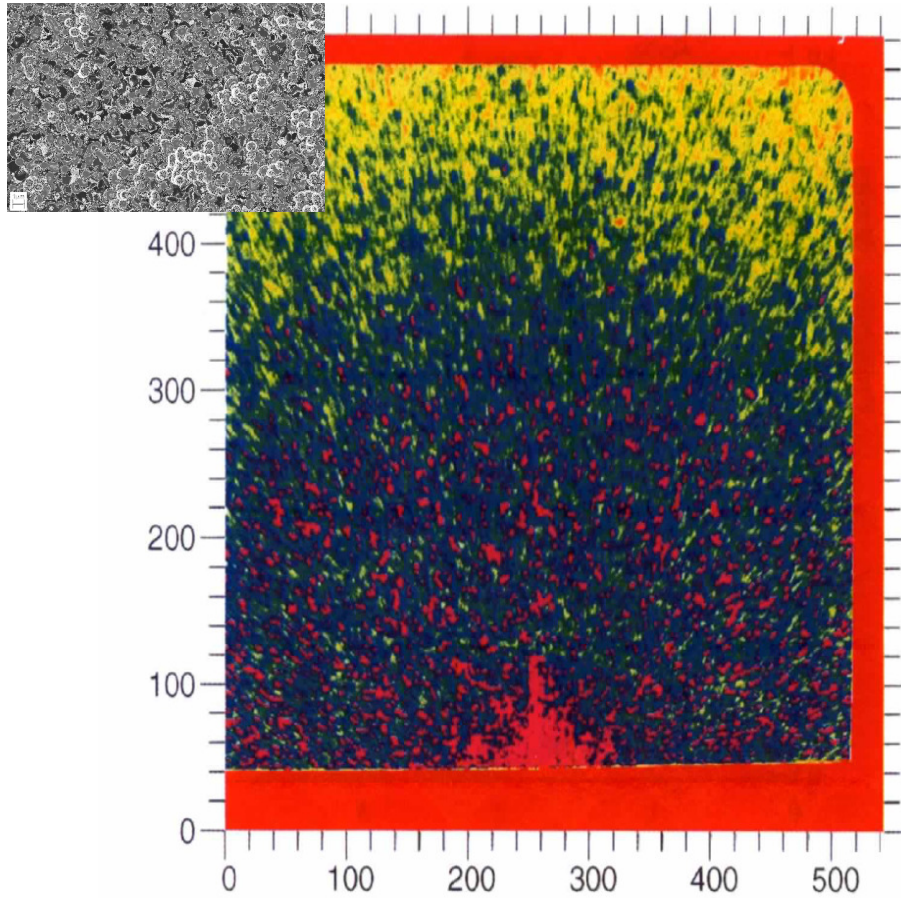
# Laser Speckles



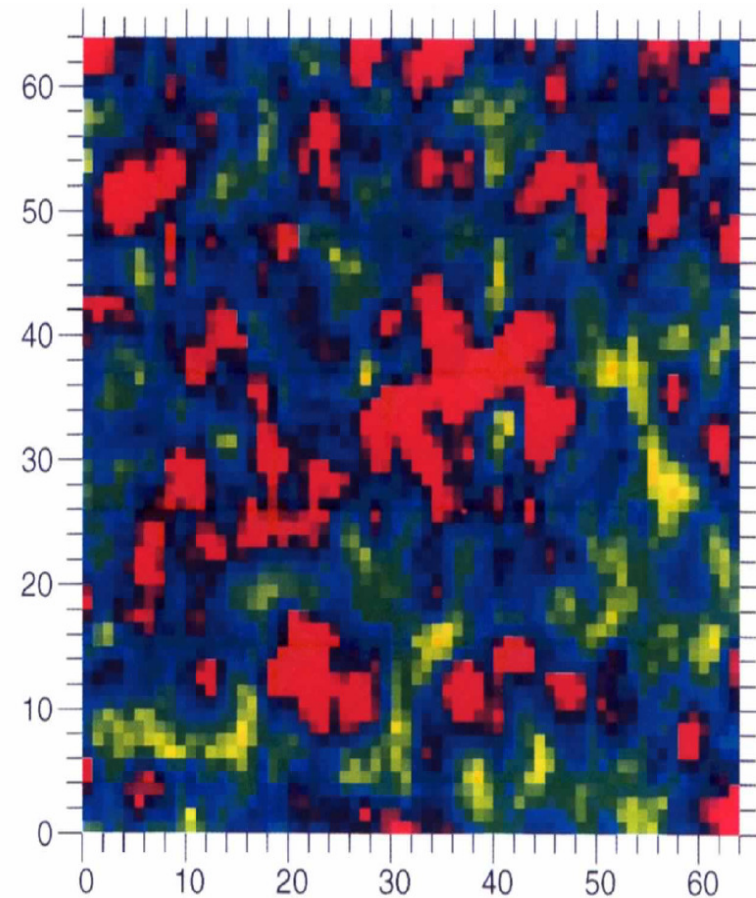


# X-ray speckle pattern

Aerogel sample



**Figure 1**  
Aerogel speckle pattern.



**Figure 2**  
Detail of the aerogel speckle pattern at  $0.006 \text{ \AA}^{-1}$ .

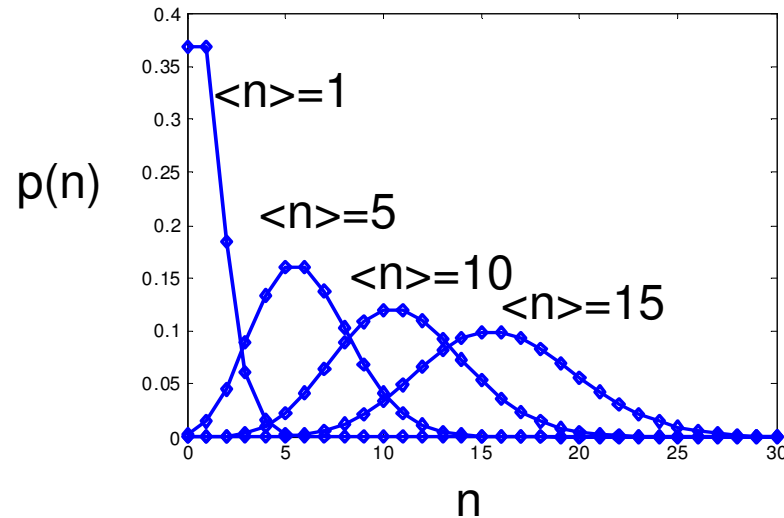
O.K.C. Tsui et al., Statistical Analysis of X-ray Speckle at the NSLS  
J. Synchrotron Rad. (1998),5, 30

