

Methoden moderner Röntgenphysik II: Streuung und Abbildung

| | |
|------------|---|
| Lecture 26 | Vorlesung zum Haupt- oder Masterstudiengang Physik, SoSe 2021 G. Grübel, O. Seeck, V. Markmann, F. Lehmkühler, <u>A. Philippi-Kobs</u> , M. Martins |
| Location | online |
| Date | Tuesdays 12:30 - 14:00 (starting 6.4.) Thursdays 8:30 - 10:00 (until 6.7.) |

Outline

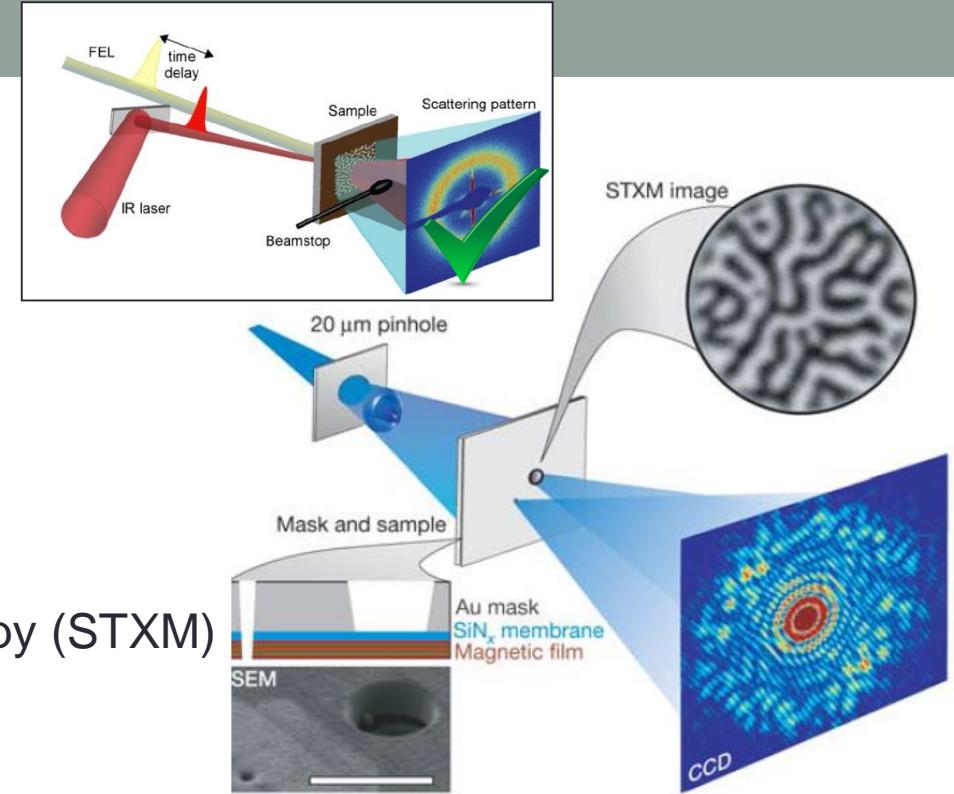
Part II/3:

Studies on Magnetic Nanostructures

by André Philippi-Kobs (AP)

[6.7.] Imaging of Magnetic Domains

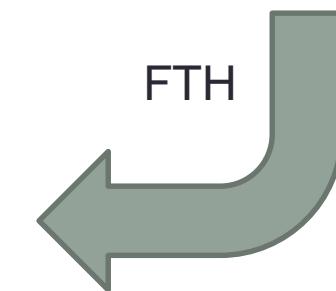
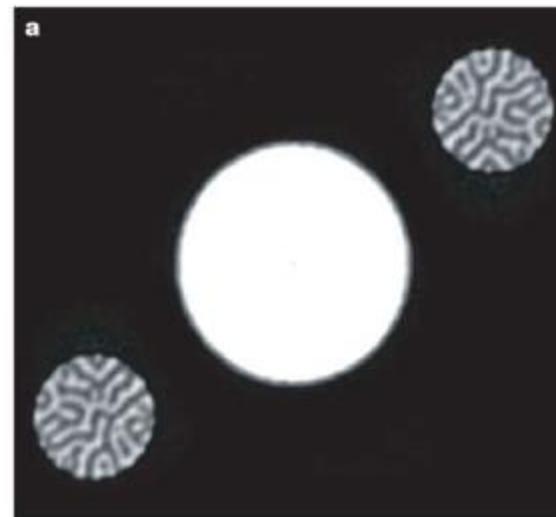
- **Fourier Transform Holography (FTH)**
- Scanning Transmission X-ray Microscopy (STXM)
- Coherent Diffraction Imaging (CDI)



Lensless imaging of magnetic nanostructures by X-ray spectro-holography

S. Eisebitt¹, J. Lüning², W. F. Schlötter^{2,3}, M. Lörgen¹, O. Hellwig^{1,4},
W. Eberhardt¹ & J. Stöhr²

NATURE | VOL 432 | 16 DECEMBER 2004 |

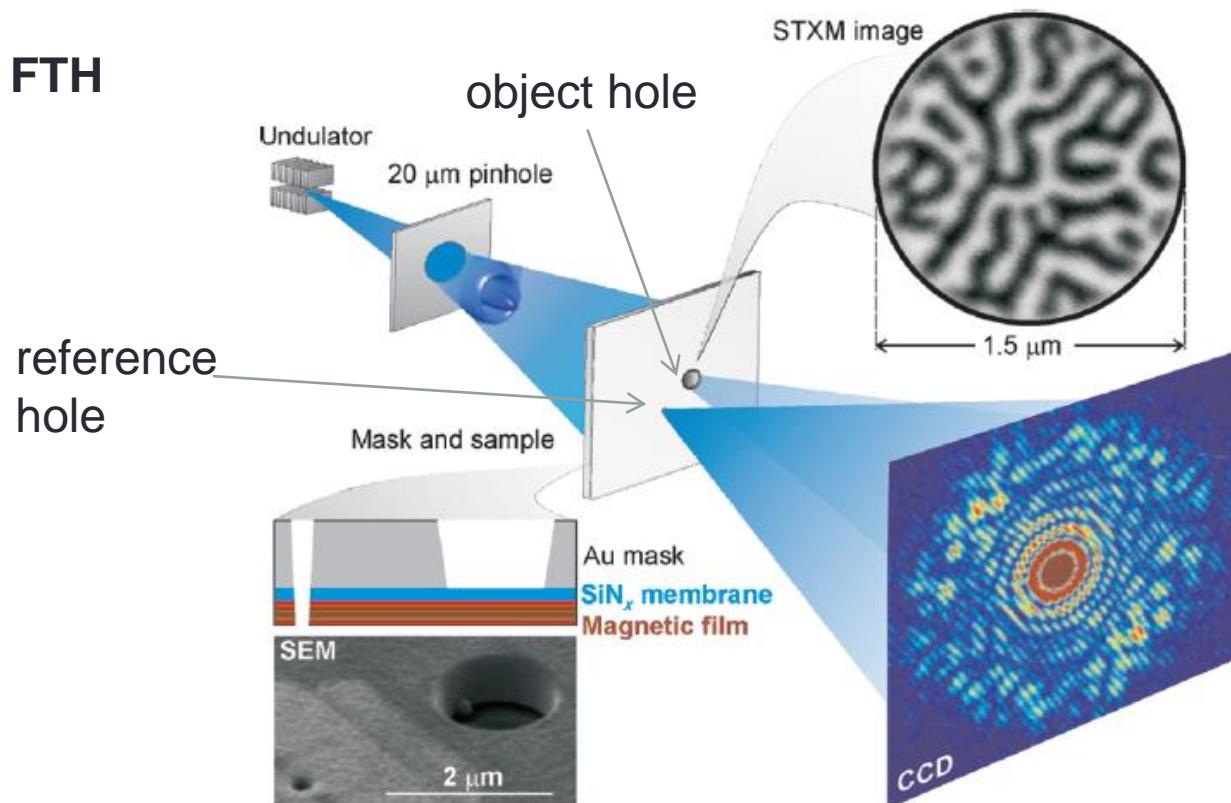


Imaging of magnetic domain patterns with X-rays

► Brief excursion on Lensless Imaging – Fourier transform Holography

| | Schlüsselement-Herstellung | Bild-Rekonstruktion | |
|-----|----------------------------|---------------------|---------------------------------|
| TXM | Zonenplatte | XXXXX | -direkt- |
| FTH | Optikmaske | XX | Einfache Fourier-Transformation |
| CDI | -direkt- | X | Phasen-Rückgewinnung |

