



Molecular Sciences

Introduction of concepts and the basics of molecular physics

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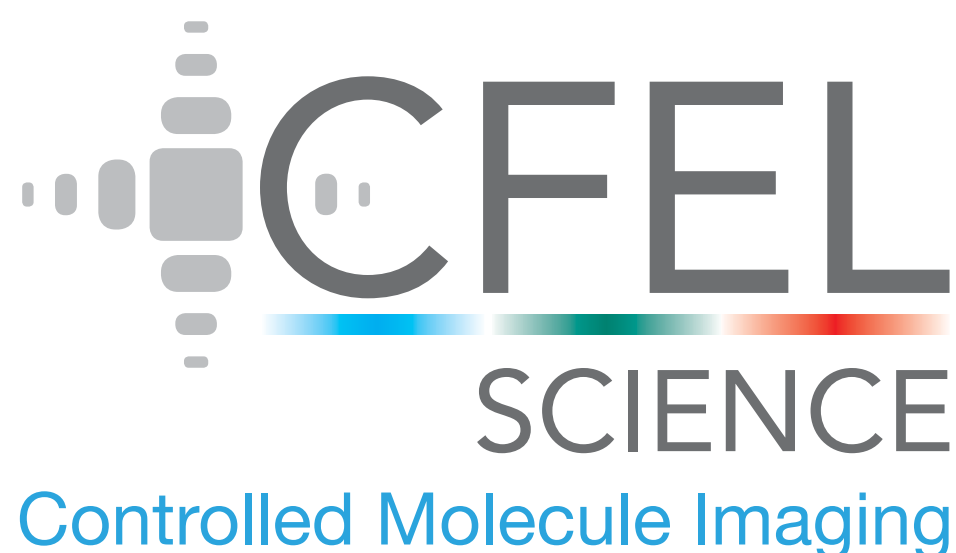
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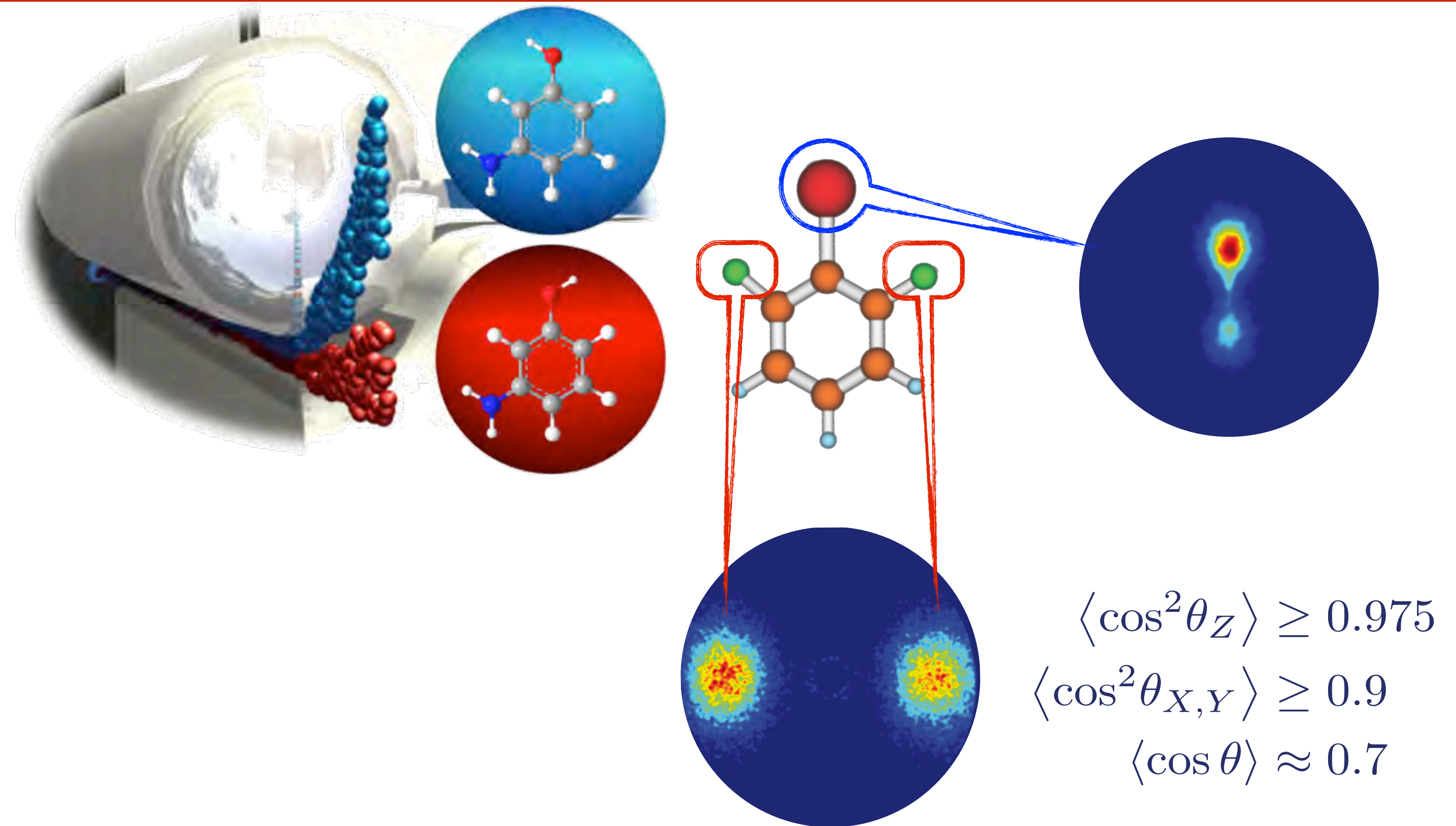
3. & 7. August 2017

based on the UHH master module “The Basis of Modern Molecular Physics”



Molecules in fields

Manipulation of translational and rotational motion

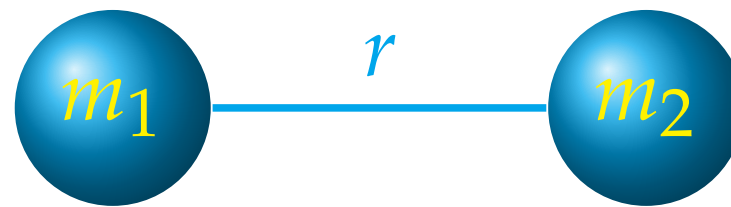


Filsinger, Erlekam, von Helden, JK, Meijer, *Phys. Rev. Lett.* **100**, 133003 (2008)

Holmegaard, Nielsen, Nevo, Stapelfeldt, Filsinger, JK, Meijer, *Phys. Rev. Lett.* **102**, 023001 (2009)

Nevo, Holmegaard, Nielsen, Hansen, Stapelfeldt, Filsinger, Meijer, JK, *Phys. Chem. Chem. Phys.* **11**, 9912 (2009)

Rigid diatomic rotor



- classical energy:

$$E_r = \frac{1}{2} I \omega^2 \quad \text{with } I = m_1 r_1^2 + m_2 r_2^2 \quad (8)$$

with the reduced mass $\mu = \frac{m_1 m_2}{m_1 + m_2}$ and bond length $r = r_1 + r_2$ one obtains

$$E_r = \frac{1}{2} \mu r^2 \omega^2 \quad (9)$$

with $L = I\omega$ this yields

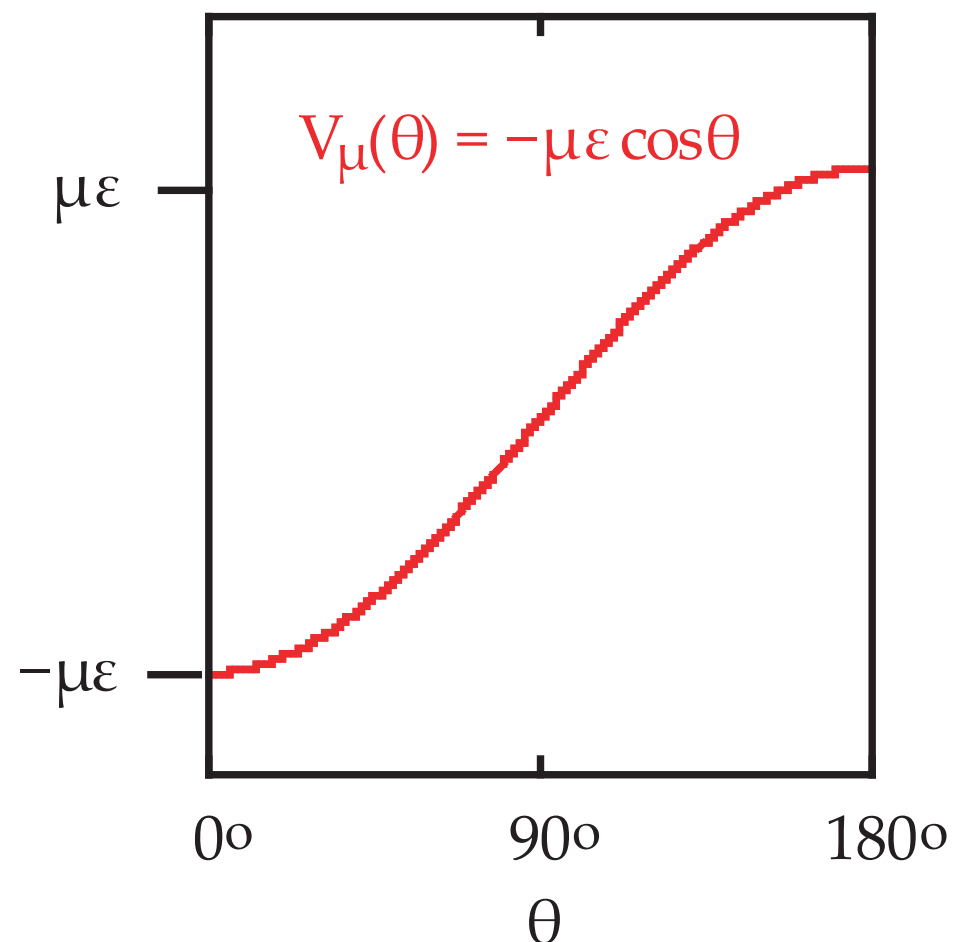
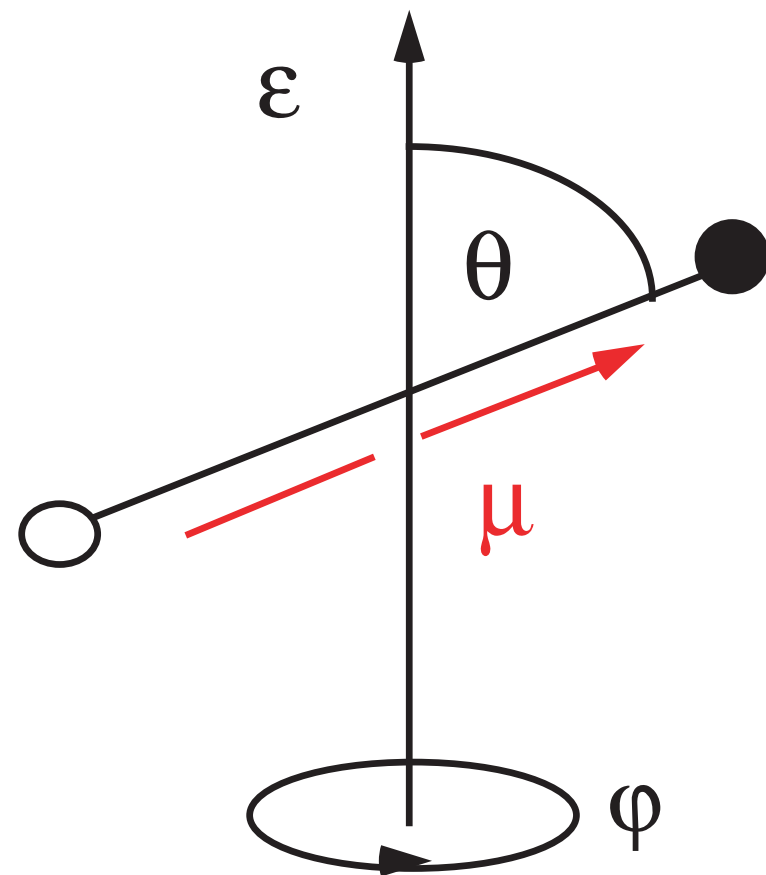
$$E_r = \frac{L^2}{2I} \quad (10)$$

- quantum mechanical energy:

$$E_r = \frac{h^2}{8\pi^2 I} J(J + 1) \quad (11)$$

- E_r is inversely proportional to I
- E_r scales with rotational quantum number as $J(J + 1)$

Orientation of a Permanent Dipole



- Hamiltonian

$$H = B\vec{J}^2 + V_\mu \quad (45)$$

- Dimensionless Hamiltonian

$$\frac{H}{B} = \vec{J}^2 - \frac{\mu\epsilon}{B} \cos \theta = \vec{J}^2 - \omega \cos \theta \quad (46)$$

- Dimensionless *interaction strength* parameter

$$\omega = \frac{\mu\epsilon}{B} \quad (47)$$

- Mixing operator — couples states with equal M but J 's differing by ± 1

$$\cos \theta \quad (48)$$

- Cylindrical symmetry around ϵ , therefore uniform in ϕ .

Schrödinger Equation

- Schrödinger Equation

$$\frac{H}{B} \psi = (\vec{J}^2 - \omega \cos \theta) \psi = \frac{E}{B} \psi \quad (49)$$

- Representation in the free rotor basis $|JK\rangle$

$$\psi = |\tilde{J}M\omega\rangle = \sum_{J=M}^{\infty} a_J^{\tilde{J}M}(\omega) |JM\rangle \quad (50)$$

- Matrix elements

$$\langle JM | \vec{J} | JM \rangle = J(J+1) \quad (51)$$

$$\langle JM | \cos \theta | J+1M \rangle = \sqrt{\frac{(J+M+1)(J-M+1)}{(2J+1)(2J+3)}} \quad (52)$$

$$\langle JM | \cos \theta | JM \rangle = 0 \quad (53)$$

$$\langle JM | \cos \theta | J-1M \rangle = \sqrt{\frac{(J+M)(J-M)}{(2J-1)(2J+1)}} \quad (54)$$

Energies and Wavefunctions

- Eigenfunctions are coherent linear superpositions of the field-free rotor basis functions: *hybridization*.
- Wavefunctions are of mixed parity (i. e., parity is different from ± 1).
With equation (50) and parity of basis functions $p_{J,M} = (-1)^J = p_J$, the parity of the mixed wavefunction is mixed as well:

$$p = \sum_J a_J p_J \quad (55)$$

- Parity conservation does not apply to state of indefinite parity. Such a state can be oriented and have a space-fixed electric dipole moment.
- In a state of definite parity, the space-fixed electric dipole moment is zero.

Example calculation

Group work “pendular states” (teams of 2, 15 min)

- set up the Hamiltonian matrix for fixed M and ω
- ensure convergence by using an appropriately sized Hamiltonian matrix
- okay, let's use $M = 0$ and $\omega = 1$, dimension (3×3) and the known matrix elements

$$H_{J,J+1} = \langle JM | \cos \theta | J+1M \rangle = \sqrt{\frac{(J+M+1)(J-M+1)}{(2J+1)(2J+3)}}$$

$$\langle JM | \cos \theta | JM \rangle = 0$$

$$H_{J,J-1} = \langle JM | \cos \theta | J-1M \rangle = \sqrt{\frac{(J+M)(J-M)}{(2J-1)(2J+1)}}$$

$$H_{J,J} = \langle JM | \vec{J} | JM \rangle = J(J+1)$$

- What can you say about the eigenstates under these conditions (numbers)?
- What is the energy of the ground state?
- What is the wavefunction of the ground state?
- What is the parity of the ground state?