# Analysis of Speckle Pattern in terms of counting statistics (1)

incoherent light

$$p(n) = \frac{\langle n \rangle^n e^{-\langle n \rangle}}{n!}$$

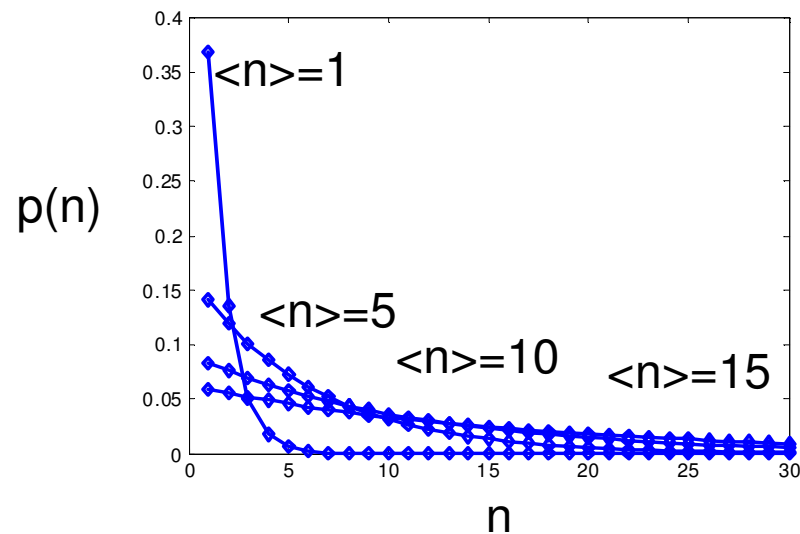
Poisson distribution



coherent light

$$p(n) = \frac{e^{-n/\langle n \rangle}}{\langle n \rangle}$$

exponential distribution



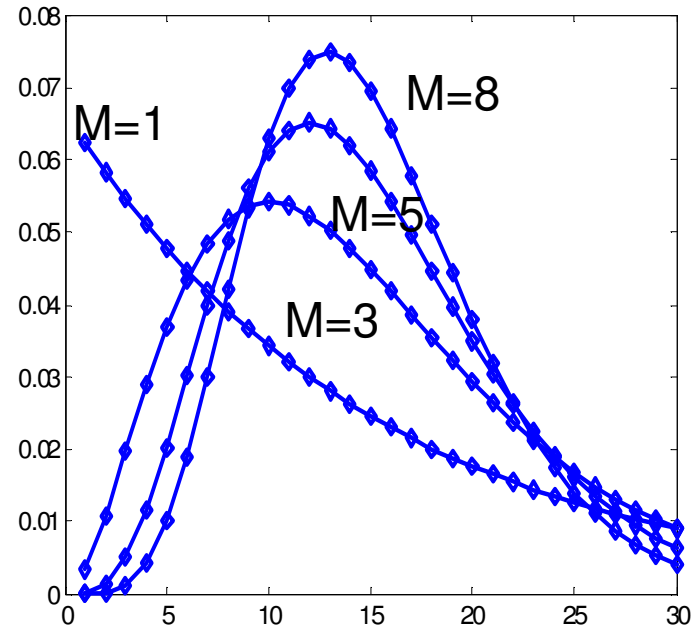