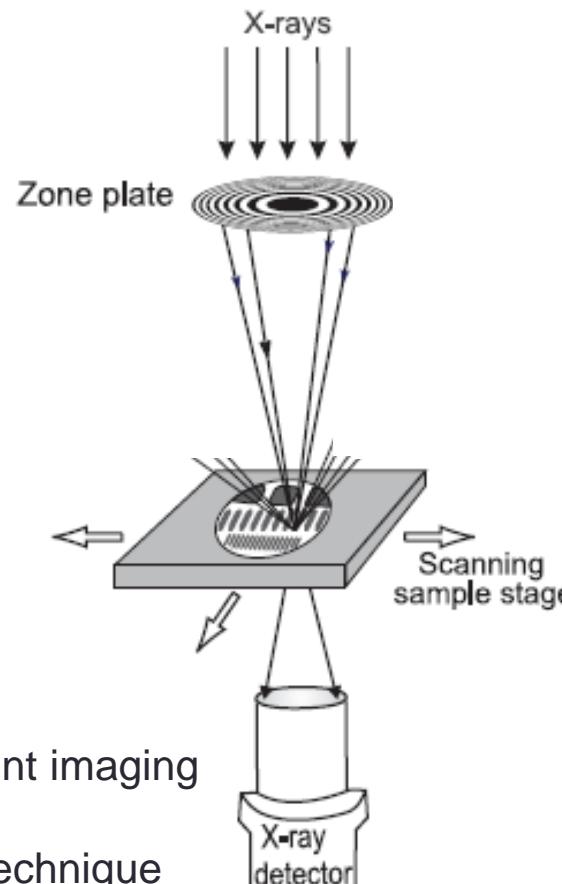
FTH



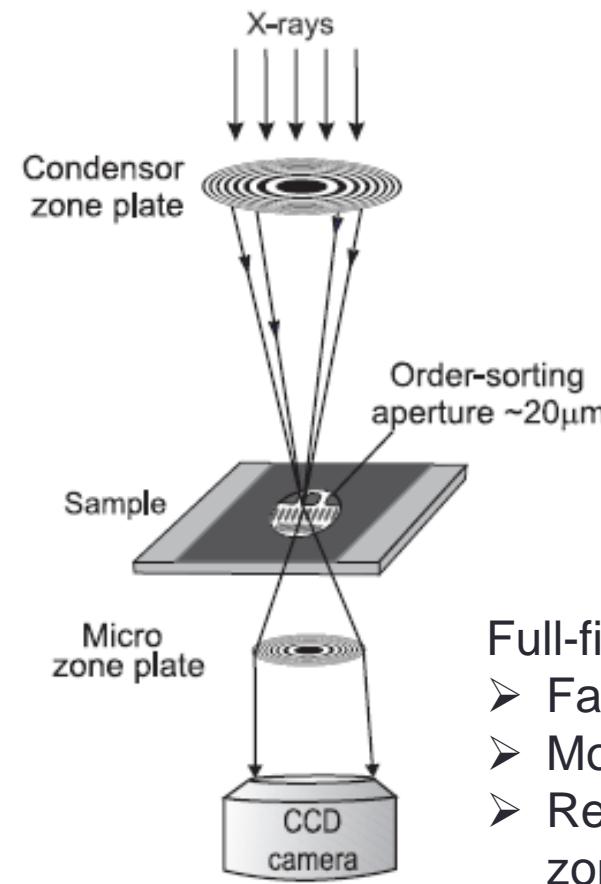
Imaging of magnetic domain patterns with X-rays

> X-ray lenses based methods

Scanning Transmission X-ray Microscopy
STXM



Transmission Imaging X-ray Microscopy
TIXM



Point-by-point imaging

- Slow
- Simple technique
- Resolution limit = focal-spot size

Full-field microscopy

- Fast (detector limited)
- More difficult
- Resolution limit by zone plate



Imaging of magnetic domain patterns with X-rays

➤ X-ray lens-based method

Fresnel Zone plates:

Condition for constructive interference at focal distance f

$$r_n = \sqrt{m\lambda f + \frac{m^2\lambda^2}{4}} \approx \sqrt{m\lambda f}$$

All zones have the same area

$$A_m = \pi(r_{m+1}^2 - r_m^2) = \pi\lambda f$$

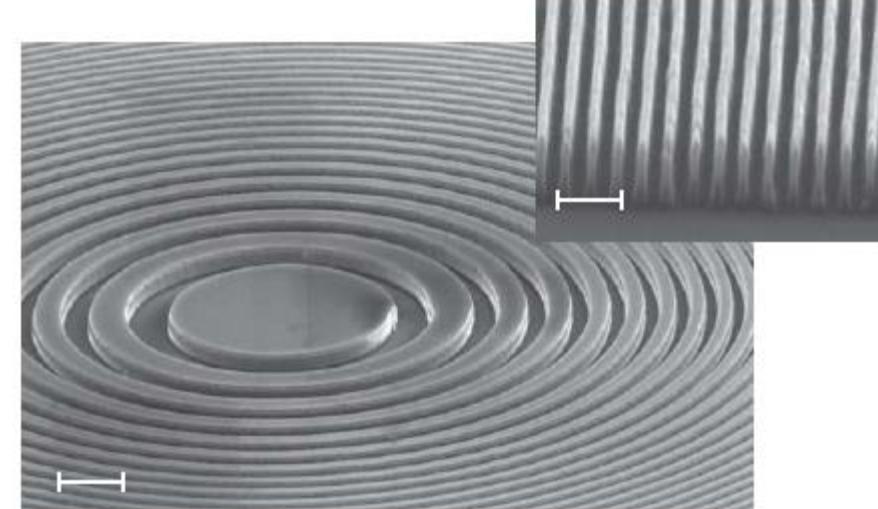
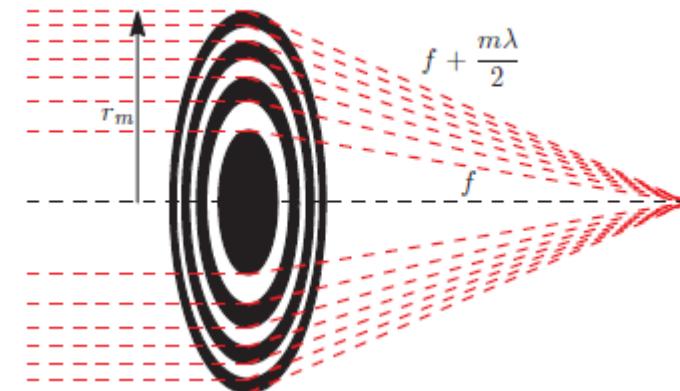
Resolution determined by width of outermost zone

$\Delta x = 1.22\Delta r_m \approx 10\text{nm}$ nowadays
 (7.8 nm FWHM shown in 2017 [Sci. Rep. 7, 43624]).

disadvantages:

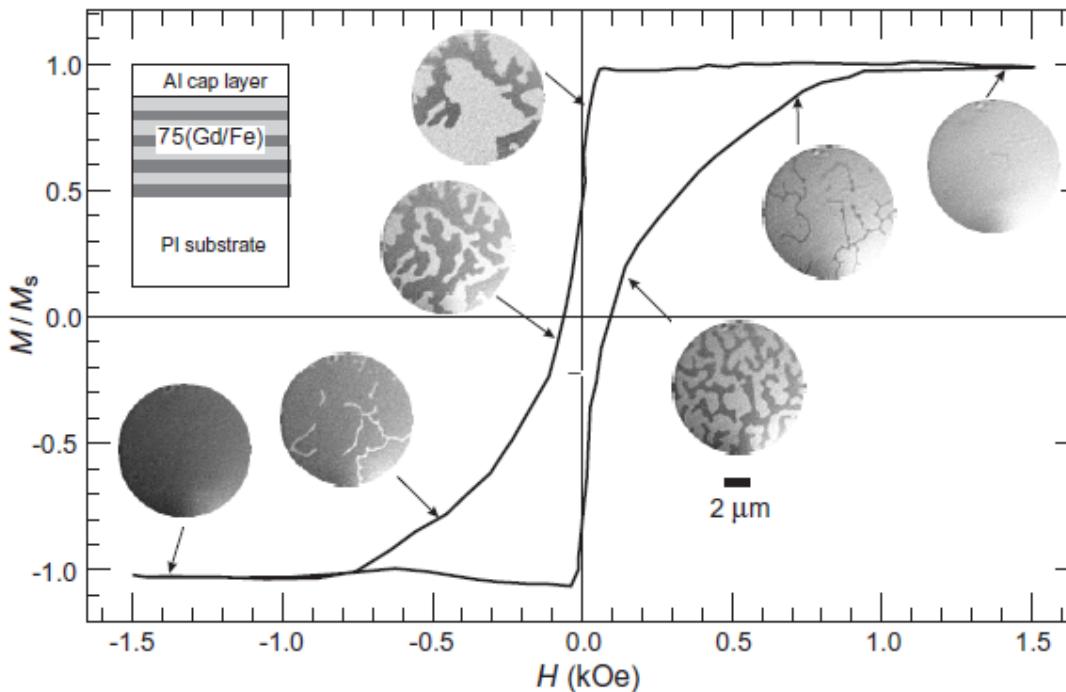
- High absorption (5-30% efficiency)
- Hard to fabricate

See, e.g., X-Ray Data Booklet, Sec 4.4
<http://www.x-ray-optics.de>



Imaging of magnetic domain patterns with X-rays

➤ X-ray lense-based method



- Element-sensitivity
- Integration of gray values for each field value
→ hysteresis