Example Calculation I

Solution

- set up a matrix for fixed M and ω and a dimension that ensures convergence of eigenproperties obtained by diagonalization

$$\hat{H} = \begin{pmatrix} J_1(J_1 + 1) & -\omega \sqrt{\frac{(J_1+M+1)(J_1-M+1)}{(2J_1+1)(2J_1+3)}} & 0 \\ -\omega \sqrt{\frac{(J_2+M)(J_2-M)}{(2J_2-1)(2J_2+1)}} & J_2(J_2 + 1) & -\omega \sqrt{\frac{(J_2+M+1)(J_2-M+1)}{(2J_2+1)(2J_2+3)}} \\ 0 & -\omega \sqrt{\frac{(J_3+M)(J_3-M)}{(2J_3-1)(2J_3+1)}} & J_3(J_3 + 1) \end{pmatrix}$$

- for instance, for $M = 0$ and $\omega = 1$ get

$$\hat{H} = \begin{pmatrix} 0(0 + 1) & -1 \sqrt{\frac{(0+0+1)(0-0+1)}{(2 \cdot 0+1)(2 \cdot 0+3)}} & 0 \\ -1 \sqrt{\frac{(1+0)(1-0)}{(2 \cdot 1-1)(2 \cdot 1+1)}} & 1(1 + 1) & -1 \sqrt{\frac{(1+0)(1-0)}{(2 \cdot 1+1)(2 \cdot 1+3)}} \\ 0 & -1 \sqrt{\frac{(2+0)(2-0)}{(2 \cdot 2-1)(2 \cdot 2+1)}} & 2(2 + 1) \end{pmatrix}$$

Example Calculation II

Solution

- this yields the (numerical) Hamiltonian matrix

$$\hat{H} = \begin{pmatrix} 0 & -0.57735 & 0 \\ -0.57735 & 2 & -0.516398 \\ 0 & -0.516398 & 6 \end{pmatrix}$$

- Diagonalization yields the eigenvalues and eigenvectors

$$\lambda_1 = -0.157653 \quad C_1 = (0.964446, 0.2639998, 0.01215001) \quad (56)$$

$$\lambda_2 = 2.09118 \quad C_2 = (-0.264, 0.956214, 0.126326) \quad (57)$$

$$\lambda_3 = 6.06648 \quad C_3 = (0.01215, -0.127666, 0.991743) \quad (58)$$

- This calculation defines the ground state with sufficient accuracy:

$$E(\tilde{J} = 0, M = 0)_{\omega=1} = -0.157653 \cdot B \quad (59)$$

$$\begin{aligned} |\tilde{J} = 0, M = 0\rangle_{\omega=1} = & 0.964446 |J = 0, M = 0\rangle + 0.2639998 |J = 1, M = 0\rangle + \\ & 0.01215001 |J = 1, M = 0\rangle \end{aligned} \quad (60)$$

- The eigenfunctions of the Hamiltonian are the coherent linear superpositions of the field-free-rotor basis functions

Example Calculation III

Solution

- we speak of *hybridization* of the wavefunctions (basis functions)
- For our ground state we get the parity

$$\begin{aligned} p &= \sum_J a_j p_j = 0.964446 \cdot (-1)^0 + 0.2639998 \cdot (-1)^1 + 0.01215001 \cdot (-1)^2 & (61) \\ &= 0.712 \\ &\neq \pm 1 \end{aligned}$$

Since the parity is not defined, such a state can be (is) oriented.
The system in this state has a space-fixed dipole moment!

Effective Dipole Moment

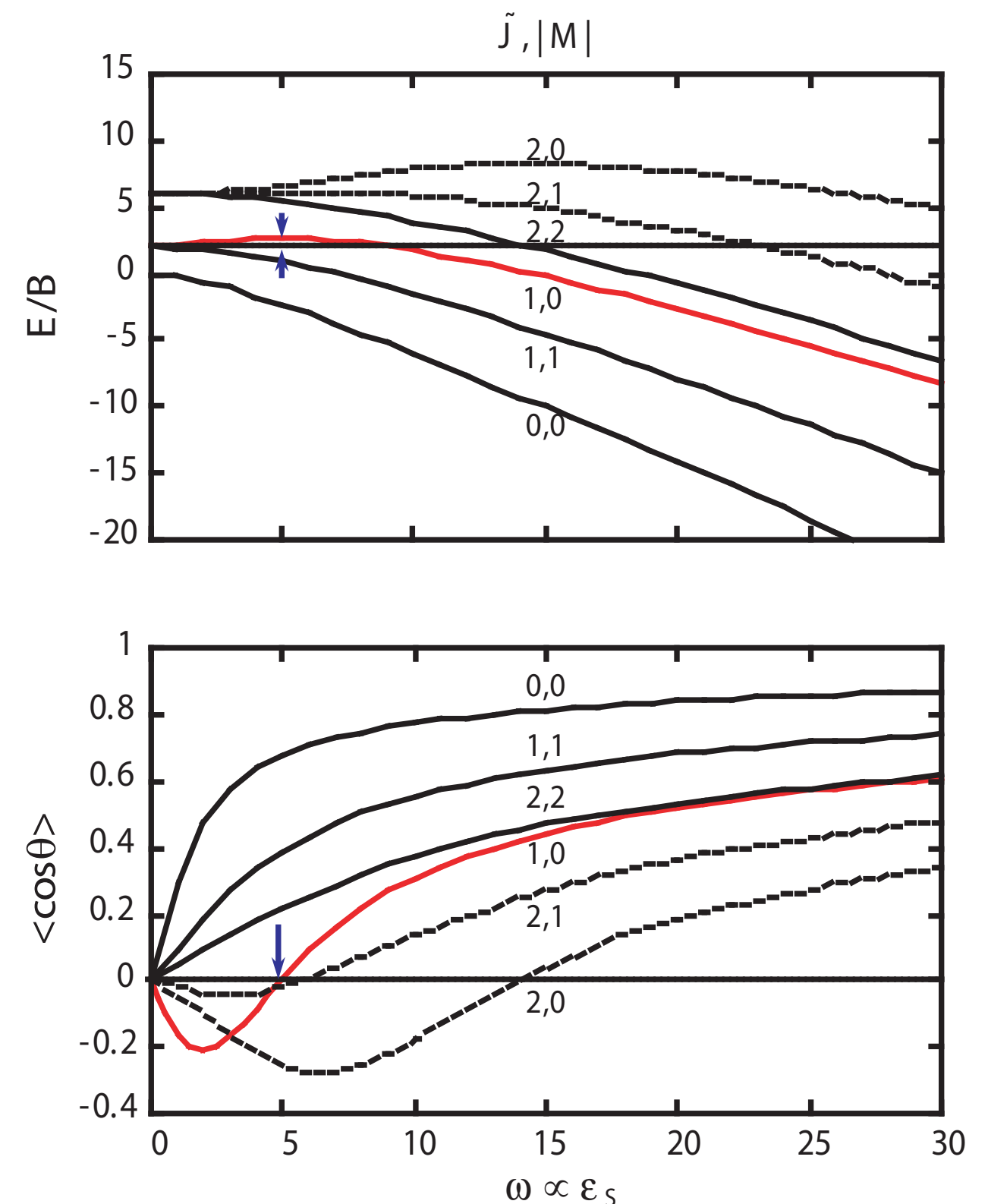
The (negative) derivative of the energy with respect to the field strength is the effective dipole moment of the quantum state:

$$\langle \cos \theta \rangle = -\frac{\partial E/B}{\partial \omega} \quad (62)$$

$$\theta_{\max} = \arccos \langle \cos \theta \rangle \quad (63)$$

low field seeker E increases with increasing ϵ .

high field seeker E decreases with increasing ϵ .



Molecules in fields

Selection of quantum states and species

$$U_{Stark} = -\frac{1}{2}E^2(\Delta\alpha\cos^2\theta + \alpha_{\perp}) - \mu E\cos\theta$$

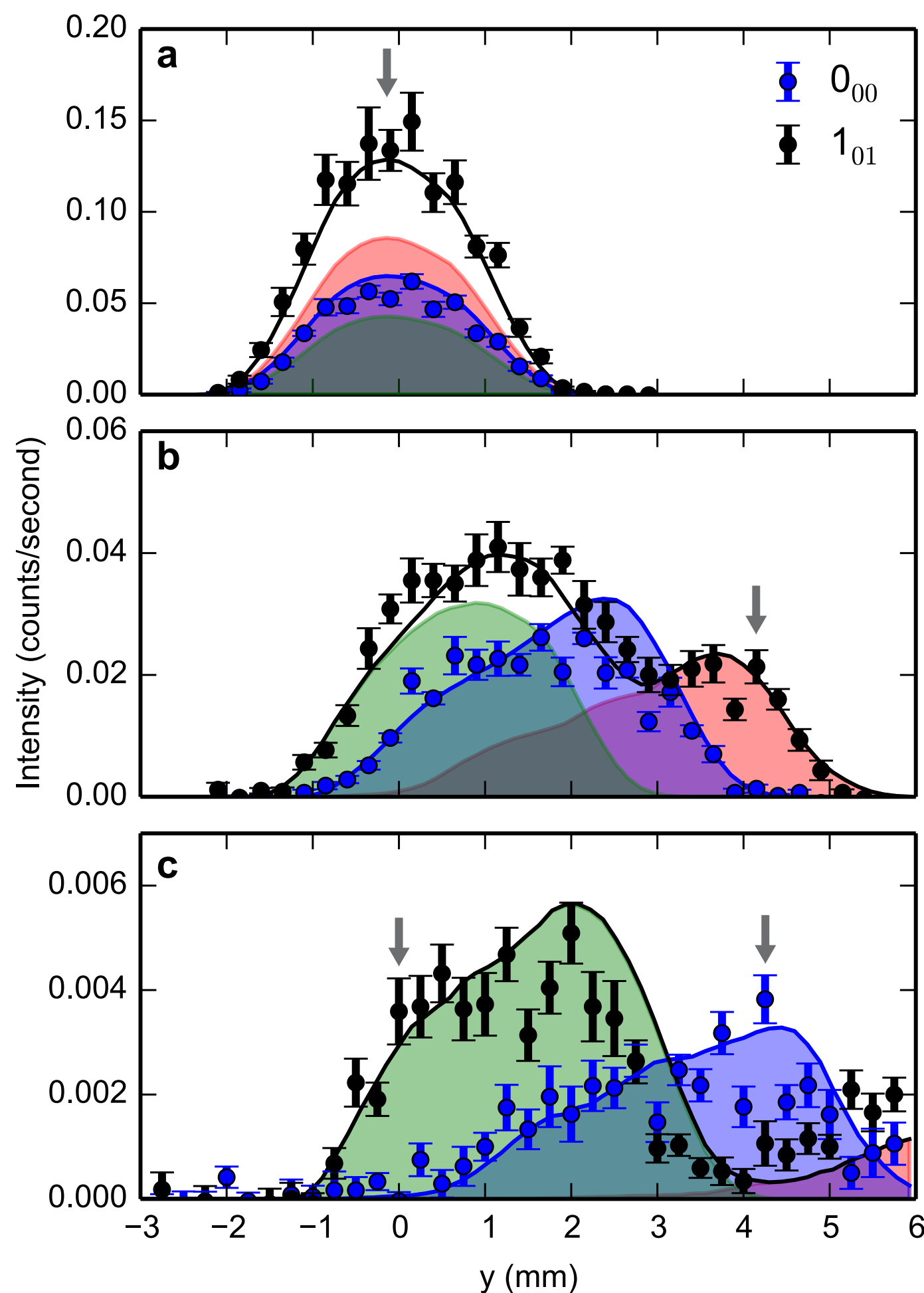
$$U_{Stark} = -\vec{\mu} \cdot \vec{E}$$

$$\vec{F}_{Stark} = -\vec{\nabla}U_{Stark}$$

$$f(\mu, \vec{\nabla}E) \longrightarrow v_z = -\frac{1}{M} \int_{-\infty}^{\infty} \left(\vec{\nabla}U \right)_z dt,$$