# Analysis of Speckle Pattern in terms of counting statistics (2)

Approximation partial coherent light

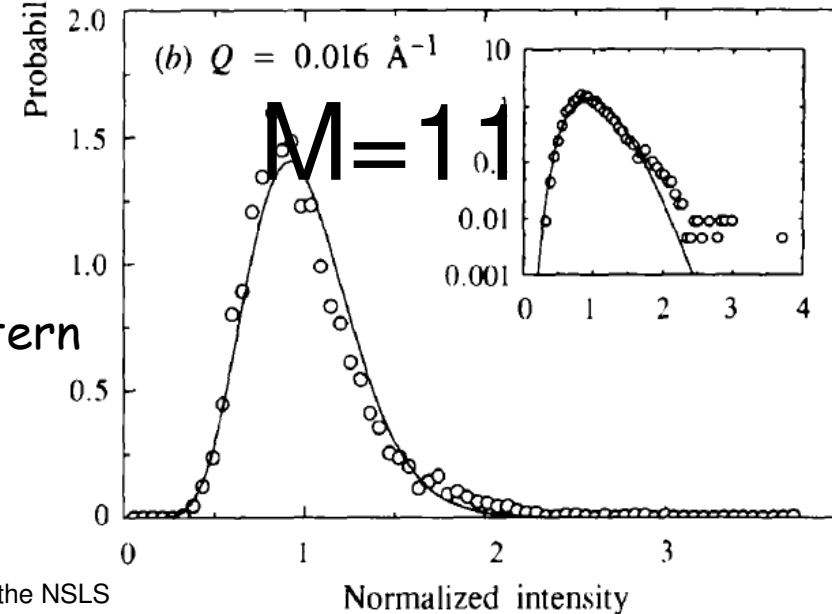
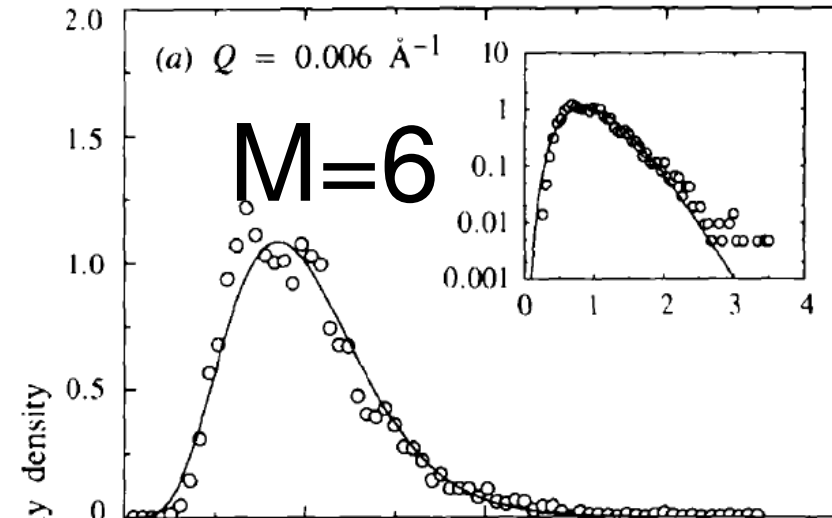
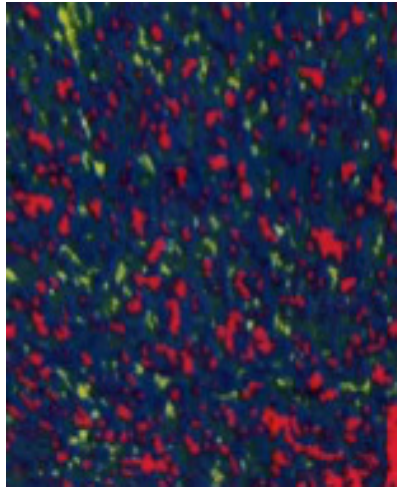
$$p(n) = M^M \left( \frac{n}{\langle n \rangle} \right)^{M-1} \frac{e^{-Mn/\langle n \rangle}}{\Gamma(M) \langle n \rangle^M}$$