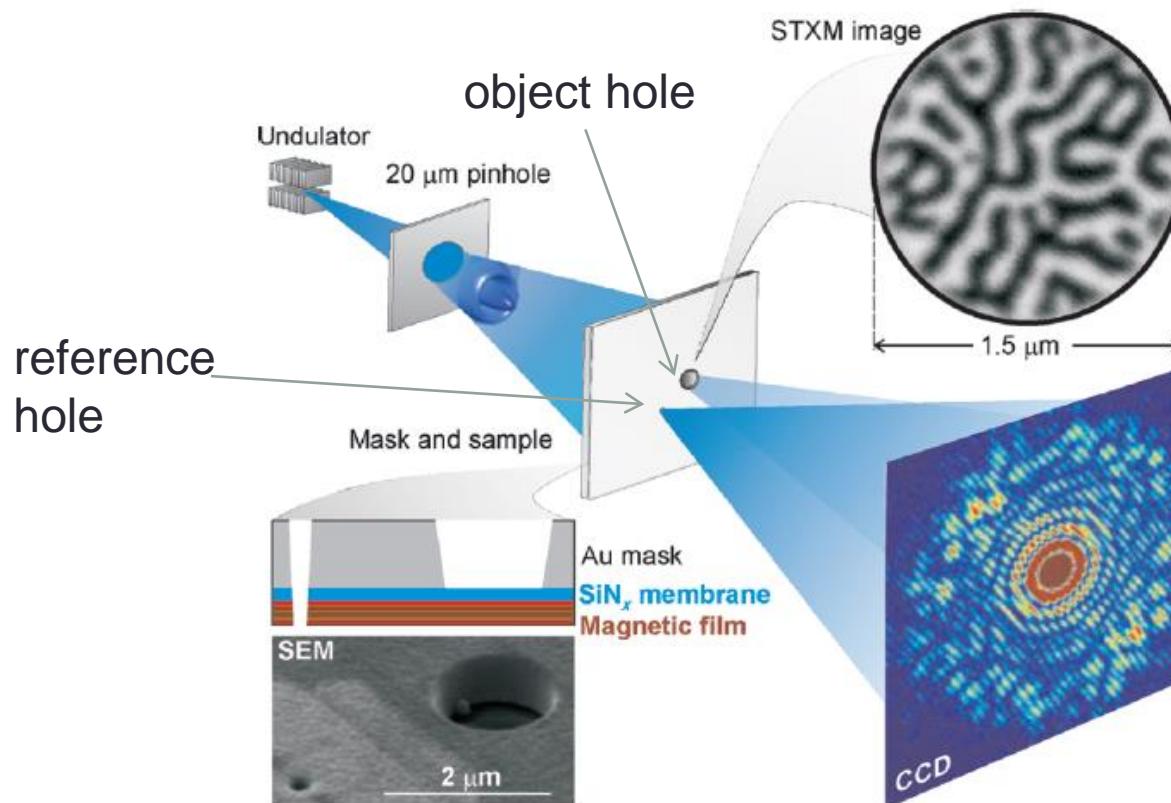
Fig. 10.22. TIXM images recorded at the FeL₃-edge as a function of applied field for a $75 \times [\text{Fe}(4.1 \text{ \AA})/\text{Gd}(4.5 \text{ \AA})]$ multilayer deposited on polyimide and capped for protection with an Al layer [463, 482]

Stöhr and Siegmann, Magnetism

Imaging of magnetic domain patterns with X-rays

► Brief excursion on Lensless Imaging – Fourier transform Holography

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Requirement:
transversal coherence length > largest length scale in sample = mask diameter

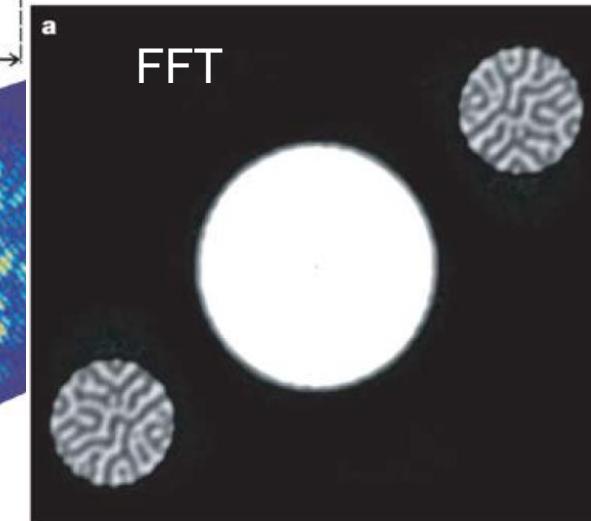
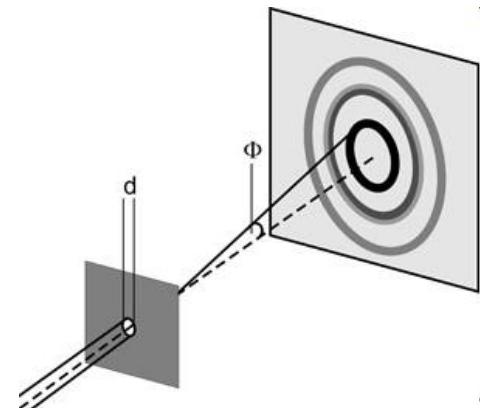
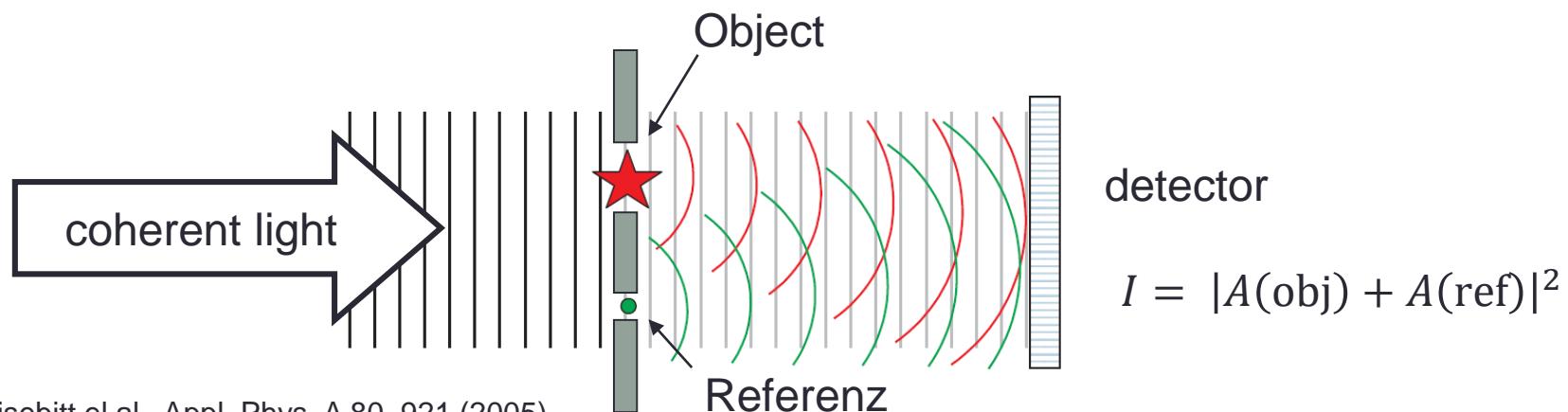


Image the nanoscale – Fourier-transform holography

Usually, in a scattering experiment the phase is lost – only the intensity is recorded



Solution: ‘Detect’ phase via interference of object wave with a reference wave

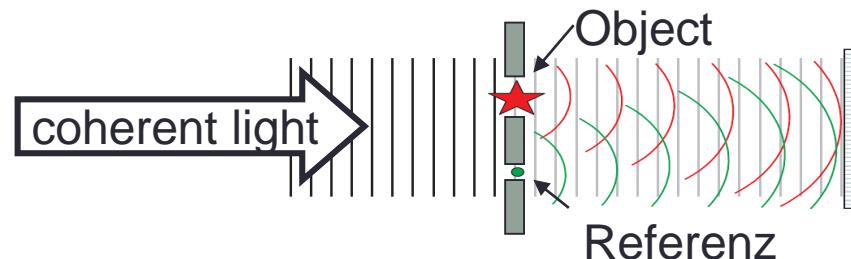


S. Eisebitt et al., Appl. Phys. A 80, 921 (2005)

Object and reference wave interfere if object and reference are illuminated coherently and the (polarization) state of the light is not changed by the scattering...