Understanding and controlling the dynamics of highly excited molecular systems

Separating para and ortho water

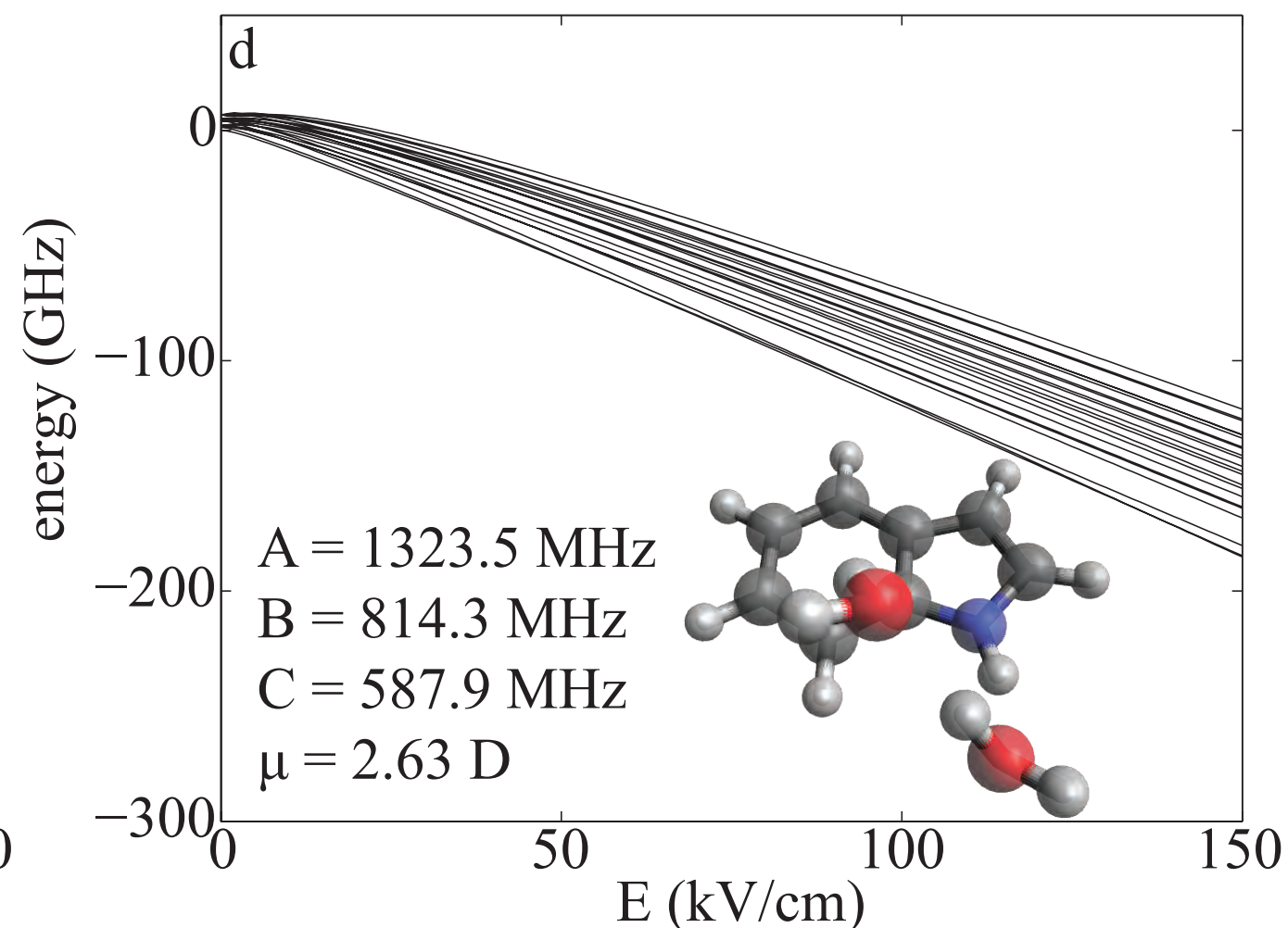
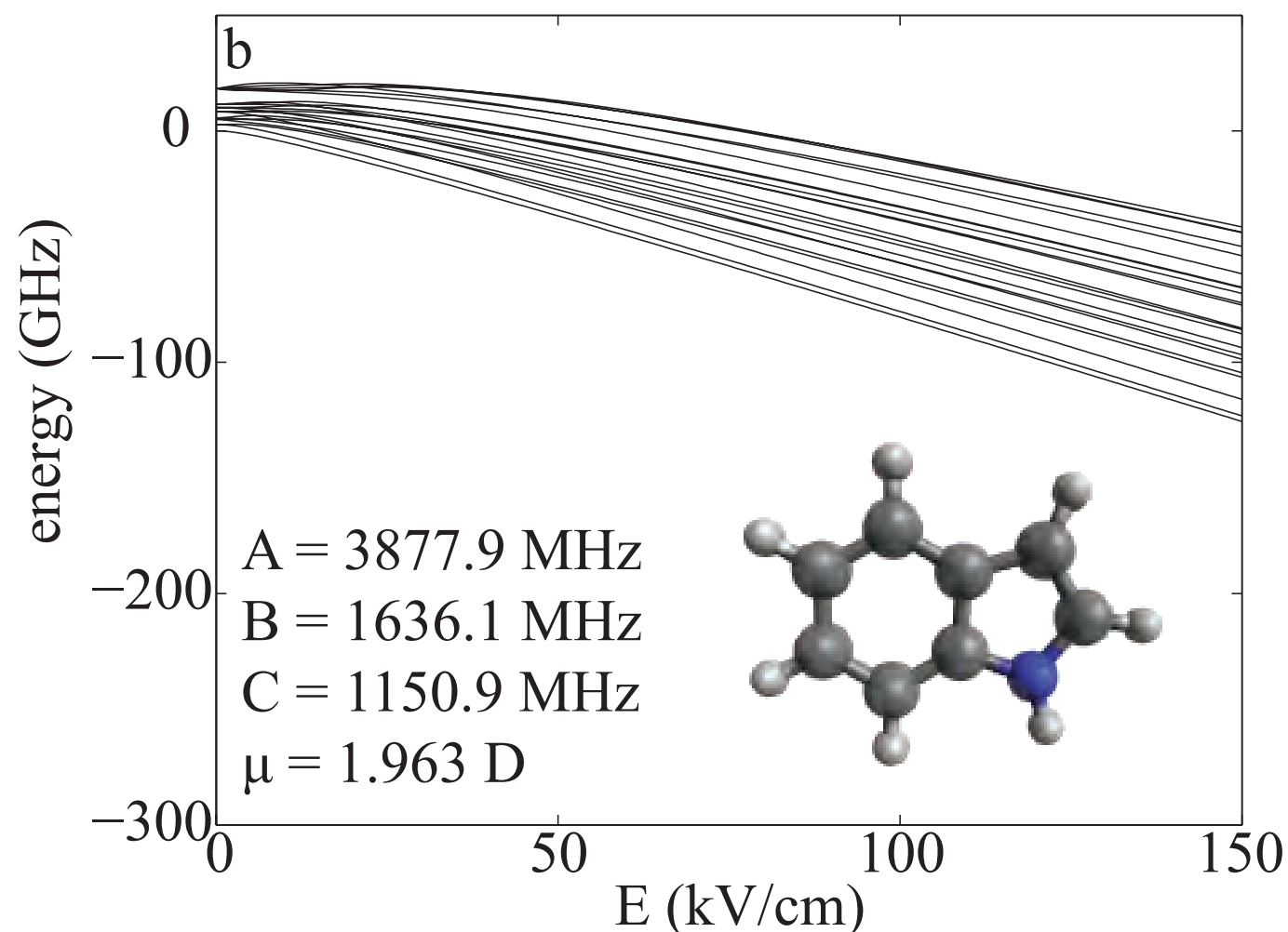
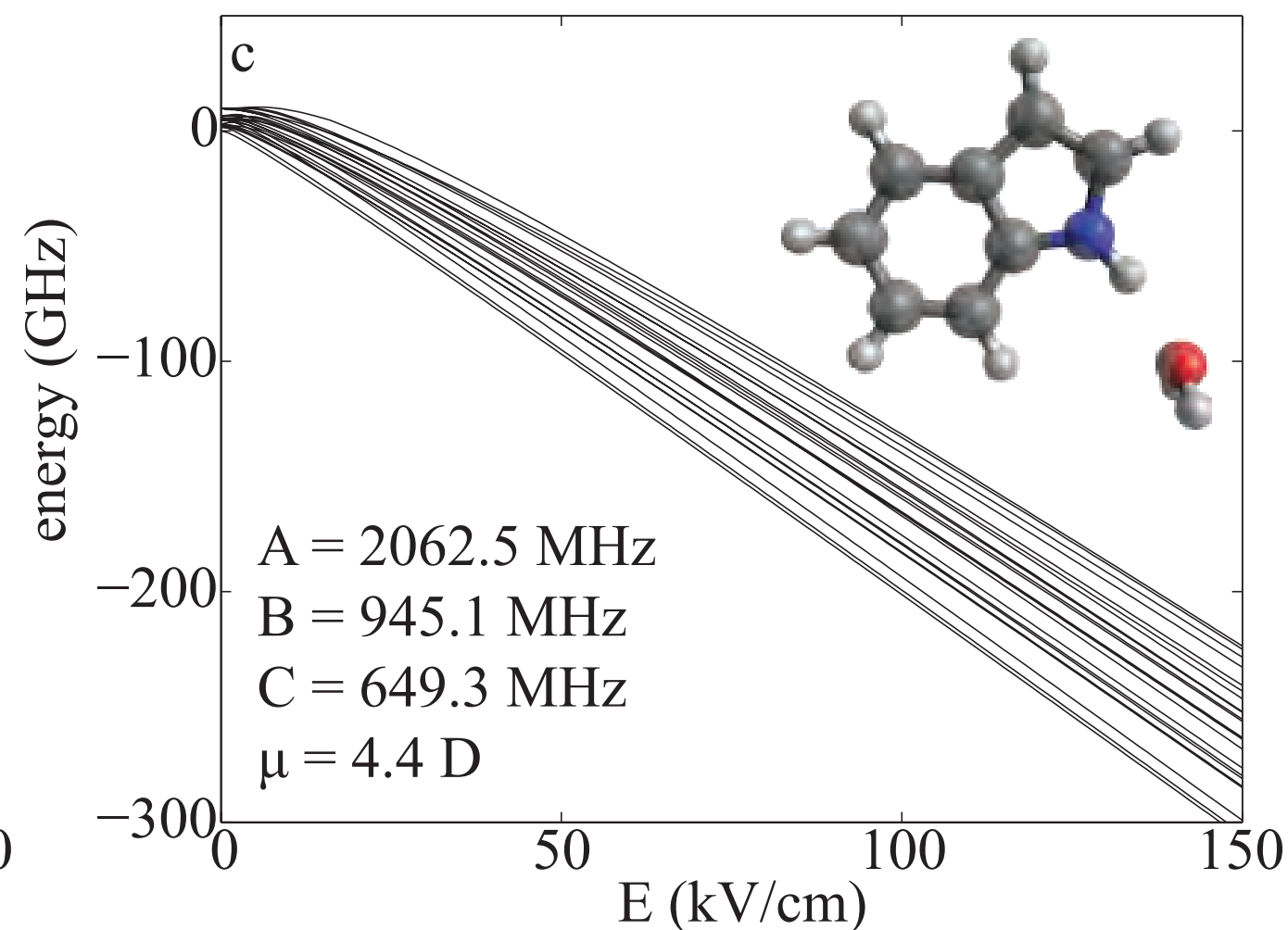
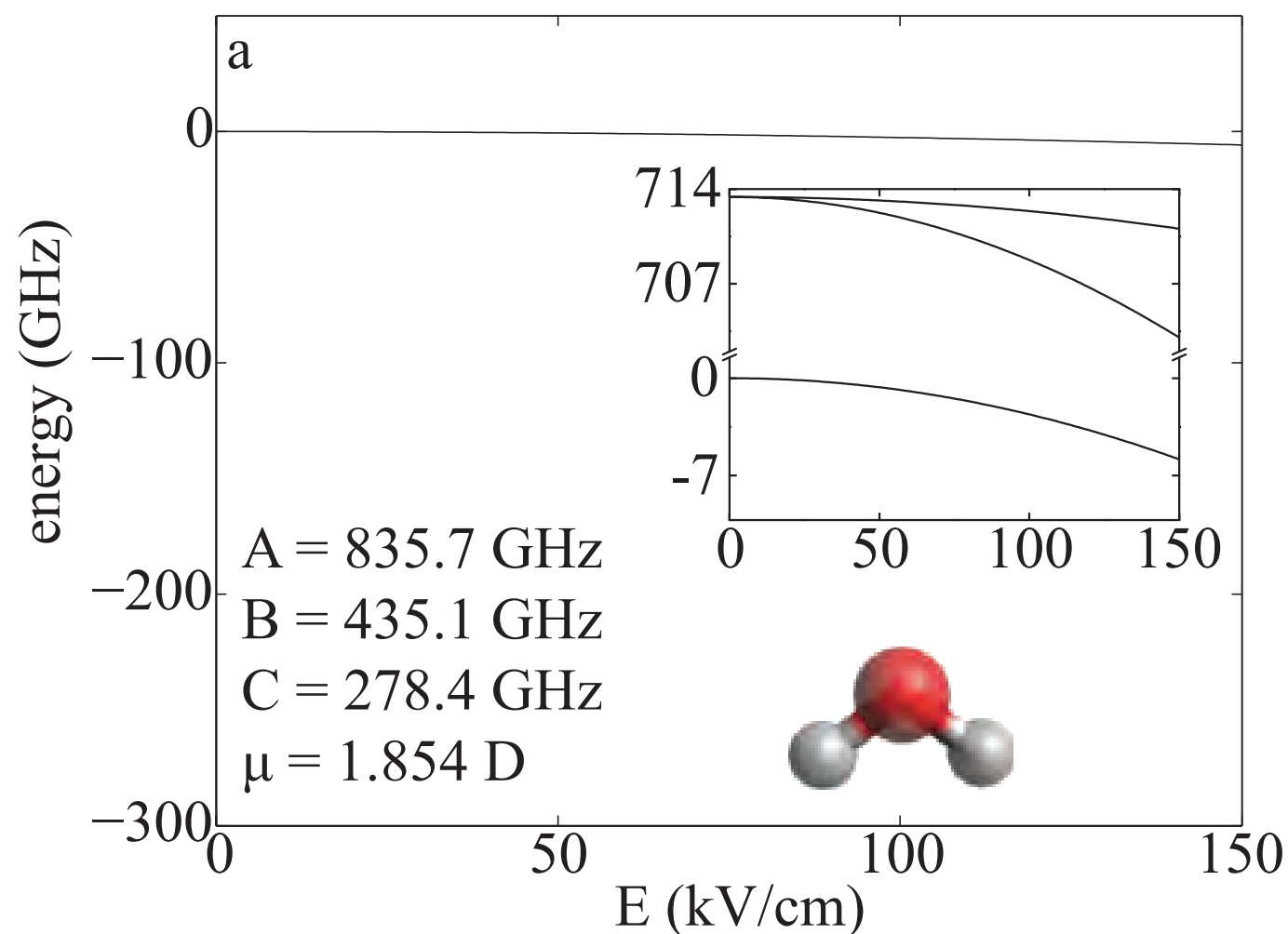


$$\vec{F}_{Stark} = 0$$

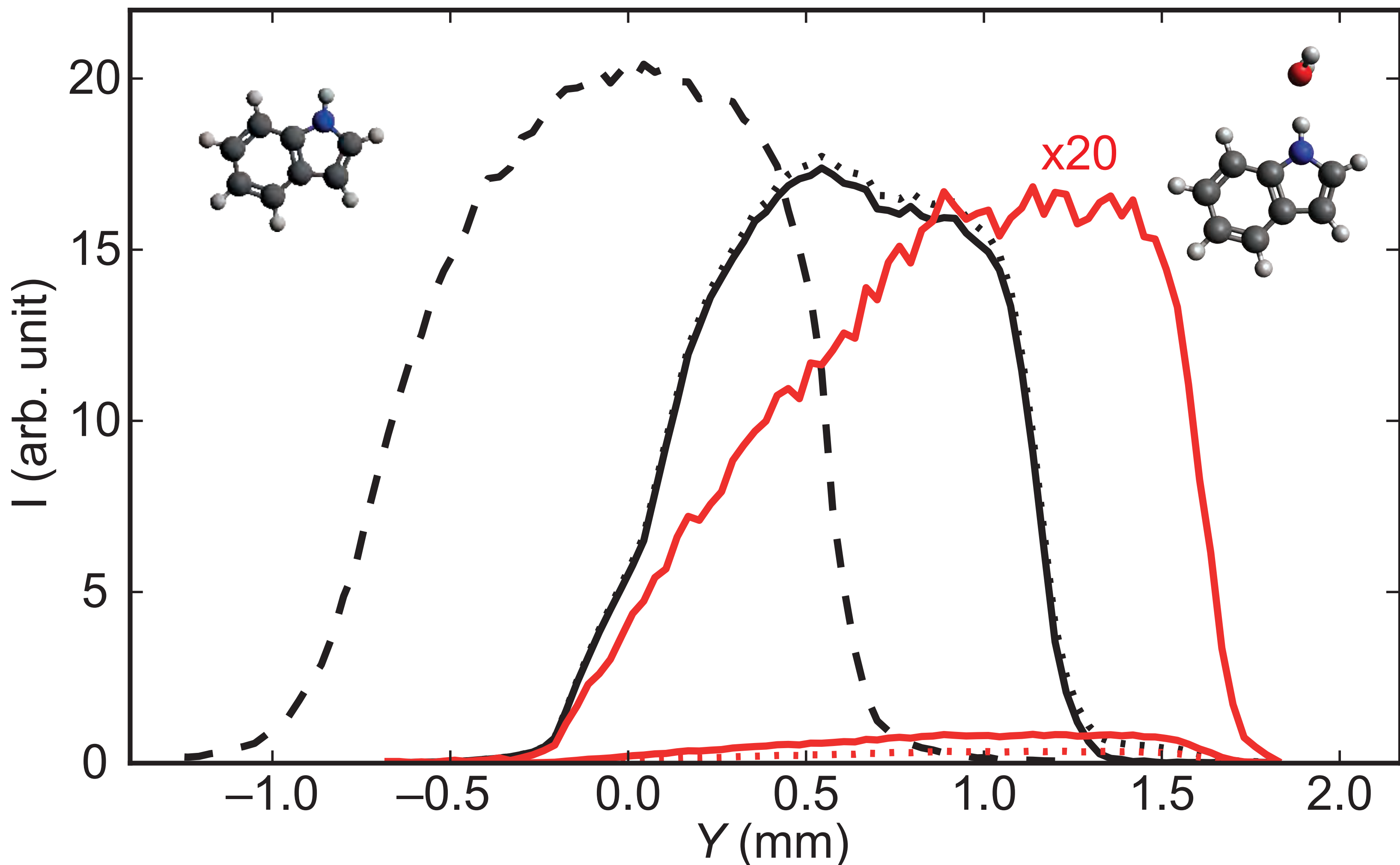
$$\vec{F}_{Stark1}$$

$$\vec{F}_{Stark2}$$

Spatial separation of neutral clusters using the m/μ deflector pure samples of indole-water ($\text{indole}-(\text{H}_2\text{O})_1$)



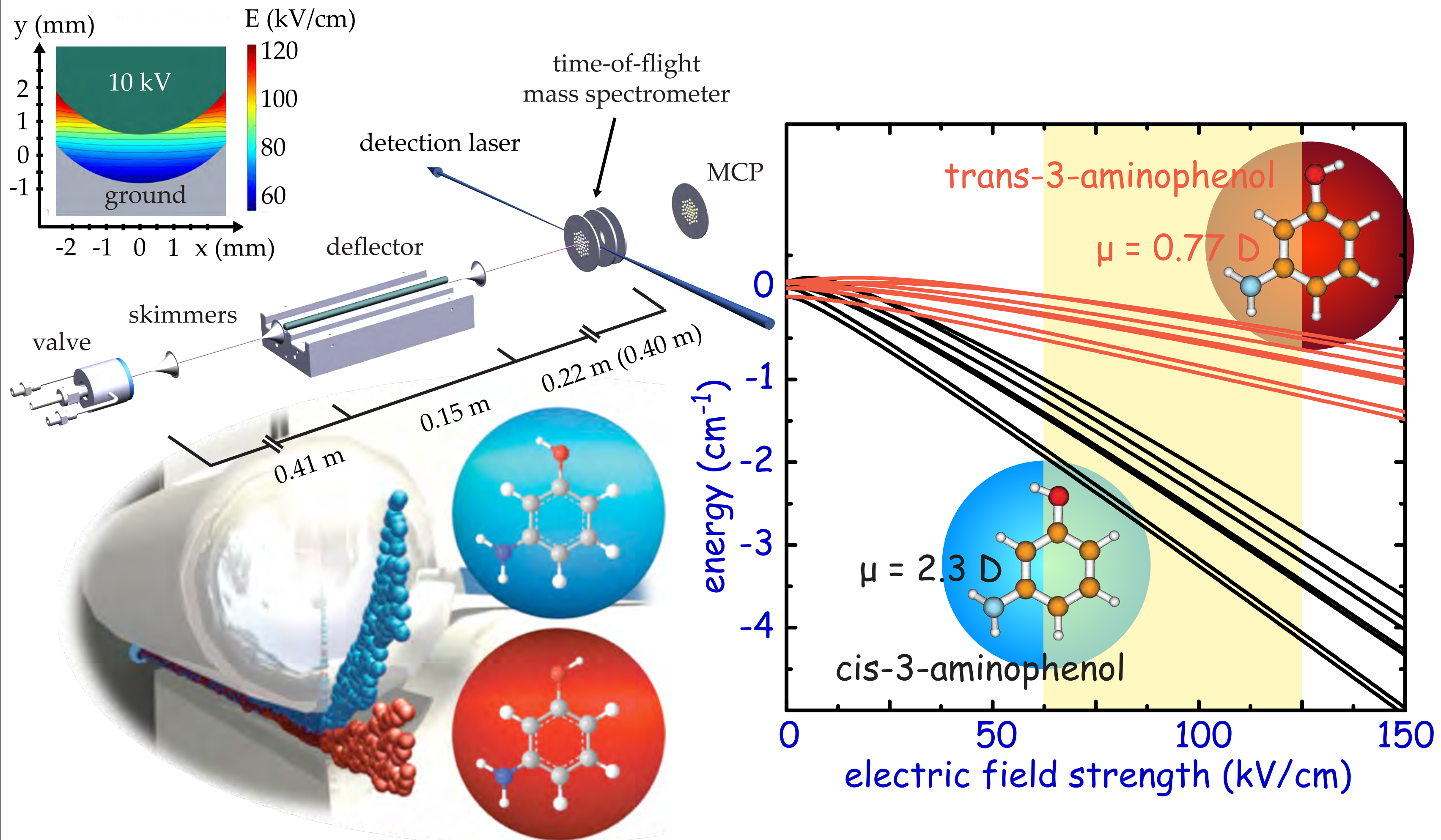
Spatial separation of neutral clusters using the m/μ deflector pure samples of indole-water ($\text{indole}-(\text{H}_2\text{O})_1$)