degree of coherence  $1/M$

$$\langle n \rangle = 15$$



# X-ray speckle pattern at synchrotron

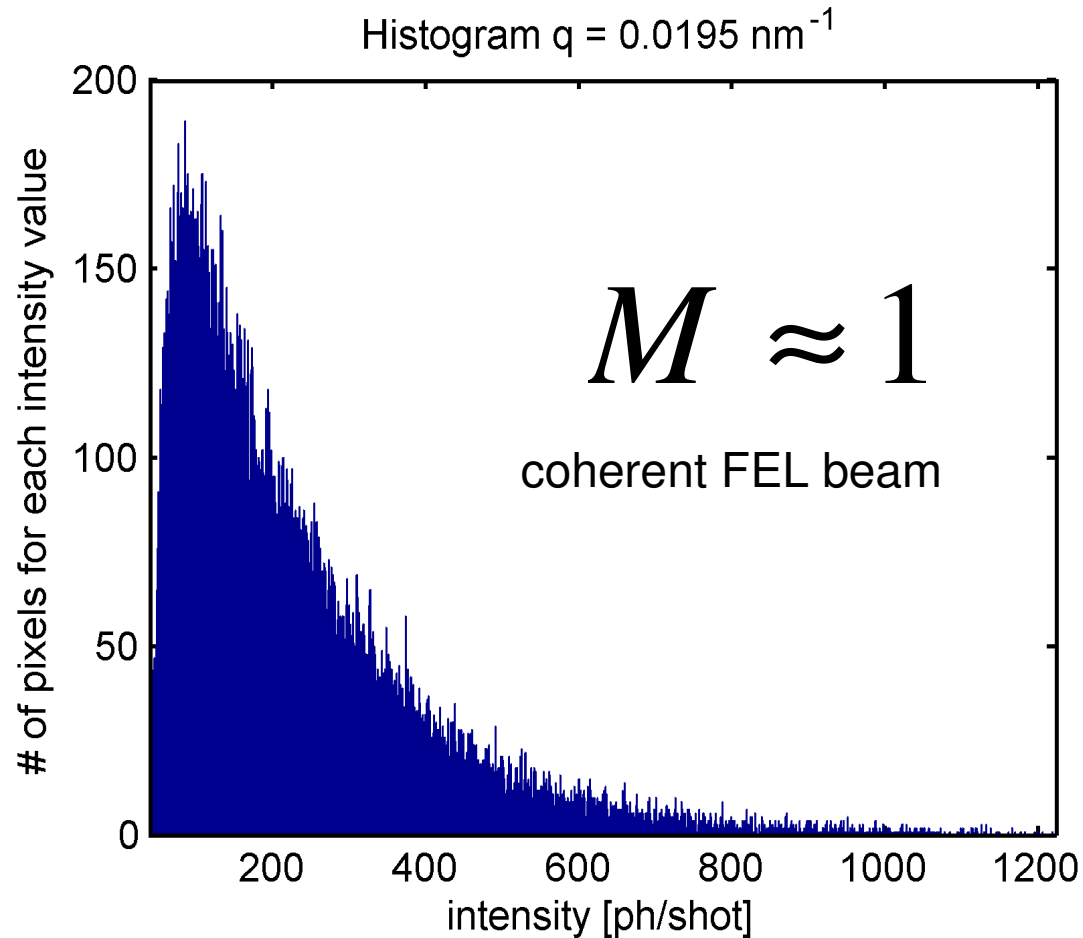


$$p(n) = M^M \left( \frac{n}{\langle n \rangle} \right)^{M-1} \frac{e^{-Mn/\langle n \rangle}}{\Gamma(M) \langle n \rangle^M}$$