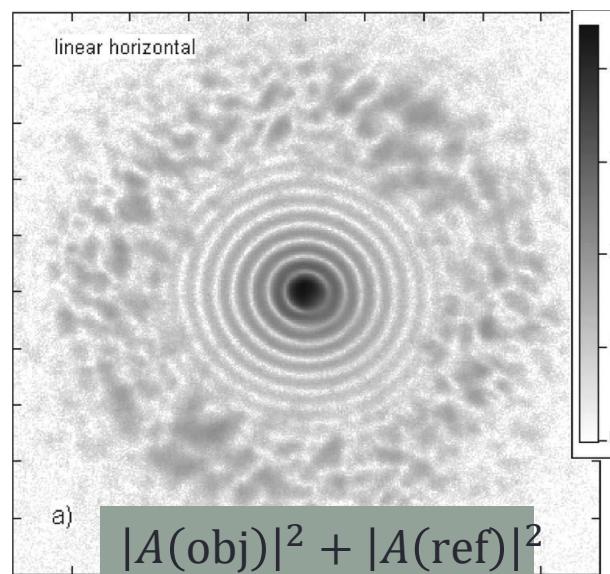
Image the magnetic nanoscale – magnetic FTH

... unfortunately, magnetic scattering does exactly that with linear light



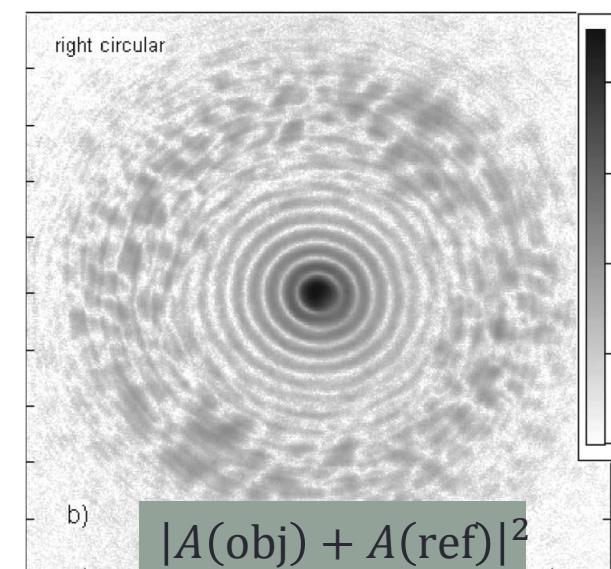
$$\text{detector } I_{\text{lin}} = |A(\text{obj})|^2 + |A(\text{ref})|^2$$

S. Eisebitt et al., Appl. Phys. A 80, 921 (2005)



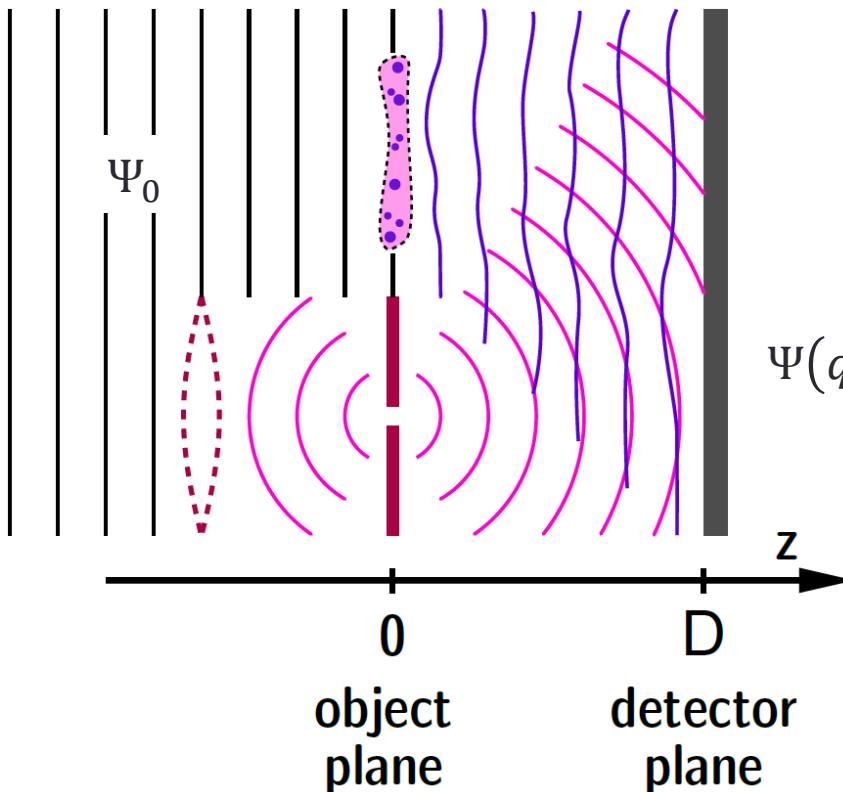
Linear light

S. Eisebitt et al., PRB, 68 104419 (2003).



Circular light

FTH image formation



$p(r')$ is the spatial autocorrelation function or the Patterson map

- Sample transmits incoming wavefield $\Psi(x, y) = t(x, y)\Psi_0$
- In the far field $\left(F \ll 1, F = \frac{\ell^2}{D\lambda}\right)$ one measures a Fraunhofer pattern described by

$$\Psi(q_x, q_y) = \iint_{-\infty}^{\infty} \Psi(x, y) \exp(-i(q_x x + q_y y)) dx dy$$

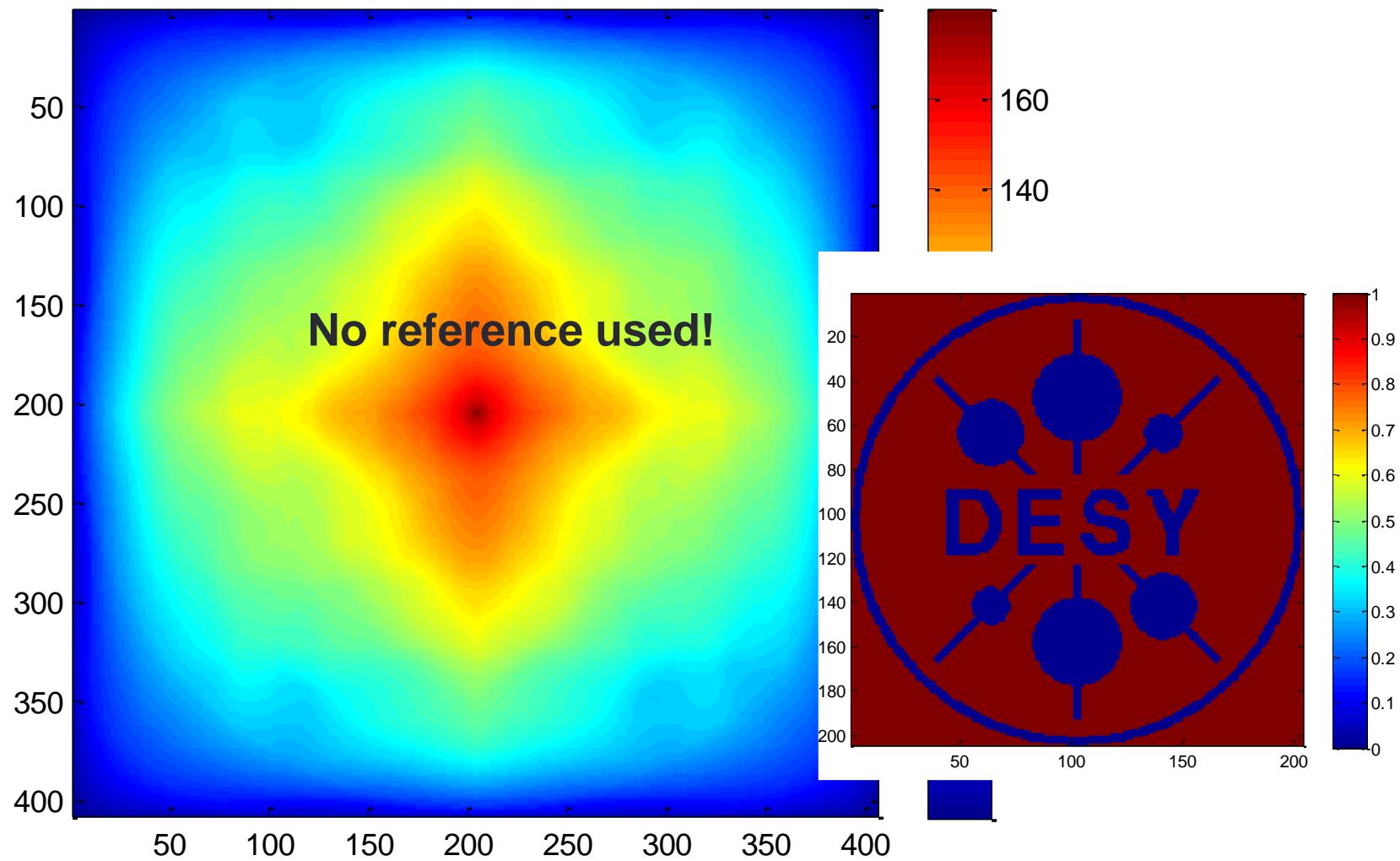
This is the Fourier transform of the transmitted wavefield which again is directly related to the sample structure.

- Unfortunately we measure the absolute square of the wavefield and the phase is lost. Hence we get

$$\begin{aligned} p(r') &= F^{-1}\{F^*(\Psi(r))F(\Psi(r))\} \\ &= \Psi^*(-r) * \Psi(r) \\ &= \iint_{-\infty}^{\infty} \Psi^*(-r)\Psi(r+r') dr \end{aligned}$$



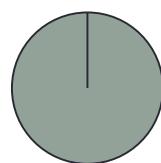
What does the Patterson map tell us?