Trippel, Chang, Stern, Mullins, Holmegaard, JK, Phys. Rev. A 86, 033202 (2012)

Chang, Horke, Trippel, JK, Int. Rev. Phys. Chem. 34, 557–590 (2015)

Conformer selection with the m/μ deflector



Filsinger, Erlekam, von Helden, JK, Meijer, *Phys. Rev. Lett.* **100**, 133003 (2008)

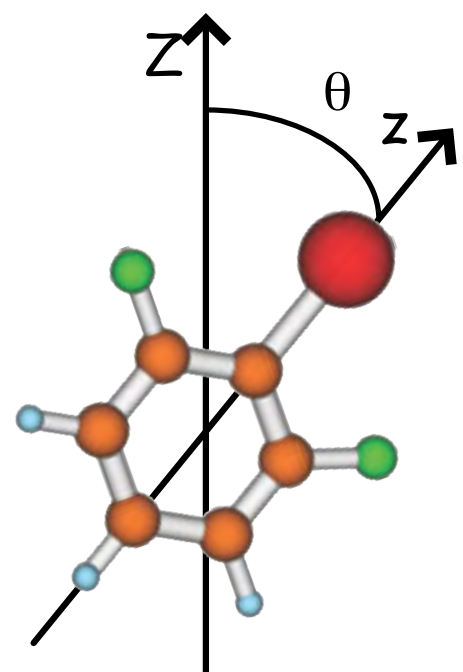
Filsinger, JK, Meijer, Hansen, Maurer, Nielsen, Holmegaard, Stapelfeldt, *Angew. Chem. Int. Ed.* **48**, 6900 (2009)

Molecules in fields

Alignment and orientation with electric fields

Molecular alignment and orientation

Connecting molecular and laboratory frame

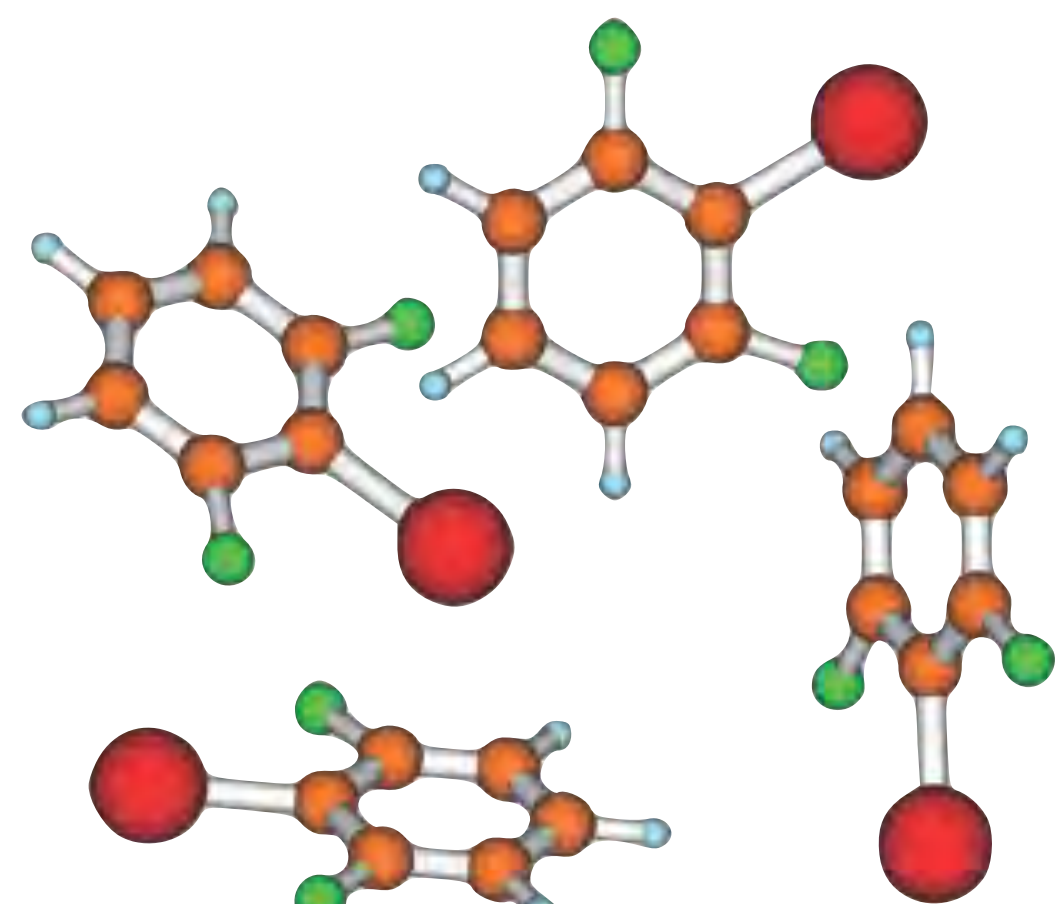


$\langle \cos \theta \rangle$ orientation
 $\langle \cos^2 \theta \rangle$ alignment

isotropic

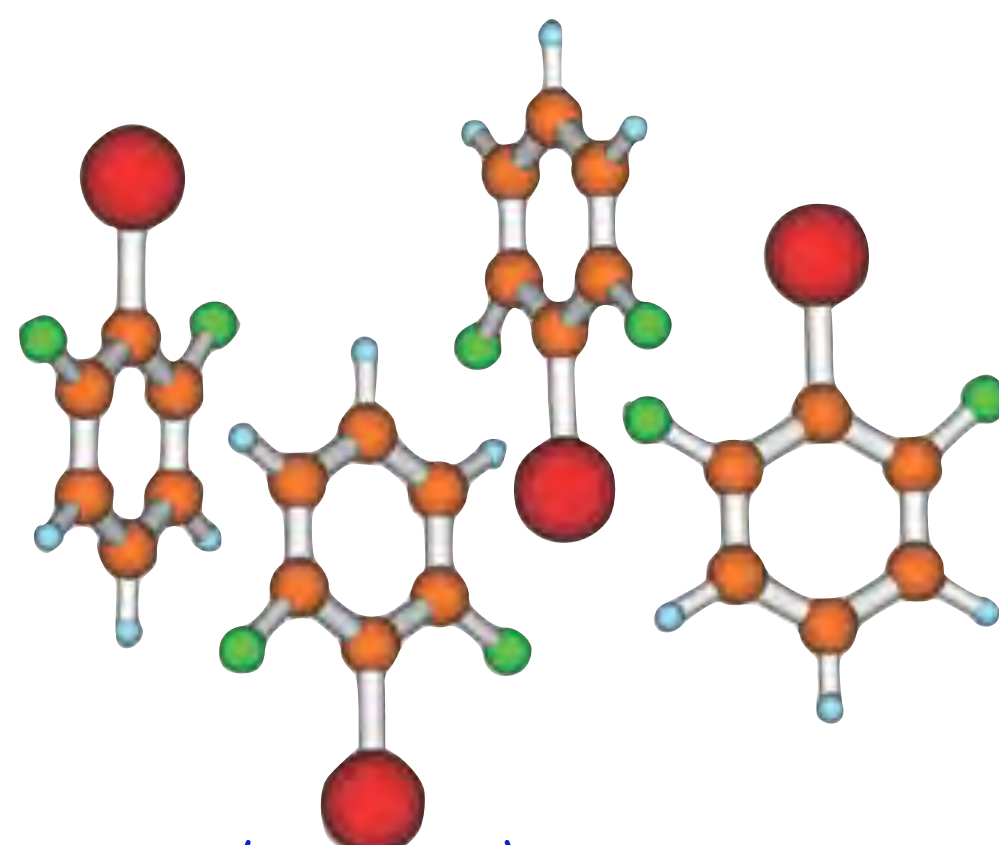
aligned

oriented



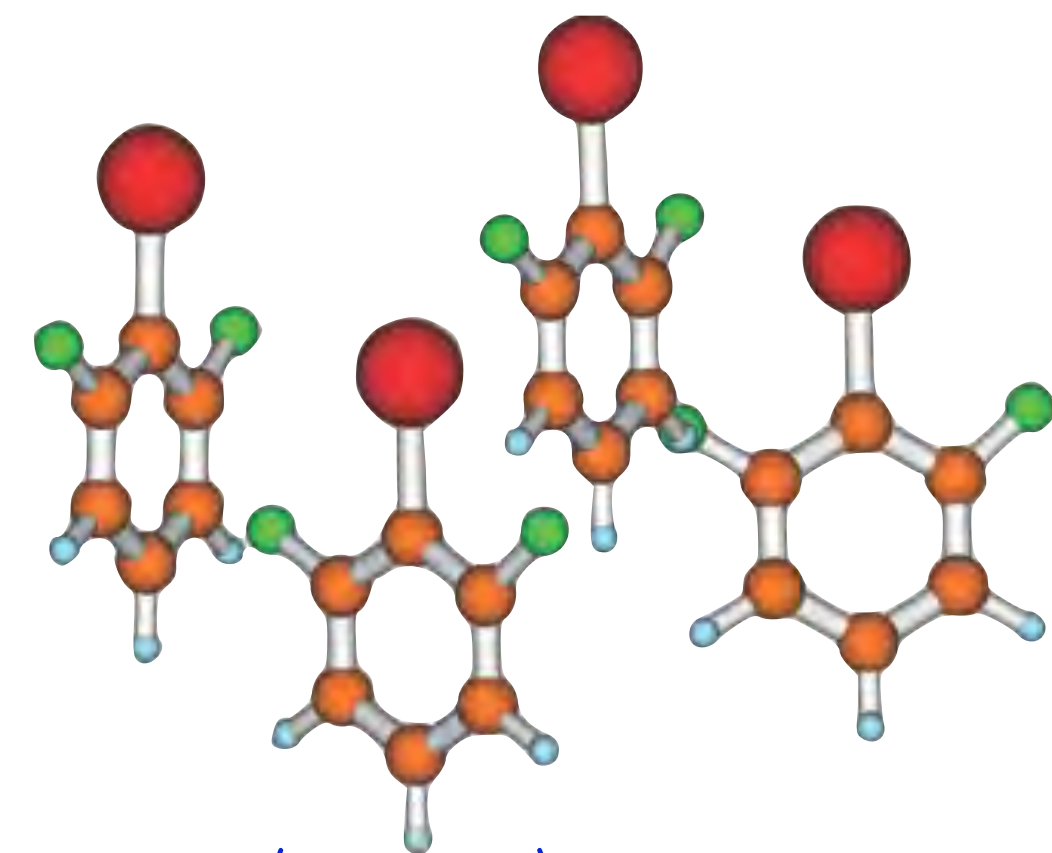
$$\langle \cos \theta \rangle = 0$$

$$\langle \cos^2 \theta \rangle = 1/3$$



$$\langle \cos \theta \rangle = 0$$

$$\langle \cos^2 \theta \rangle > 1/3$$

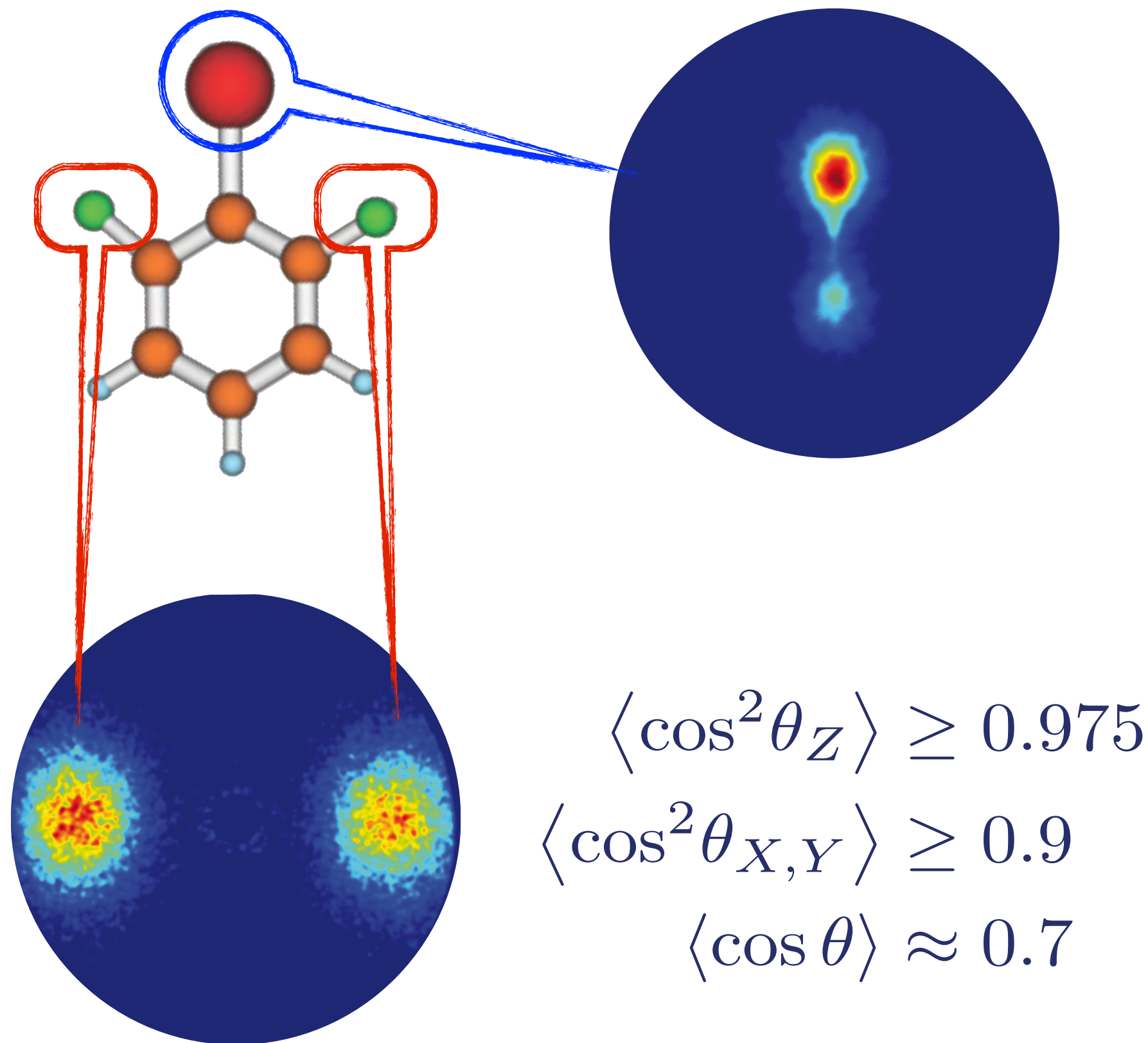


$$\langle \cos \theta \rangle > 0$$

$$\langle \cos^2 \theta \rangle > 1/3$$

Molecular alignment and orientation

Connecting molecular and laboratory frame

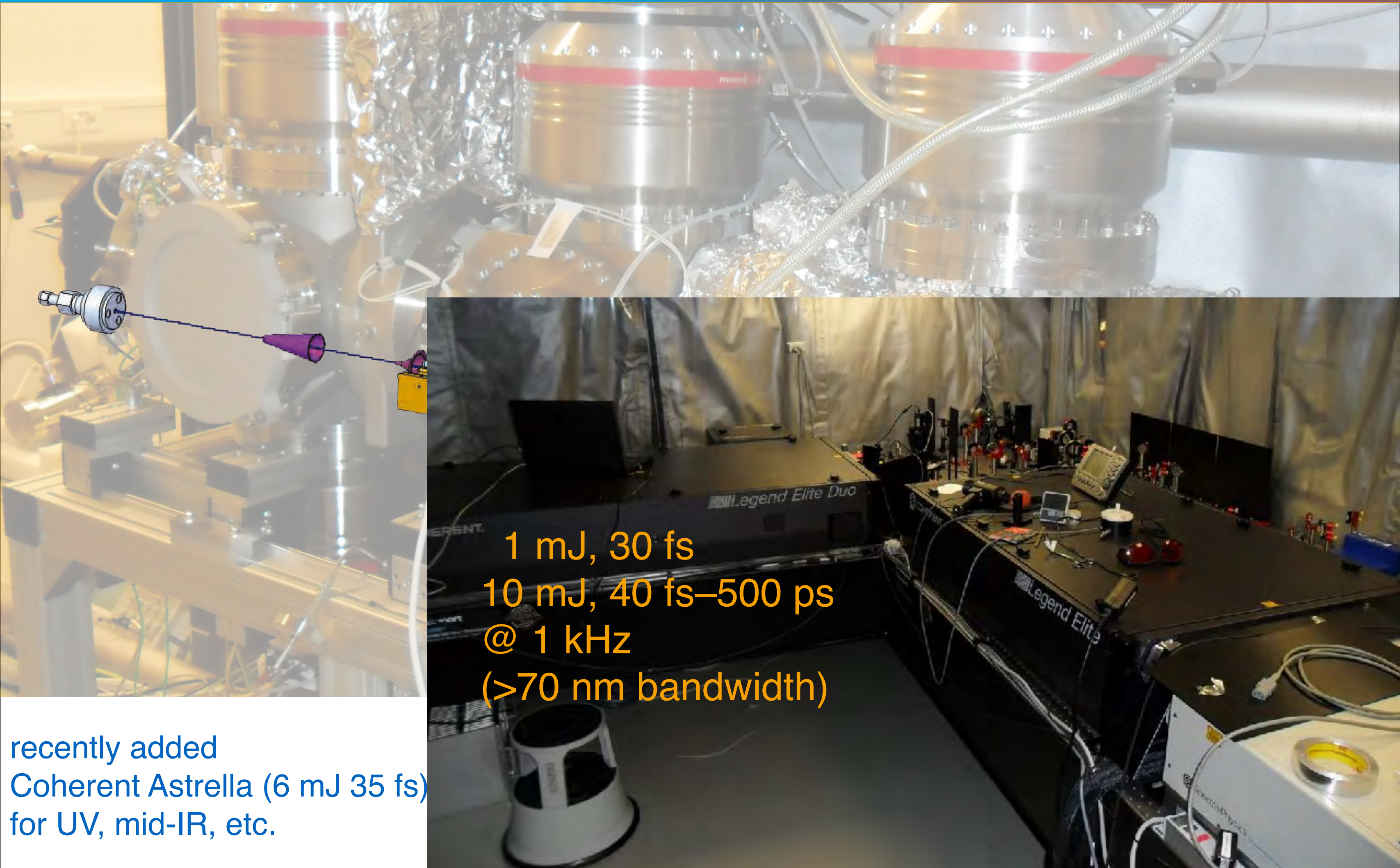


Filsinger, Erlekam, von Helden, JK, Meijer, *Phys. Rev. Lett.* **100**, 133003 (2008)