Contrast ('Visibility') of speckle pattern

$$C = \frac{\sigma}{\langle n \rangle} \propto \frac{1}{\sqrt{M}}$$

# Soft X-ray speckle pattern at FLASH



$$p(n) = M^M \left( \frac{n}{\langle n \rangle} \right)^{M-1} \frac{e^{-Mn/\langle n \rangle}}{\Gamma(M) \langle n \rangle}$$

# Applications

## Coherent X-ray Scattering

### Structure

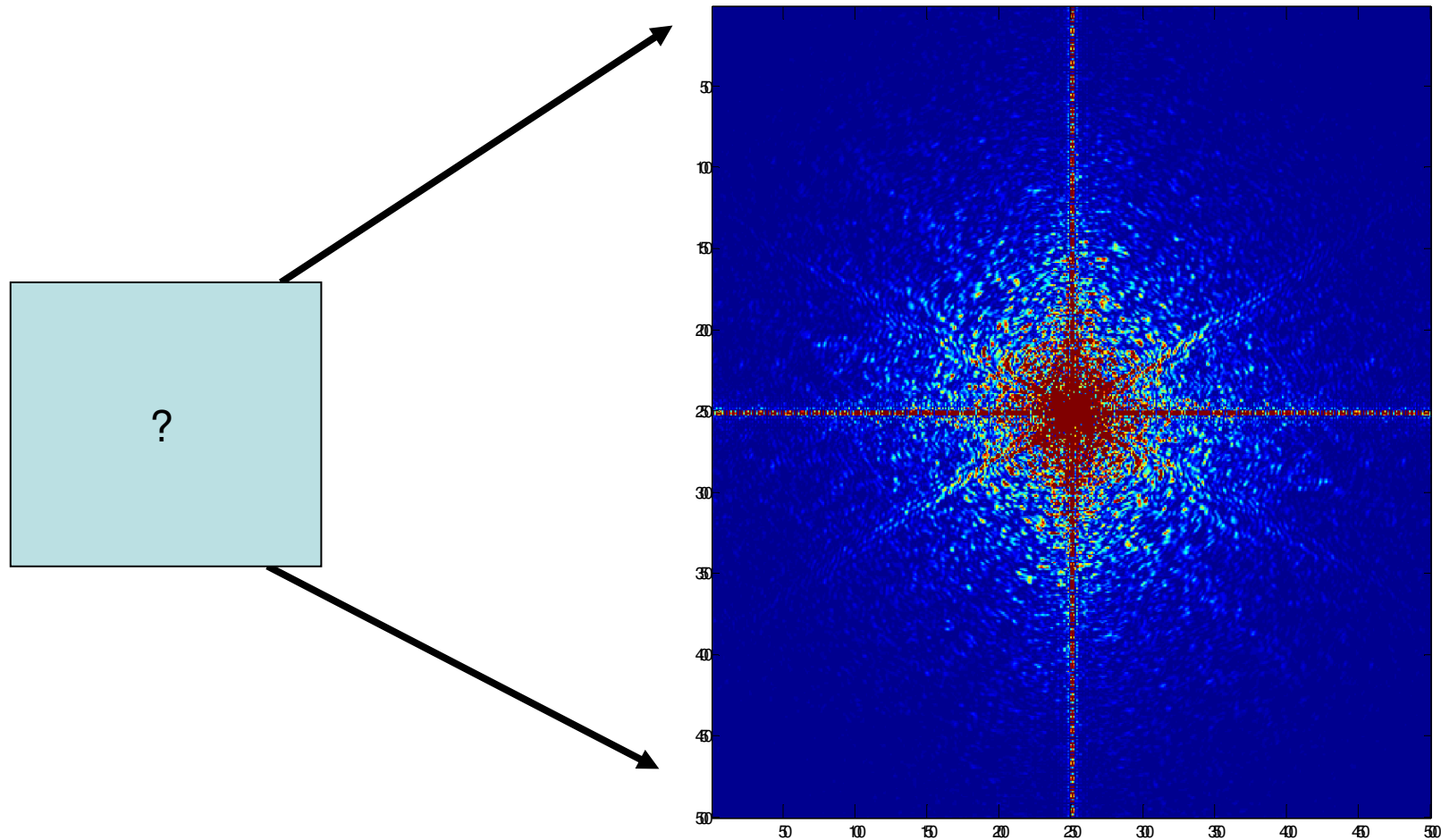
- Coherent Diffractive Imaging
- Holography

### Dynamics

- X-Ray Photon Correlation Spectroscopy



# Speckle Pattern for Coherent Diffractive Imaging



Happy X-mas - see you 08.01.2009