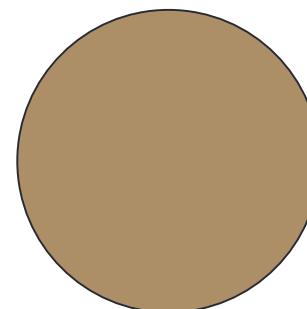
The trick of the reference

- No reference

Object



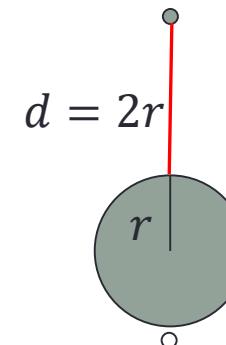
Patterson Map



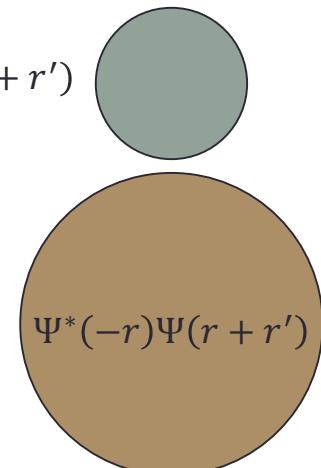
$$\Psi^*(-r)\Psi(r + r')$$

Needs iterative phase retrieval (CDI)
Not necessarily 'easy'

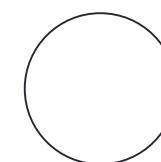
- With reference



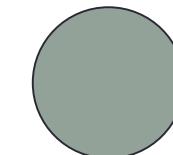
$$\Psi^*(-r)R(r + r')$$



$$\Psi^*(-r)\Psi(r + r')$$



$$R^*(-r)\Psi(r + r')$$

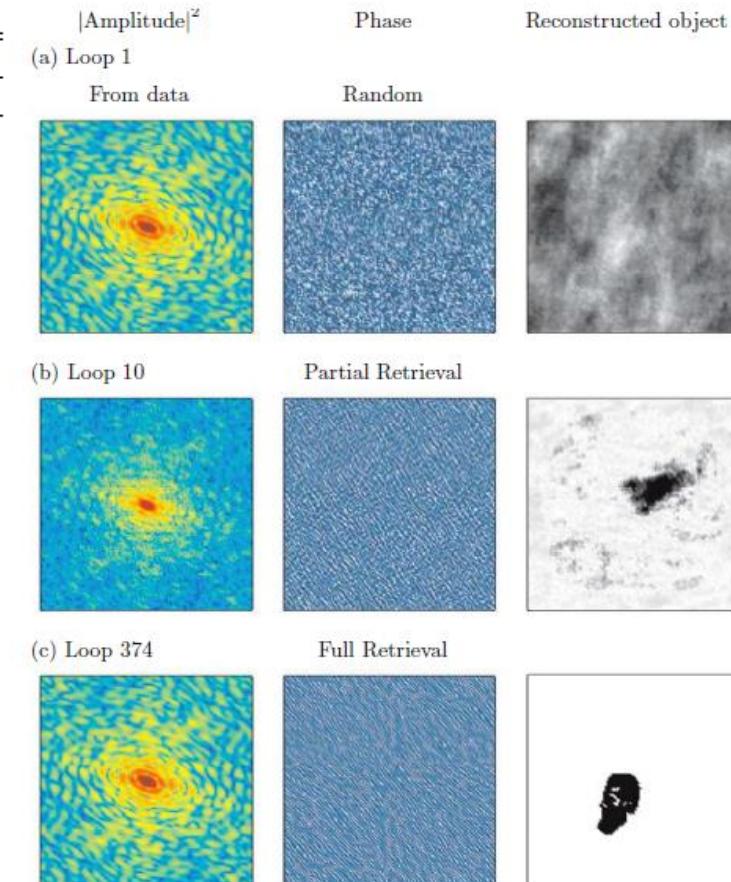
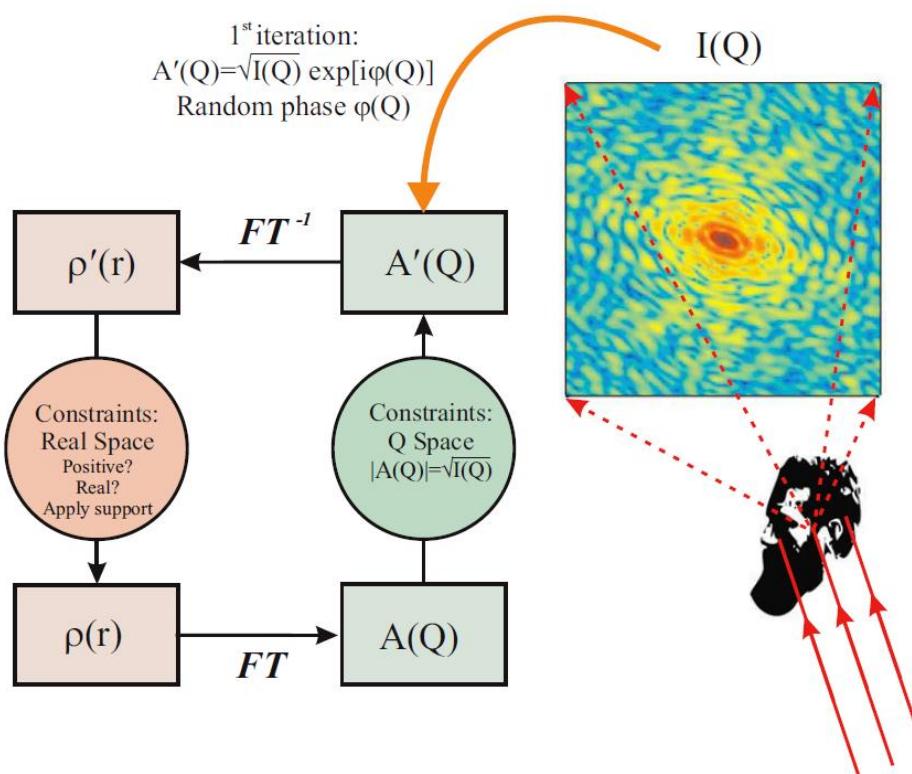


Real space image comes from FFT
Resolution limited by reference size

Imaging of magnetic domain patterns with X-rays

> Lensless Imaging – Coherent Diffraction Imaging (CDI)

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| FTH | Optikmaske | XX | Einfache Fourier-Transformation | XX |
| CDI | -direkt- | X | Phasen-Rückgewinnung | XXXXX |

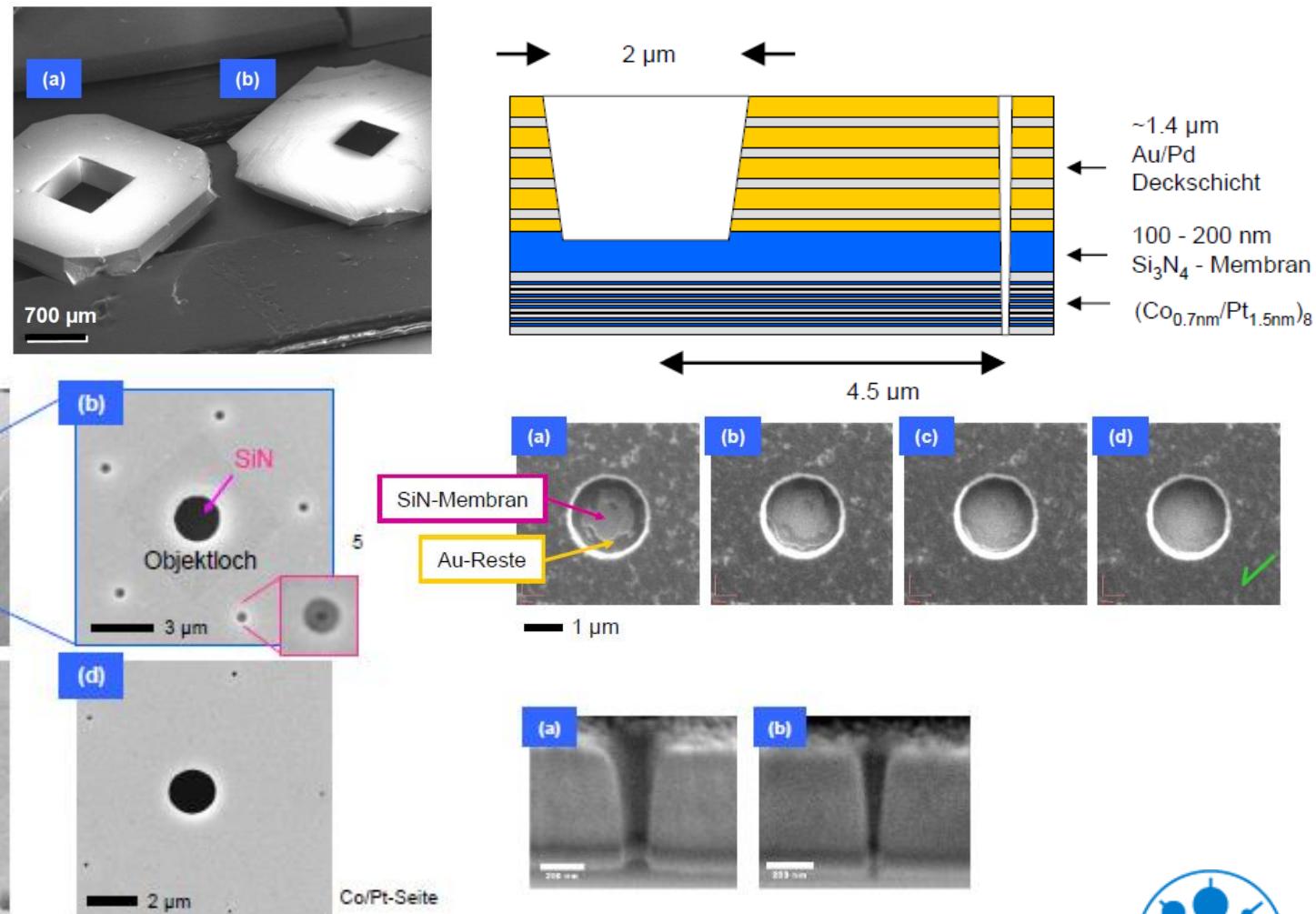


Imaging of magnetic domain patterns with X-rays

> Lensless Imaging – Fourier transform Holography

Mask and sample:

Preparation
by focused ion beam
technique



Imaging of magnetic domain patterns with X-rays

> Lensless Imaging – Fourier transform Holography (FTH)

Principle:

$$\tilde{f}(\vec{Q}) \quad | \quad \tilde{f}_j(\vec{Q}) = \mathcal{FT}(f_j(r))$$

- Intensity on detector: $I(\vec{Q}) = \left| \sum_j f_j(\vec{Q}) e^{i\vec{Q} \cdot \vec{r}_j} \right|^2$

- Scattering factor for circularly polarized light and $\mathbf{M} \parallel \mathbf{L}_{ph}$:

$$f = \vec{\epsilon} \cdot \vec{\epsilon}' F^c - i (\vec{\epsilon} \times \vec{\epsilon}') \cdot \vec{n} F^n$$

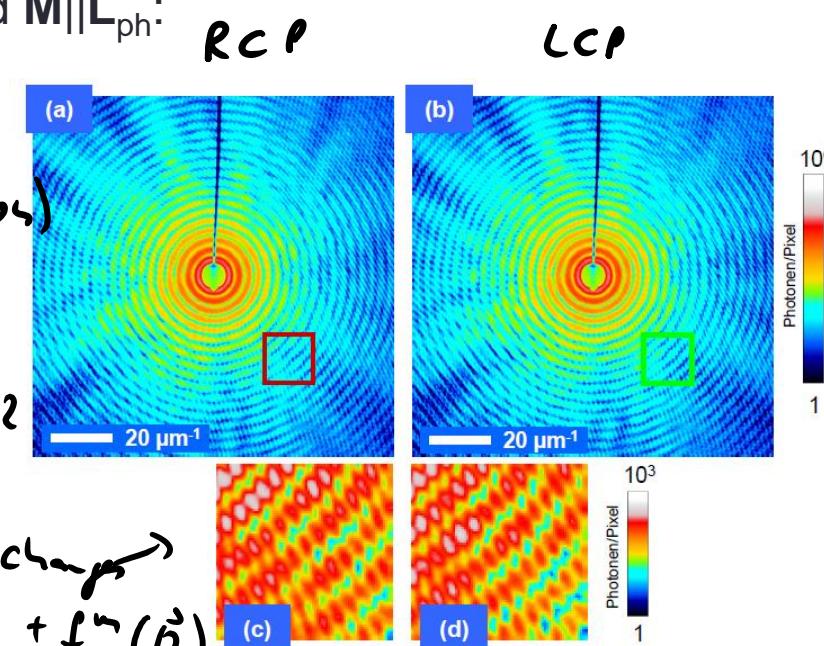
$$f = f^c(\vec{Q}) \pm f^n(\vec{Q}) \quad (+RCP, -LCP)$$

$$f^c(\vec{Q}) = f_0^c(Q) + f_R^c(\vec{Q})$$

- "Hologram" (= $I(\vec{Q})$) with RCP and LCP

$$I(\vec{Q}) = \left| \tilde{f}_0^c(\vec{Q}) + \tilde{f}_R^c(\vec{Q}) \pm \tilde{f}_0^n(\vec{Q}) \right|^2$$

phase change
due to $\pm \tilde{f}_0^n(\vec{Q})$



Imaging of magnetic domain patterns with X-rays

> Lensless Imaging – Fourier transform Holography (FTH)

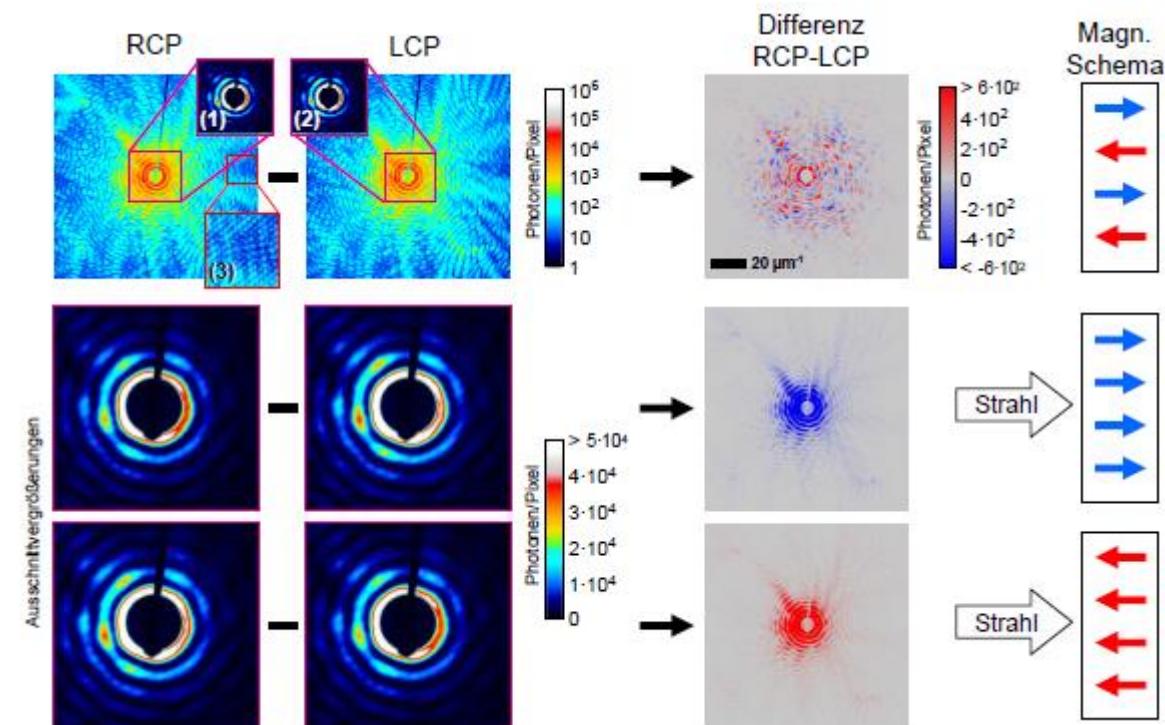
Principle:

- Difference hologram: $\Delta \mathcal{I}(\vec{Q})$

$$= \mathcal{I}^+(\vec{Q}) - \mathcal{I}^-(\vec{Q})$$

$$= \tilde{f}_o^{\perp \times} \cdot \tilde{f}_o^c + \tilde{f}_o^{\perp} \cdot f_o^{c \times}$$

$$+ \tilde{f}_o^{\perp \times} \tilde{f}_R^c + \tilde{f}_o^{\perp} \cdot \tilde{f}_R^{c \times}$$



Imaging of magnetic domain patterns with X-rays

> Lensless Imaging – Fourier transform Holography (FTH)

Principle: Note: $\widehat{FT}^{-1}(\tilde{f}(\vec{q})) = \widehat{FT}^{-1}\widehat{FT}(f(\vec{r}))$

- Reconstruction = Fourier transformation: $= f(\vec{r})$

$\widehat{FT}^{\gamma}(\Delta I(\vec{q}))$ ^{invers}

$$= \widehat{FT}^{-1}(\tilde{f}_o^{\gamma} \cdot \tilde{f}_o^c) + \widehat{FT}^{-1}(\tilde{f}_o^{\gamma} \cdot \tilde{f}_o^{c+})$$

Autocorrelation

$$+ \widehat{FT}^{-1}(\tilde{f}_o^{\gamma*} \cdot \tilde{f}_R^c) + \widehat{FT}^{-1}(\tilde{f}_o^{\gamma} \cdot \tilde{f}_R^{c+})$$

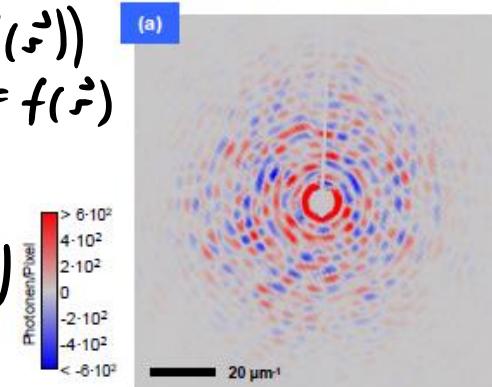
Reconstruction

$$\textcircled{1} = \widehat{FT}^{-1}(\tilde{f}_o^{\gamma}) \otimes \widehat{FT}^{-1}(f_R^{c+})$$