Holmegaard, Nielsen, Nevo, Stapelfeldt, Filsinger, JK, Meijer, *Phys. Rev. Lett.* **102**, 023001 (2009)

Nevo, Holmegaard, Nielsen, Hansen, Stapelfeldt, Filsinger, Meijer, JK, *Phys. Chem. Chem. Phys.* **11**, 9912 (2009)

Toward time-resolved *imaging of chemical dynamics* kHz-rate manipulation experiments

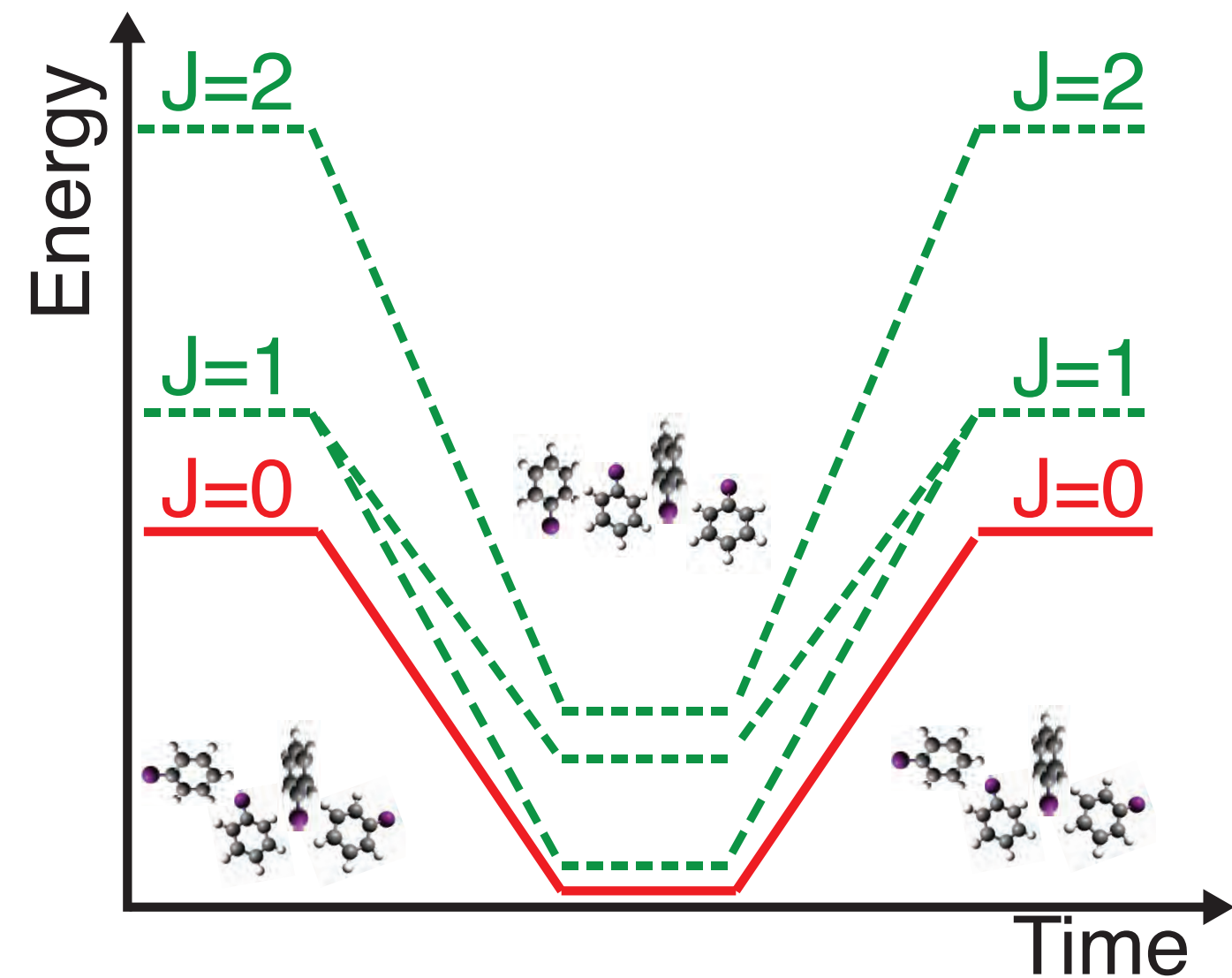
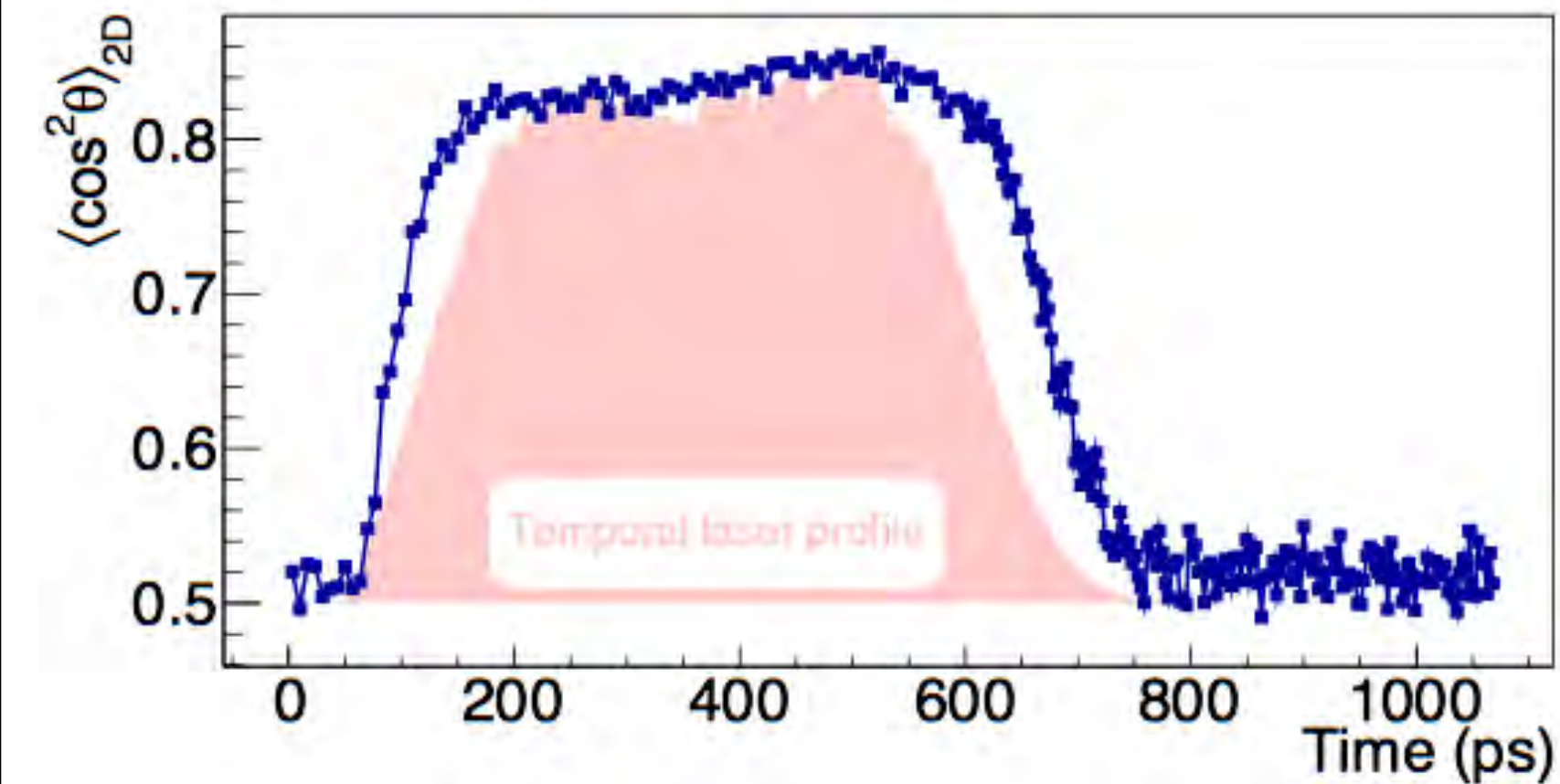
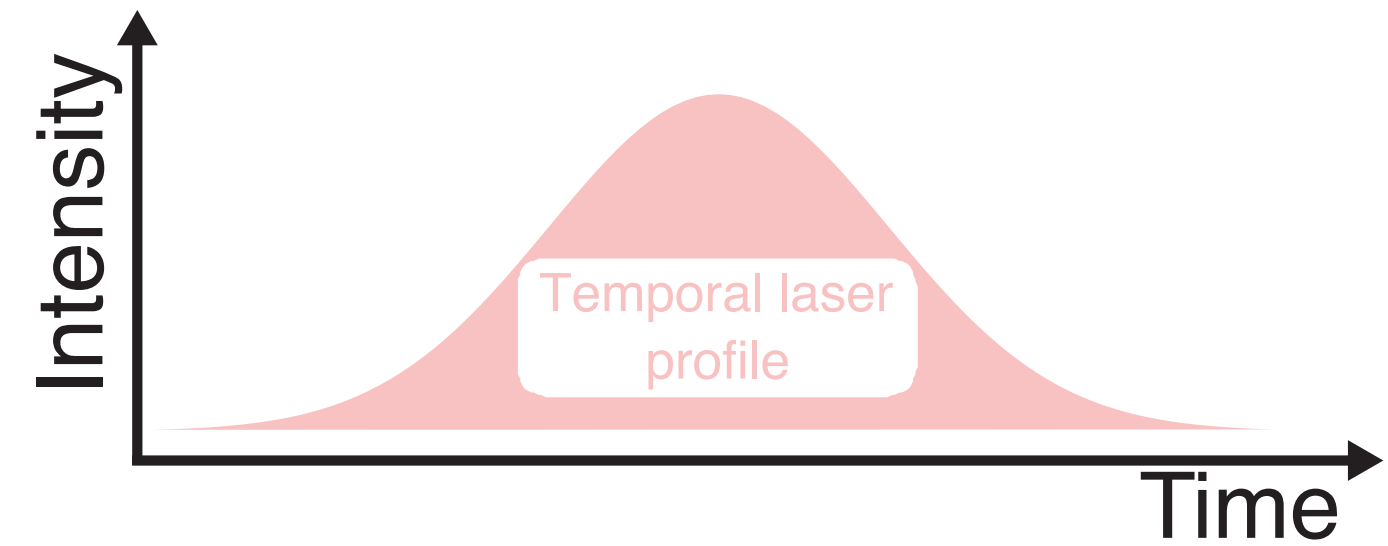
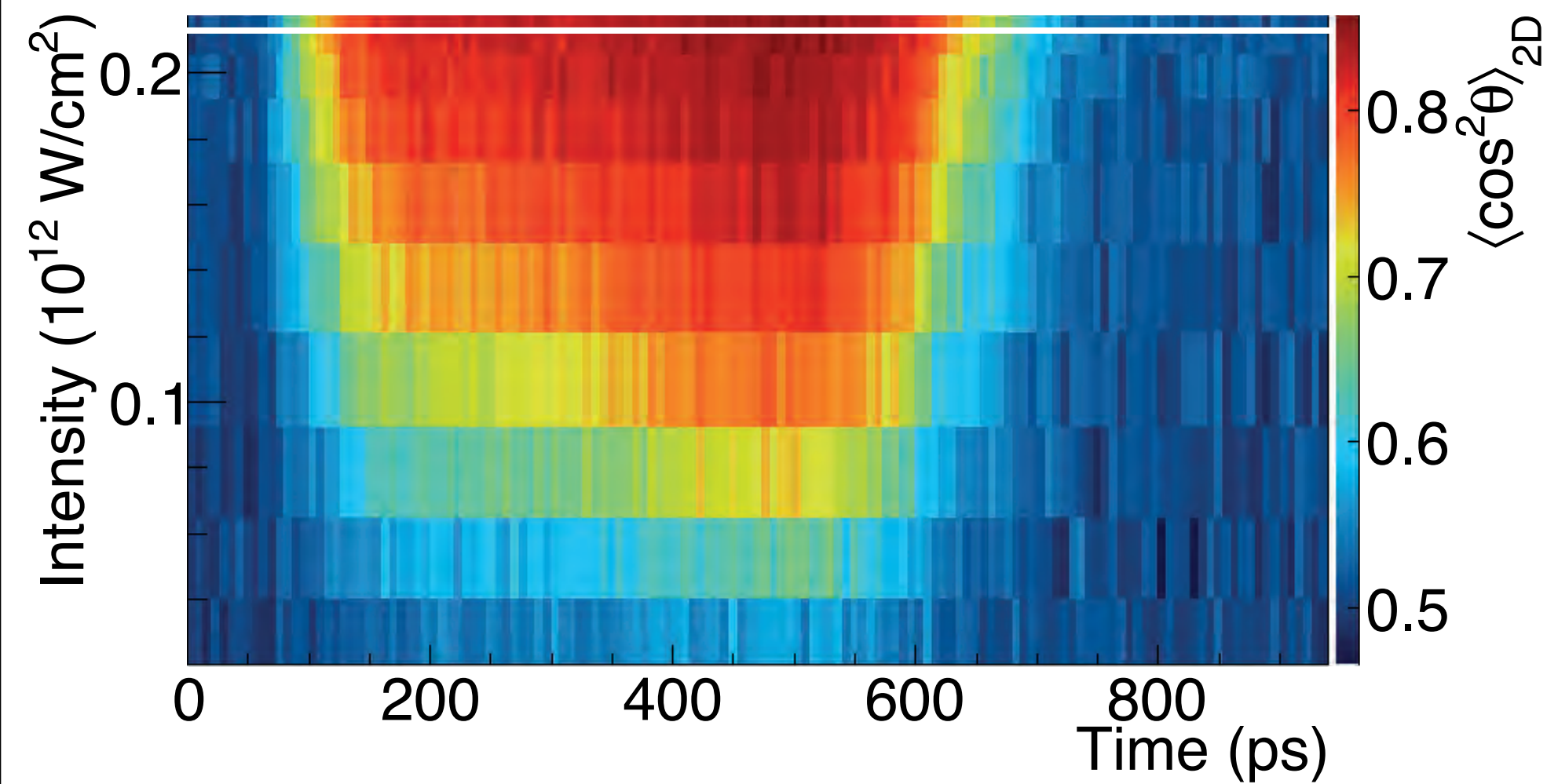


1 mJ, 30 fs
10 mJ, 40 fs–500 ps
@ 1 kHz
(>70 nm bandwidth)

recently added
Coherent Astrella (6 mJ 35 fs)
for UV, mid-IR, etc.

Scenarios of rotational dynamics in OCS ($X, v=0, J=0$)

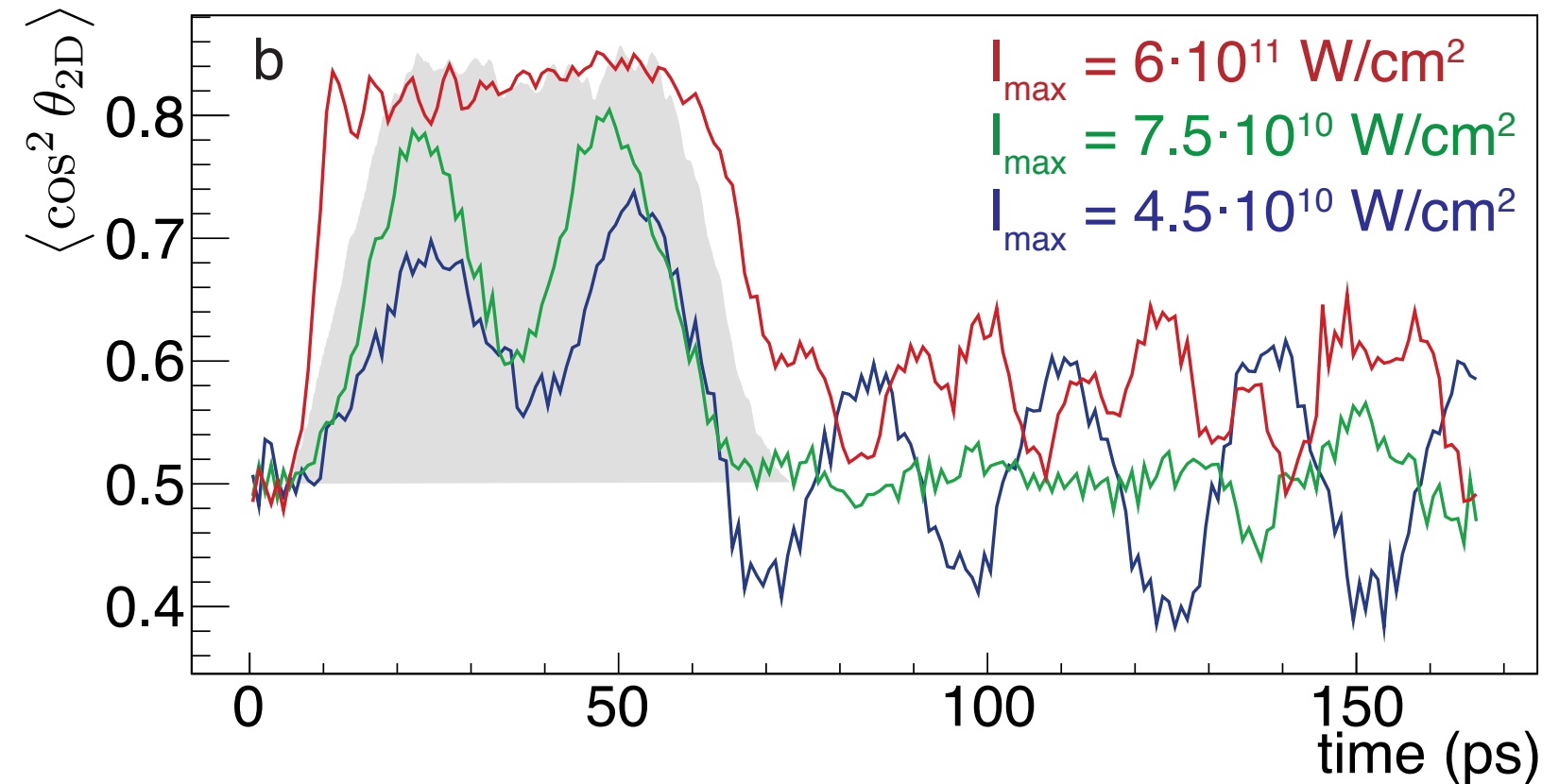
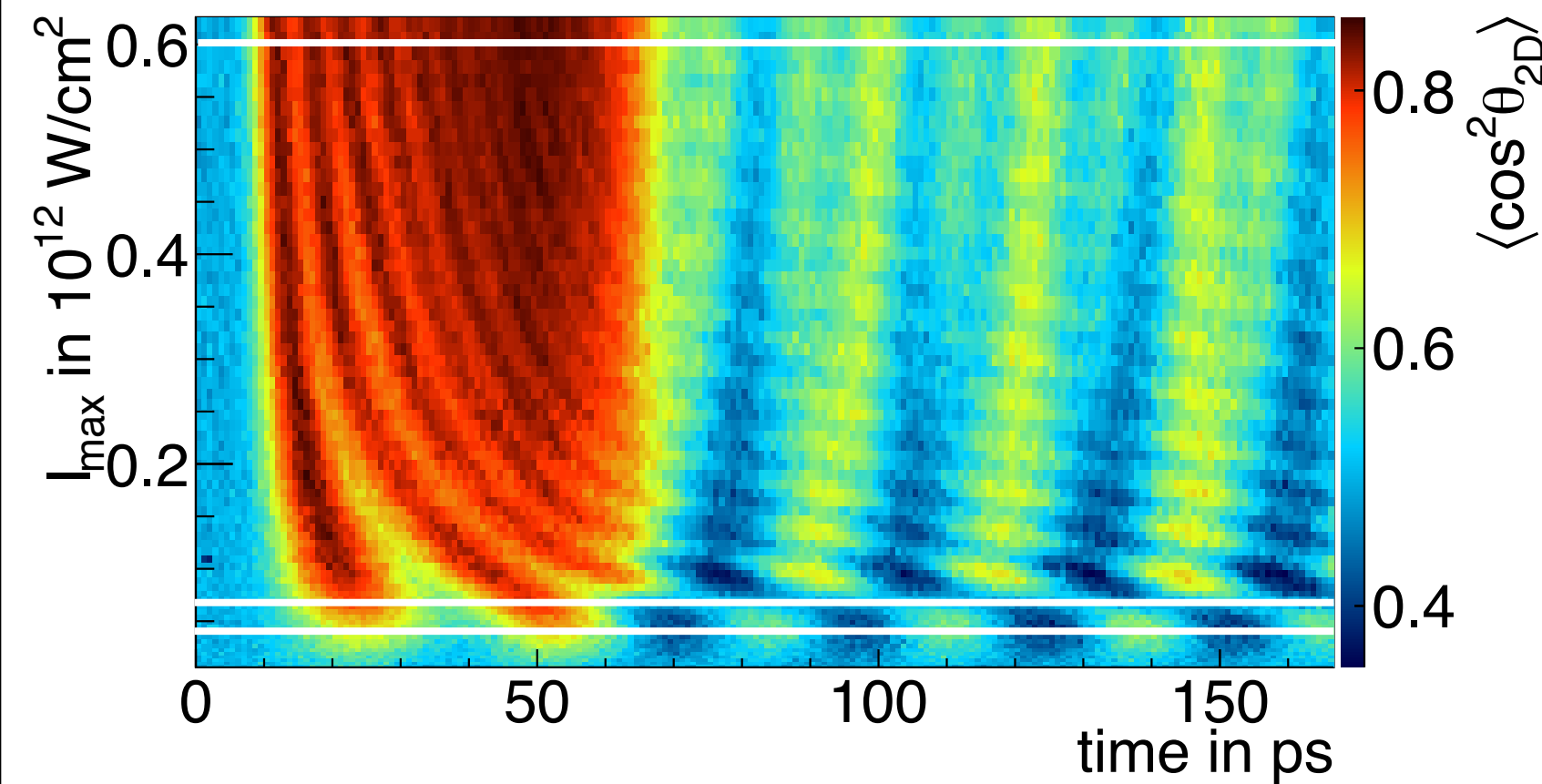
Adiabatic alignment with a 485 ps pulse



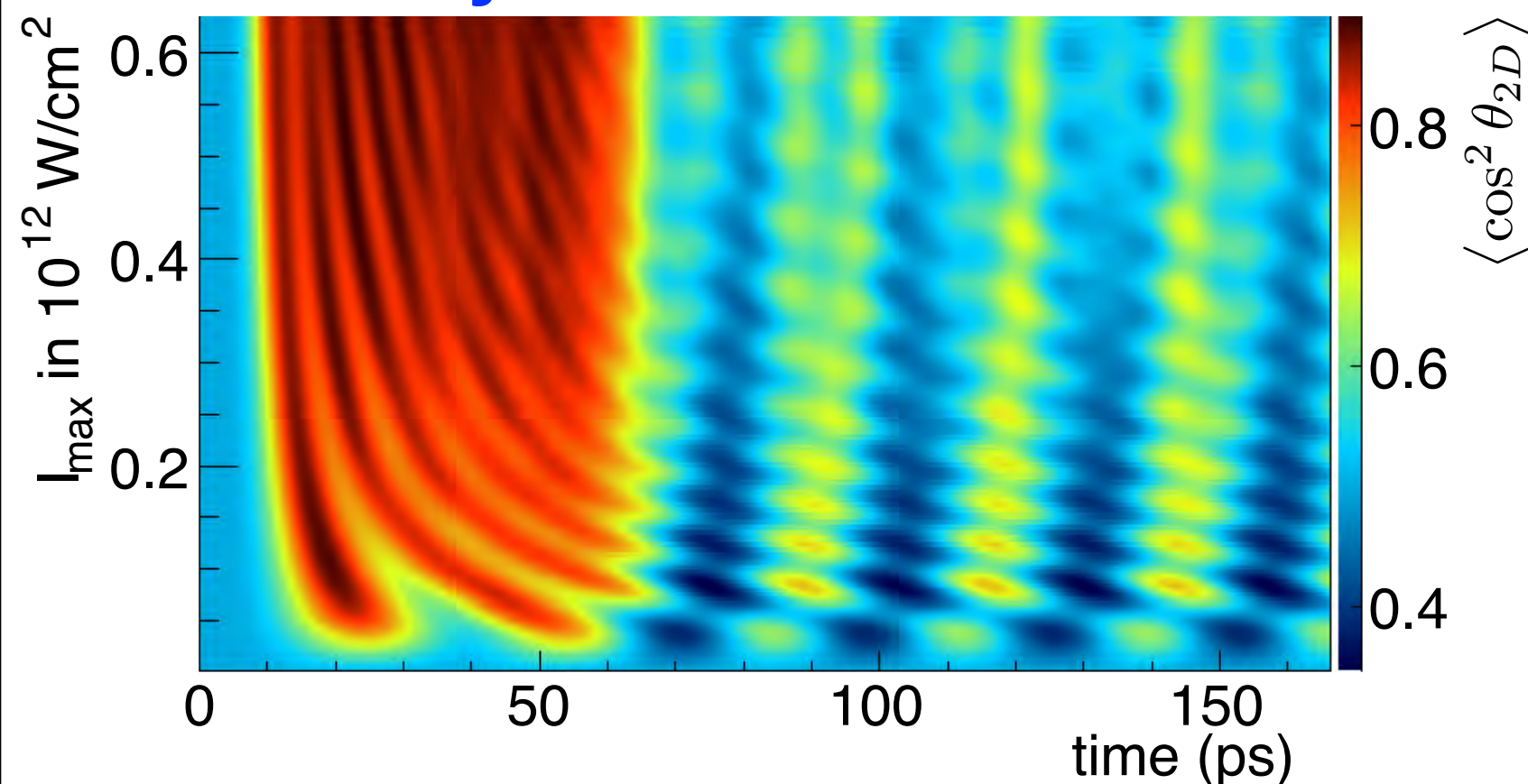
Scenarios of rotational dynamics in OCS ($X, v=0, J=0$)

Intermediate-case alignment with a 50 ps pulse

experiment



theory

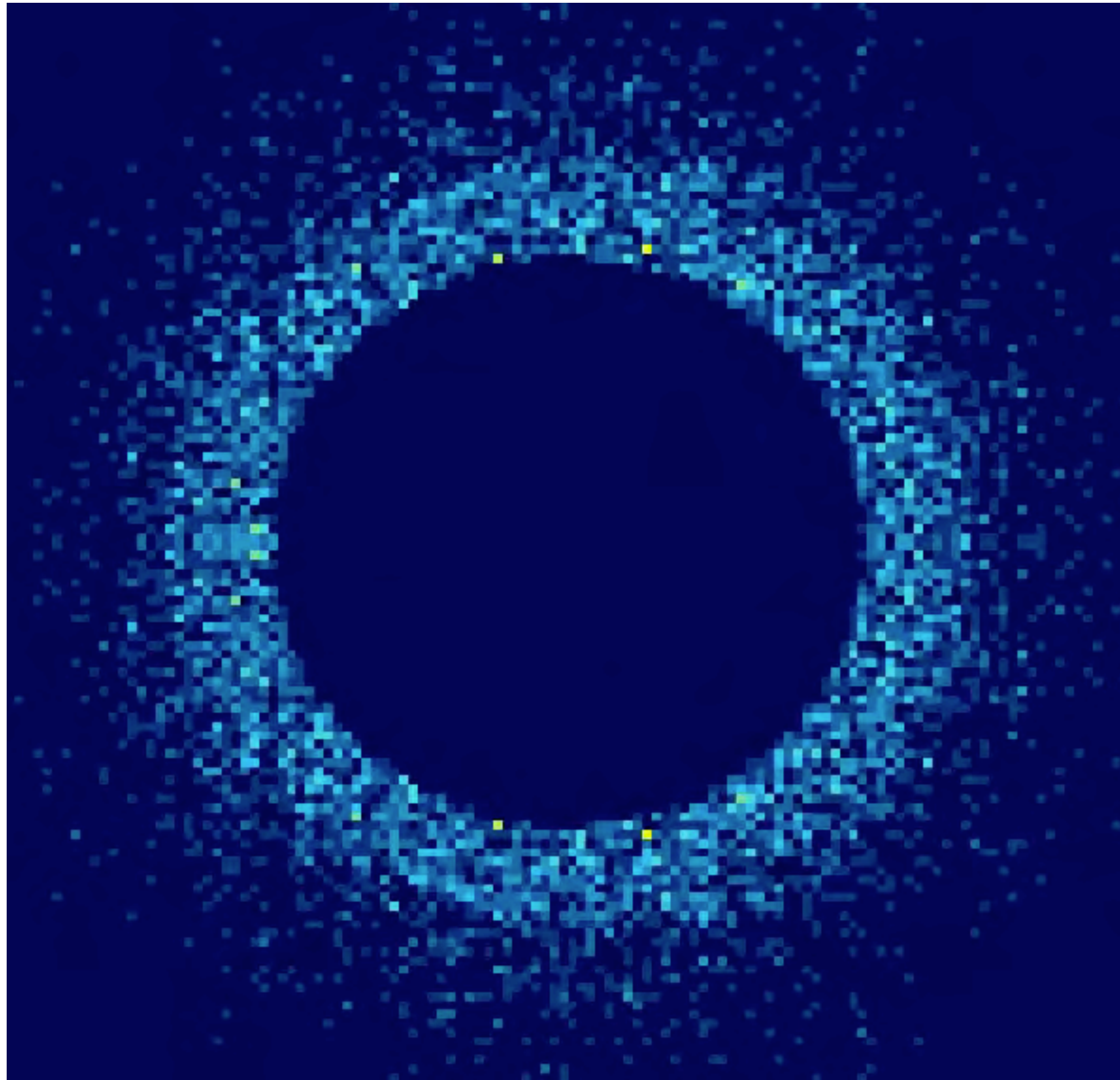


A simple two state wave packet,
a working coherent-control experiment
and a strongly-driven quantum pendulum

Achievable degree of alignment
is comparable to adiabatic case!