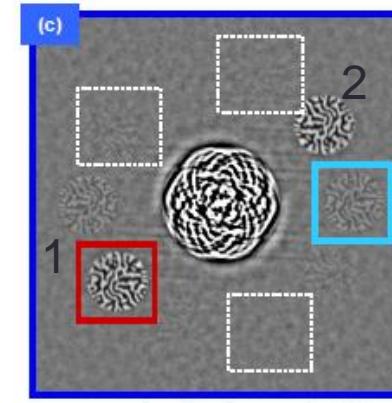
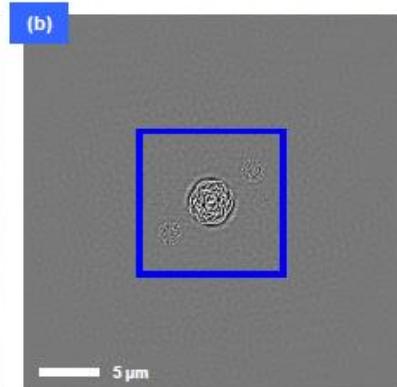
$$= f_o^{\gamma}(\vec{r}) \otimes f_R^{c+}(\vec{r}) \leftarrow \delta - \text{Filter}$$

$$= \int f_o^{\gamma}(\vec{r}) \delta(\vec{r} - \vec{r}') d\vec{r}' = \underline{f_o^{\gamma}(\vec{r})}$$

Differenzhologramm



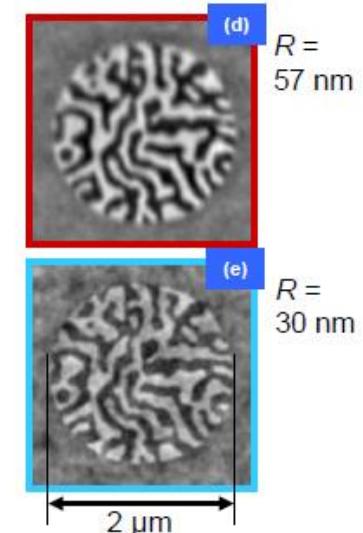
Fourier-Transformation



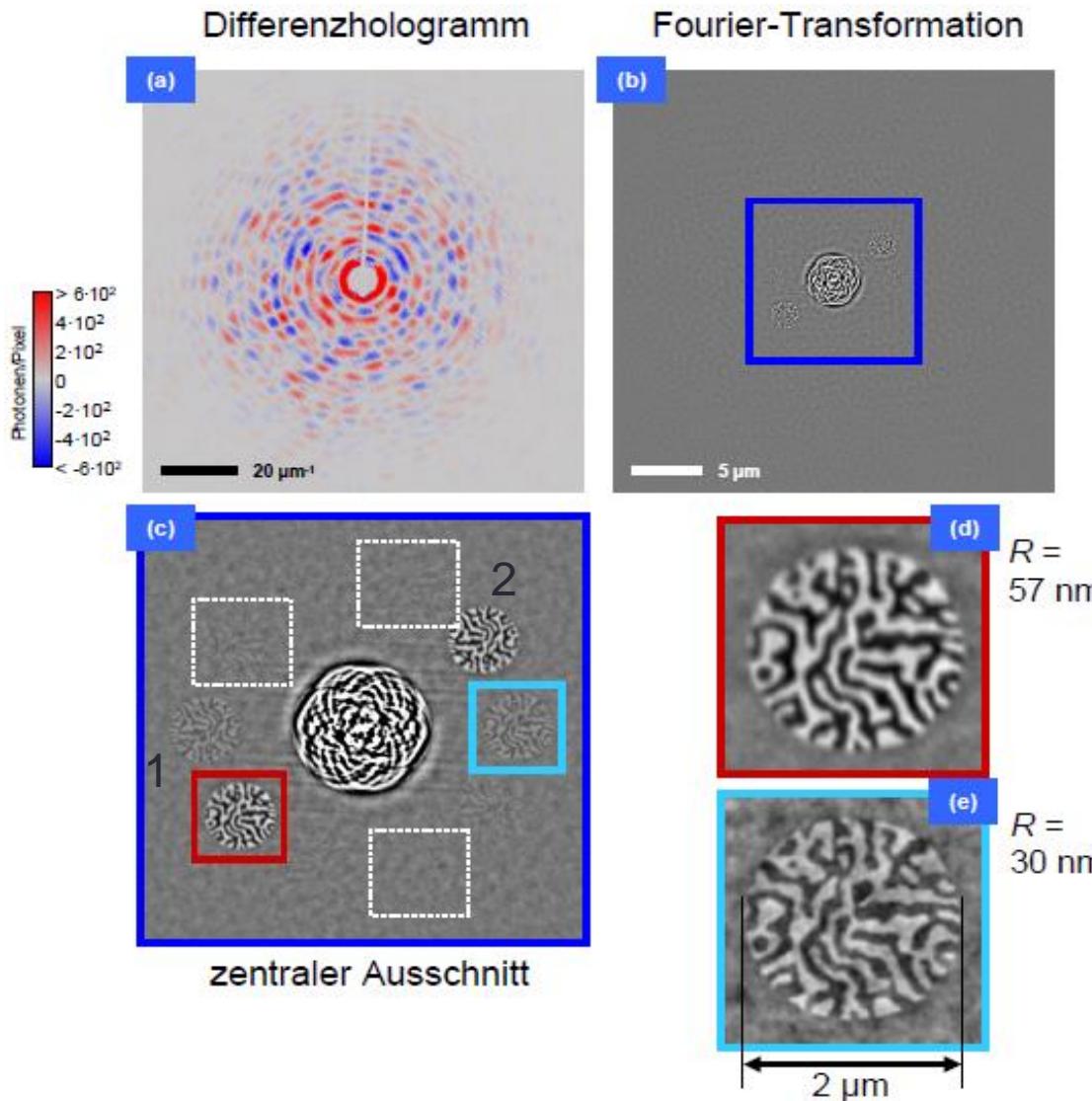
zentraler Ausschnitt

$$\textcircled{1} = -\tilde{f}_o(-\vec{r}')$$

$$R = 57 \text{ nm}$$



Imaging of magnetic domain patterns with X-rays



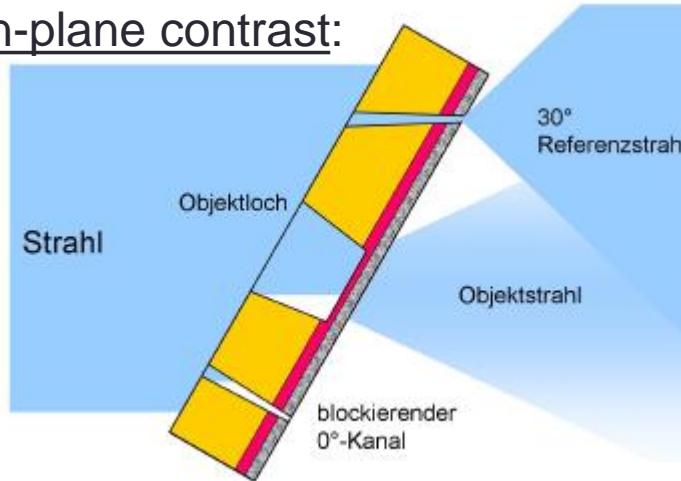
The FT of the Hologram shows the Patterson map of the Holographic mask with the magnetic structure. In the center there is the object-hole autocorrelation. The 10 reference-hole-object-hole cross correlations (2 for each reference hole) are arranged outside of that area. Large references yield high contrast but low resolution whereas small reference holes yield high resolution at low contrast.

Imaging of magnetic domain patterns with X-rays

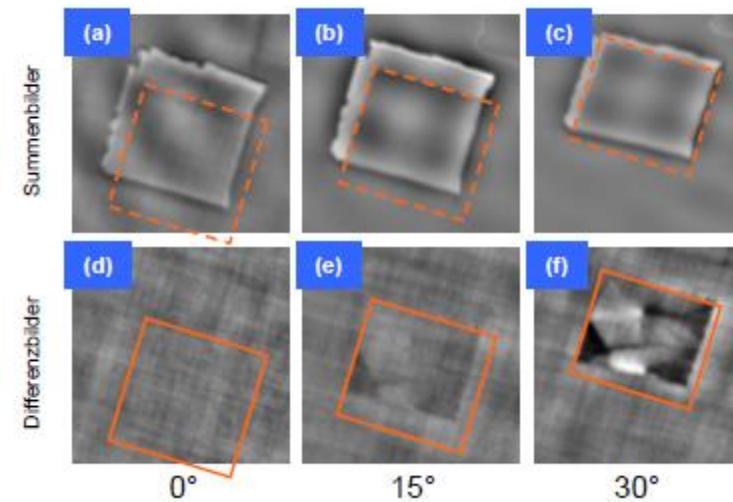
> Lensless Imaging – Fourier transform Holography (FTH)

References holes with different inclinations

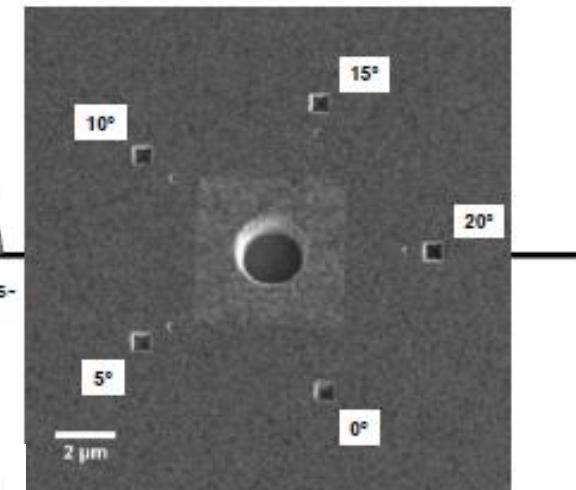
In-plane contrast:



In-plane magnetized
20 nm thick Co film

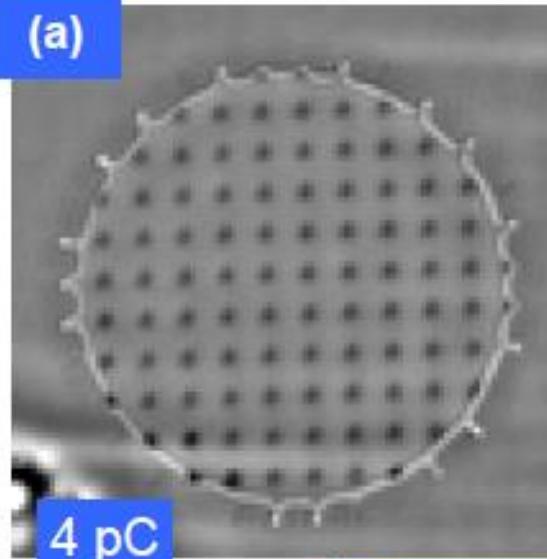


$$\Delta I_{XMCD} \propto \vec{M} \cdot \vec{L}_\gamma$$

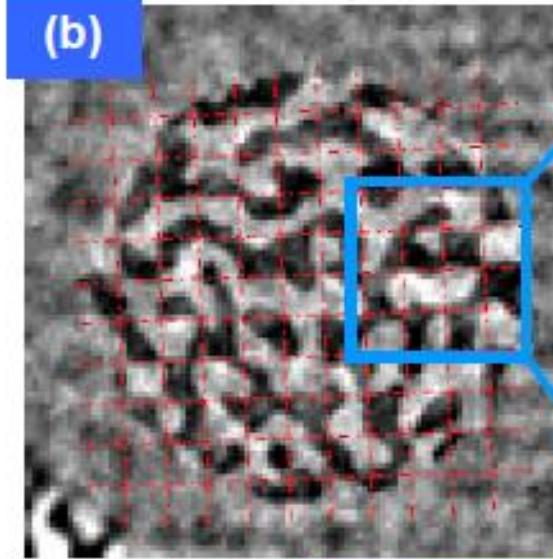


> Lensless Imaging – Fourier transform Holography (FTH)

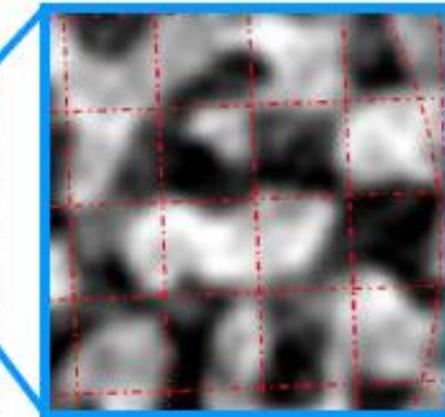
(a)



(b)

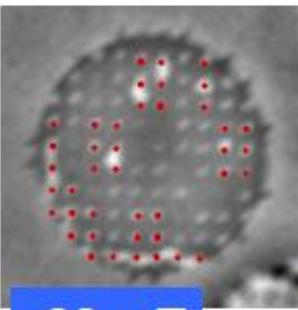


Differenz

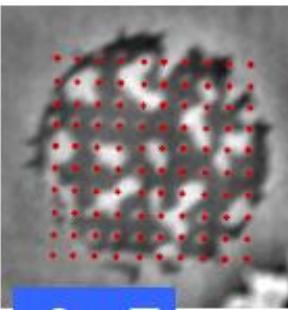


Summe

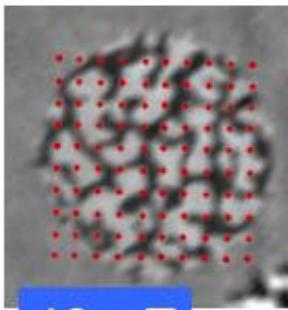
Differenz



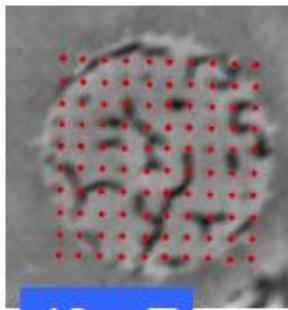
-20 mT



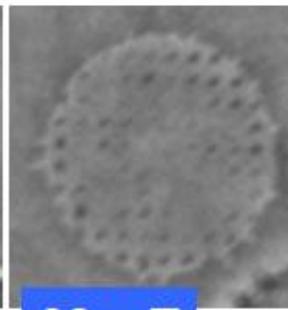
-8 mT



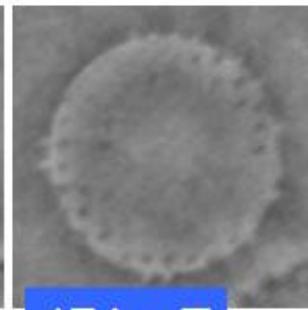
16 mT



40 mT



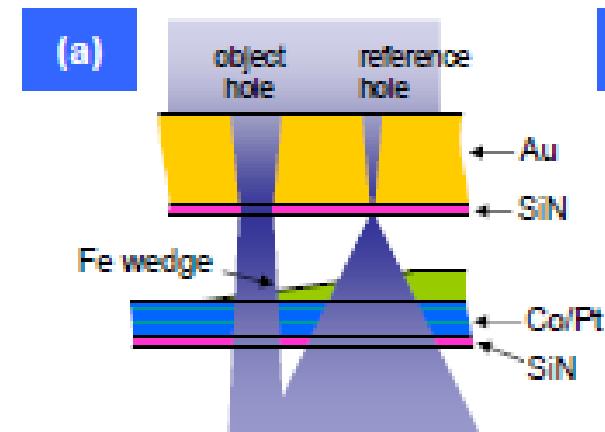
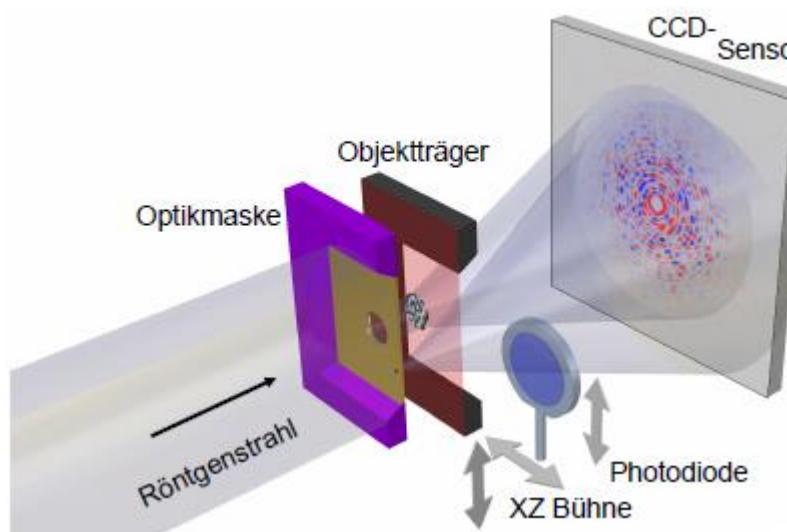
83 mT



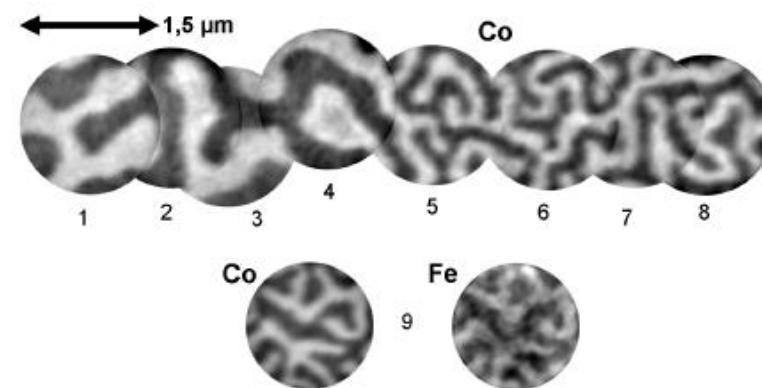
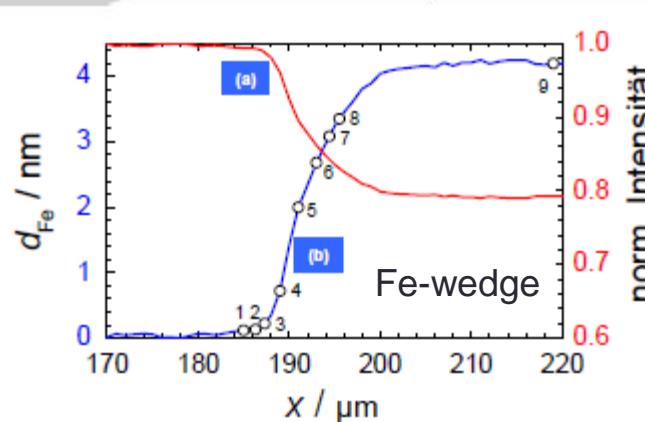
174 mT

Imaging of magnetic domain patterns with X-rays

> Lensless Imaging – Fourier transform Holography (FTH)



D. Stickler et al., Appl. Phys. Lett. **96**, 042501 (2010)



Element-selectivity