Scenarios of rotational dynamics in OCS (X , $v=0$, $J=0$)

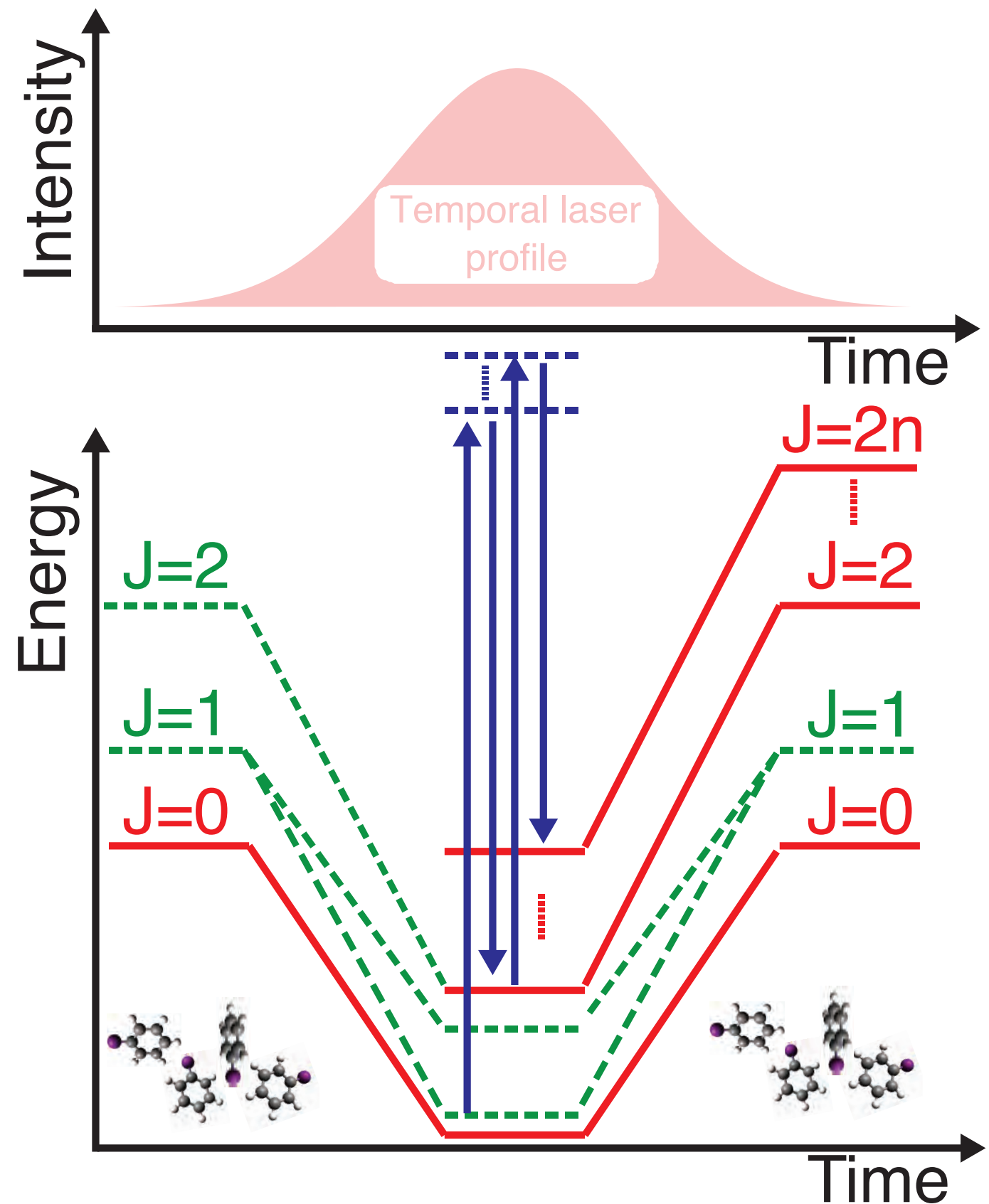
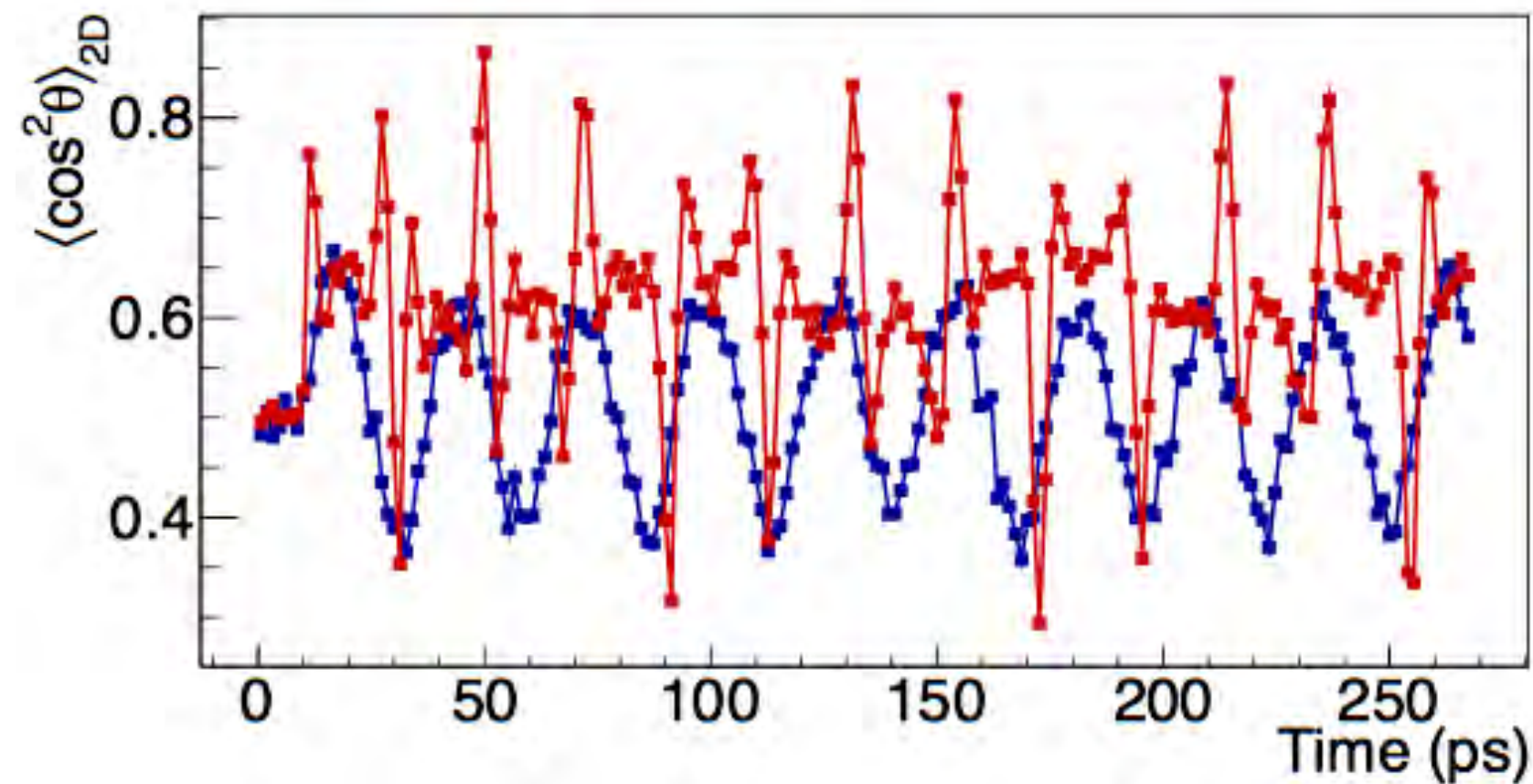
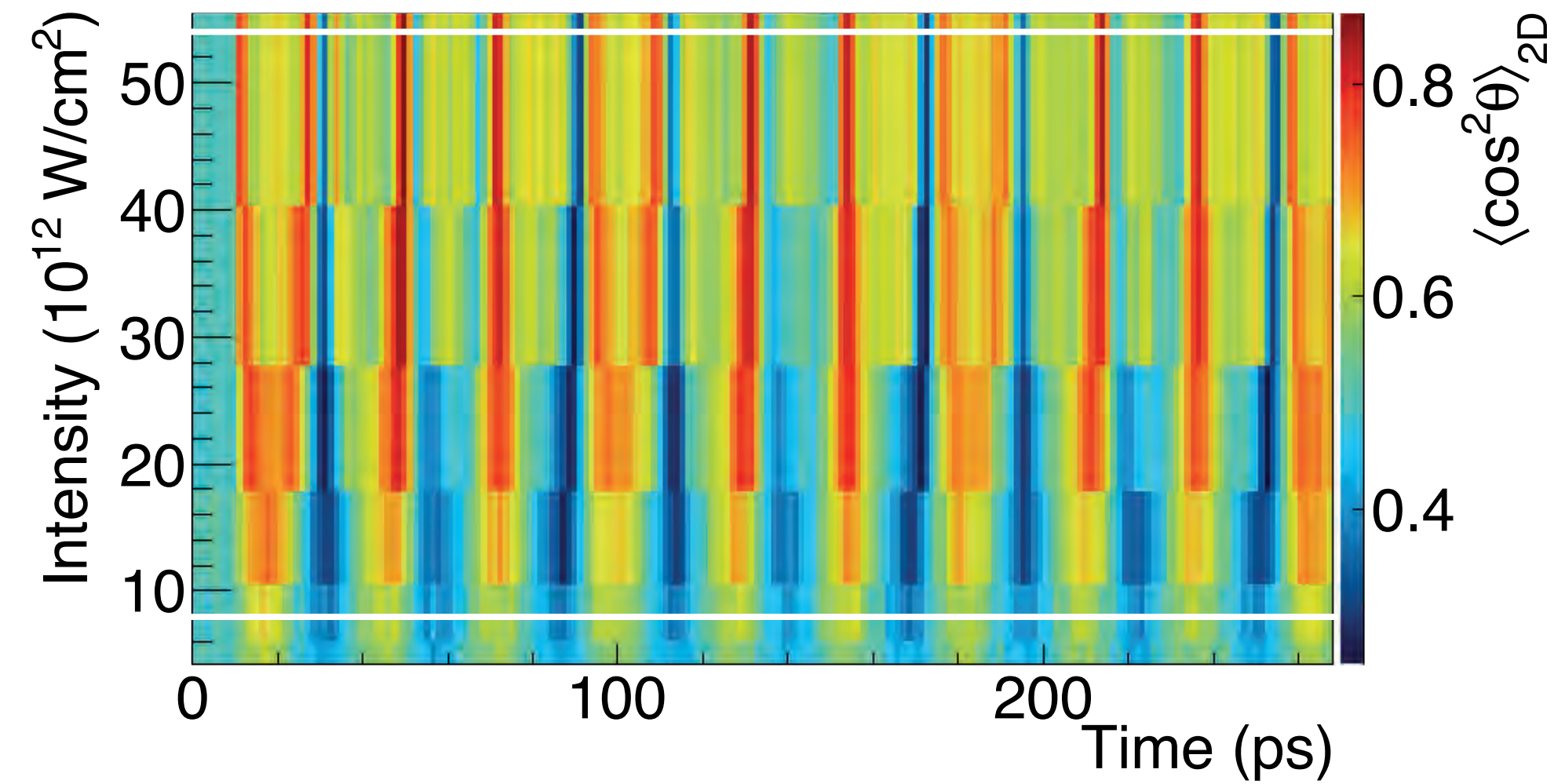
Intermediate-case alignment with a 50 ps pulse



A “molecular movie” at 10^{12} slowdown (80 ps in 80 s)

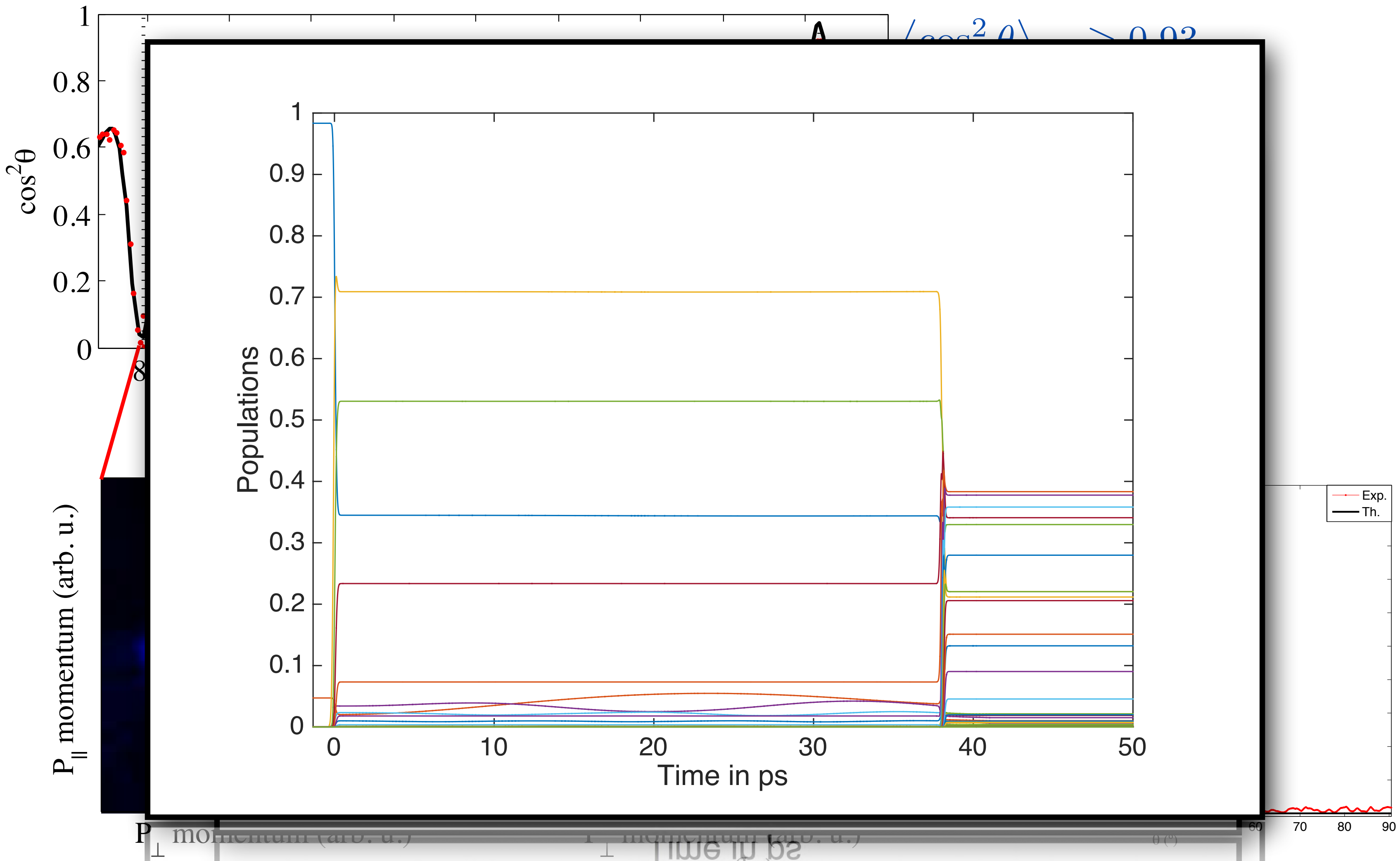
Scenarios of rotational dynamics in OCS ($X, v=0, J=0$)

Impulsive alignment with a 50 fs pulse



Scenarios of rotational dynamics in OCS ($X, v=0, J=0$)

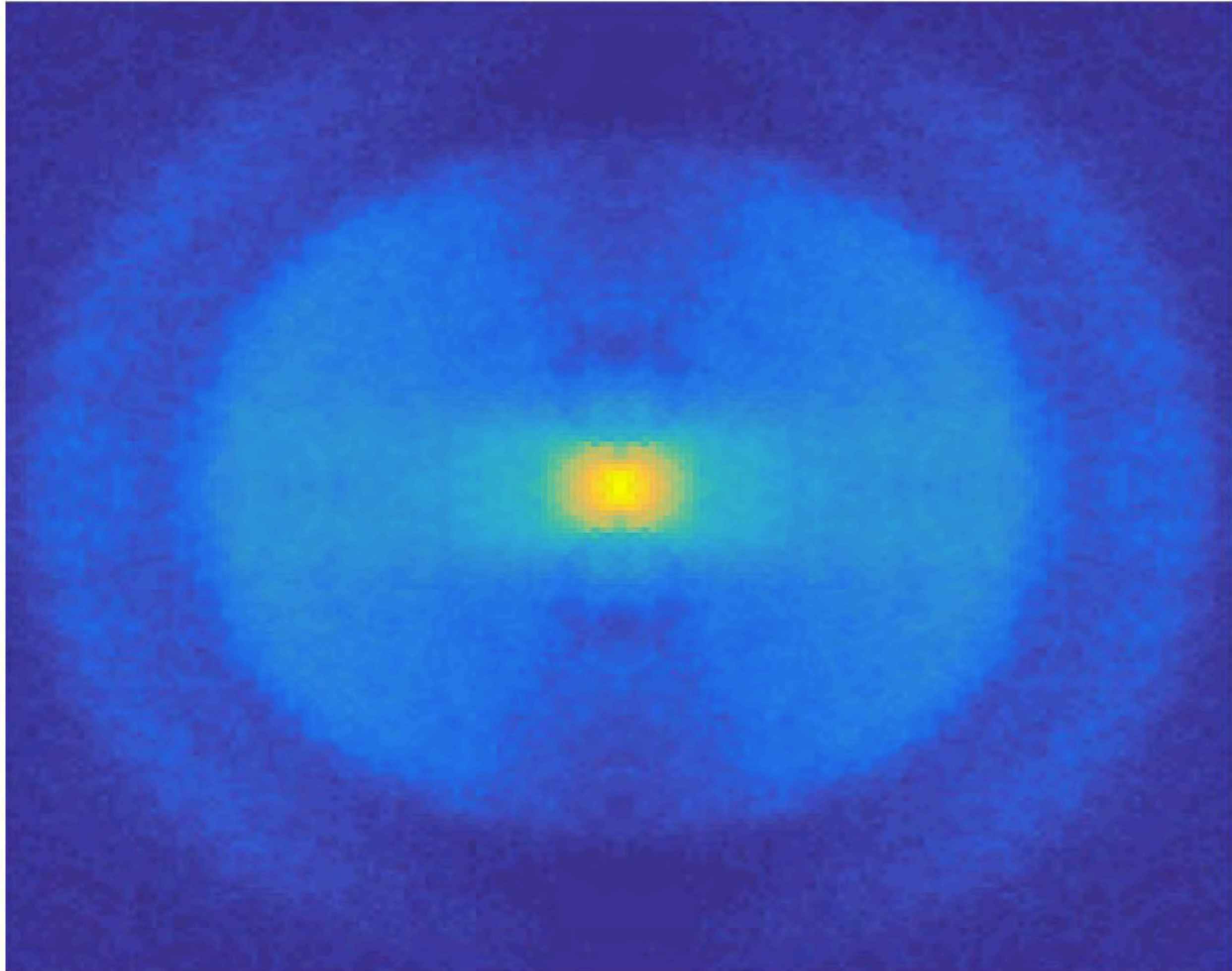
Very strong alignment using two short pulses



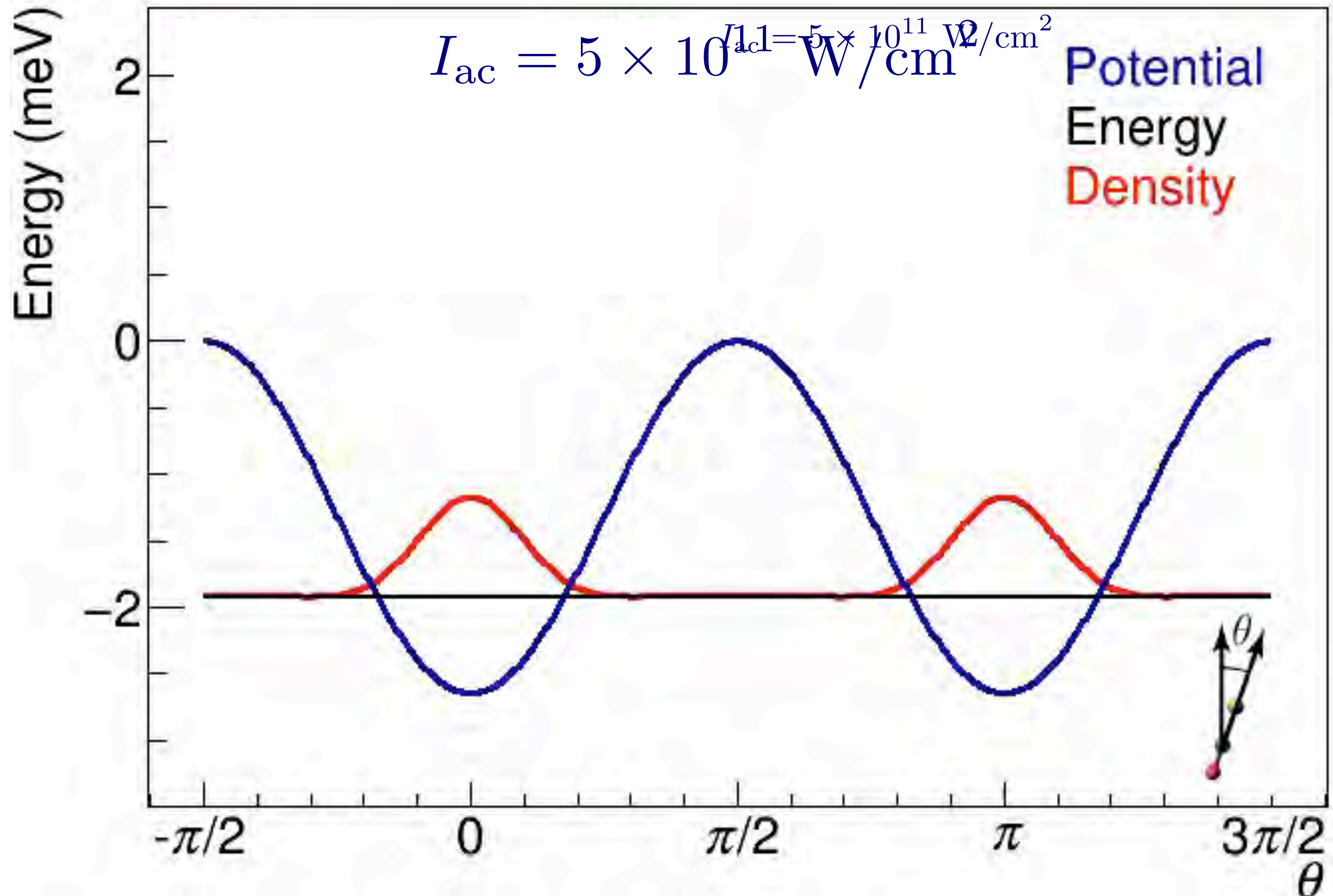
Scenarios of rotational dynamics in OCS ($X, v=0, J=0$)

Imaging quantum-revivals

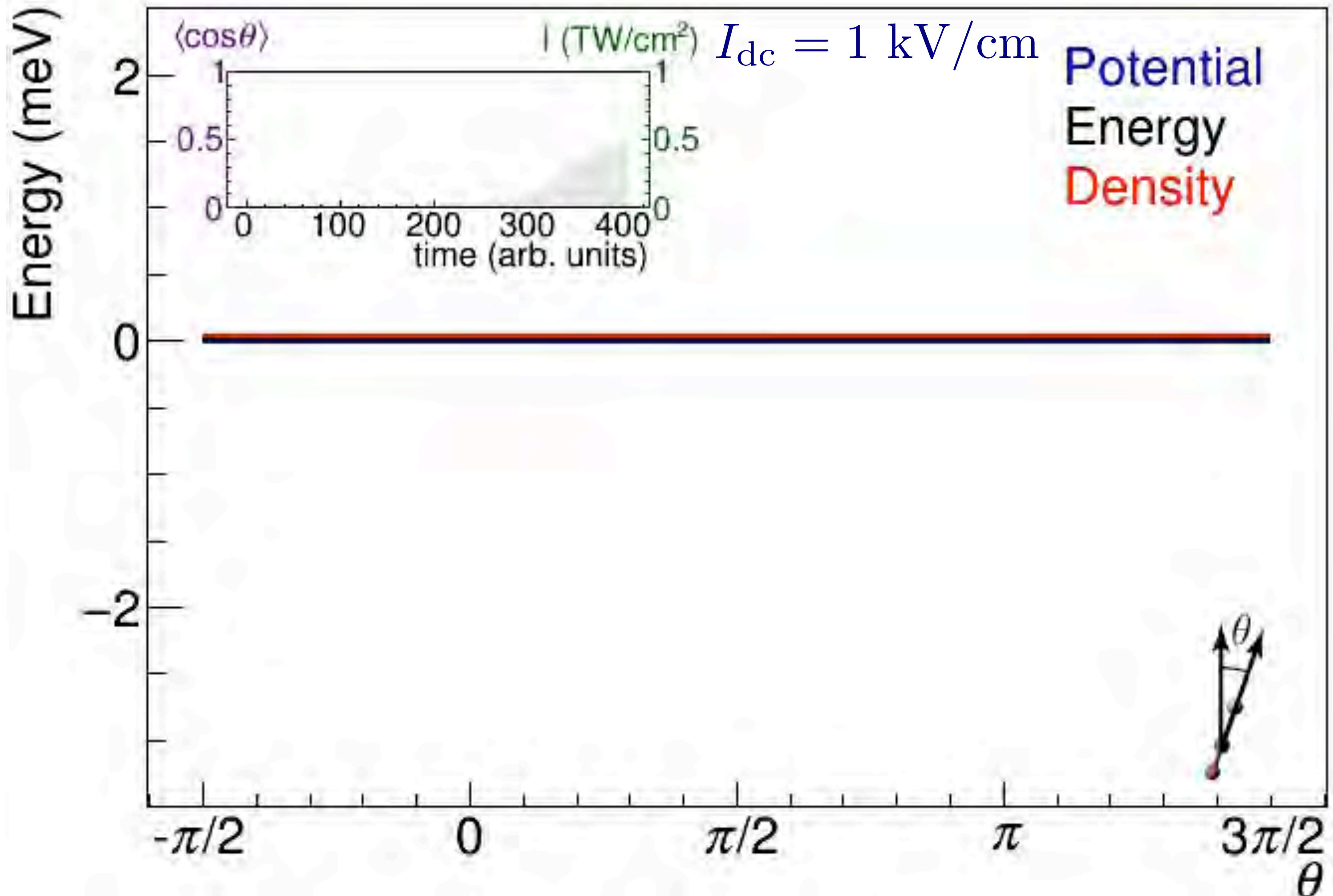
-0.75355 ps



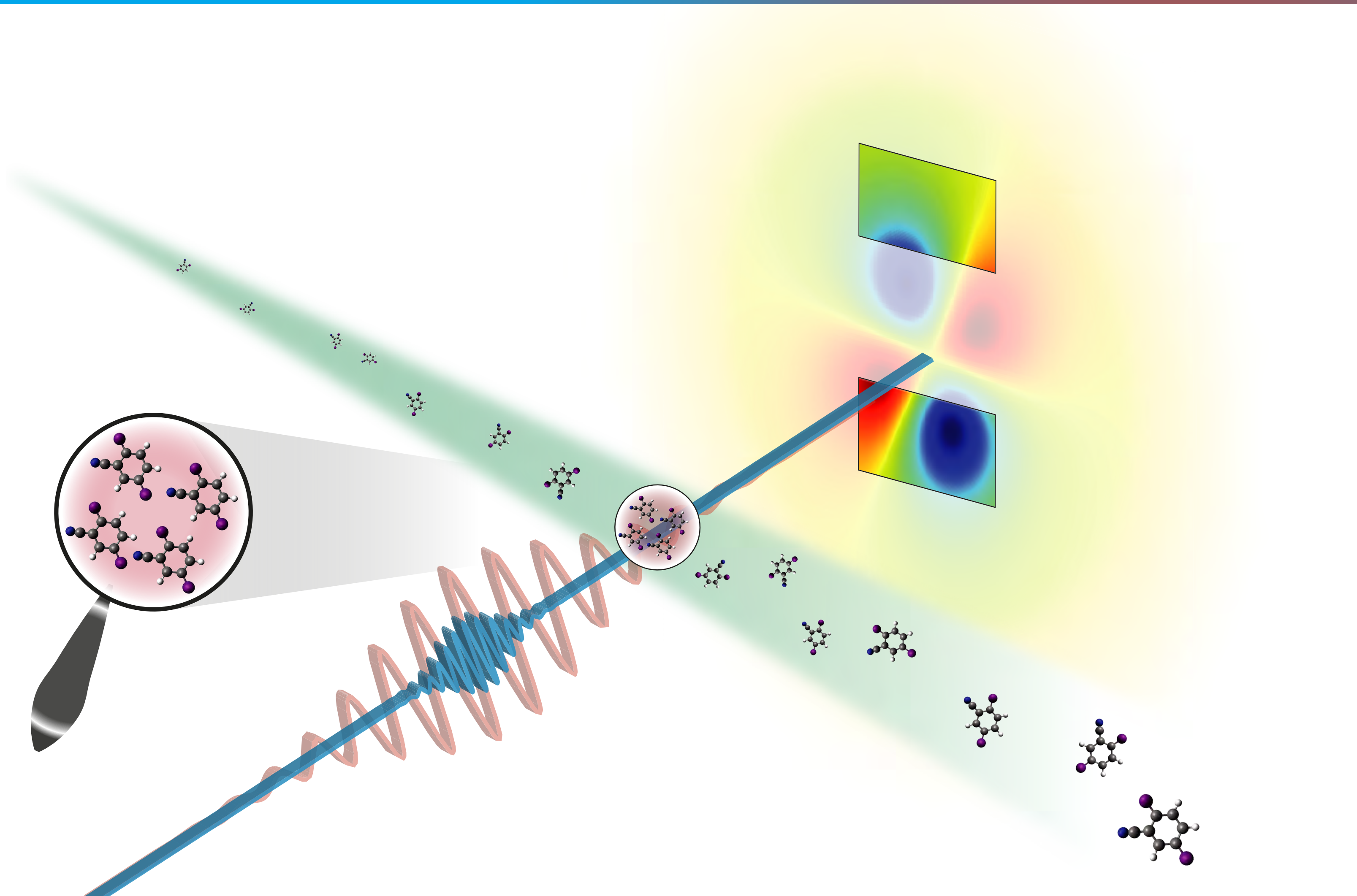
Mixed-field orientation: *The flea on Schrödinger's cat* or how tiny perturbations localize the wave function



Building a double-minimum potential with pertinent asymmetry

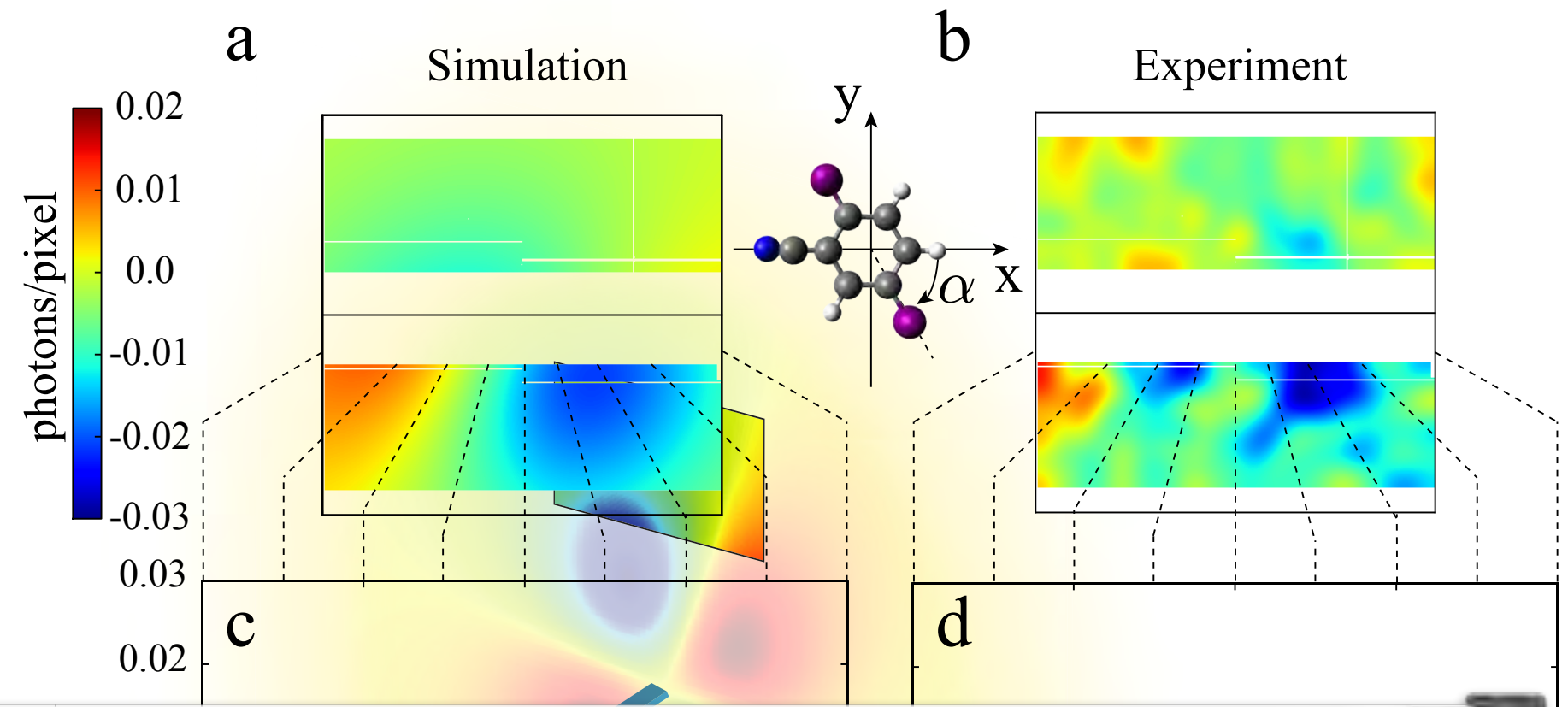
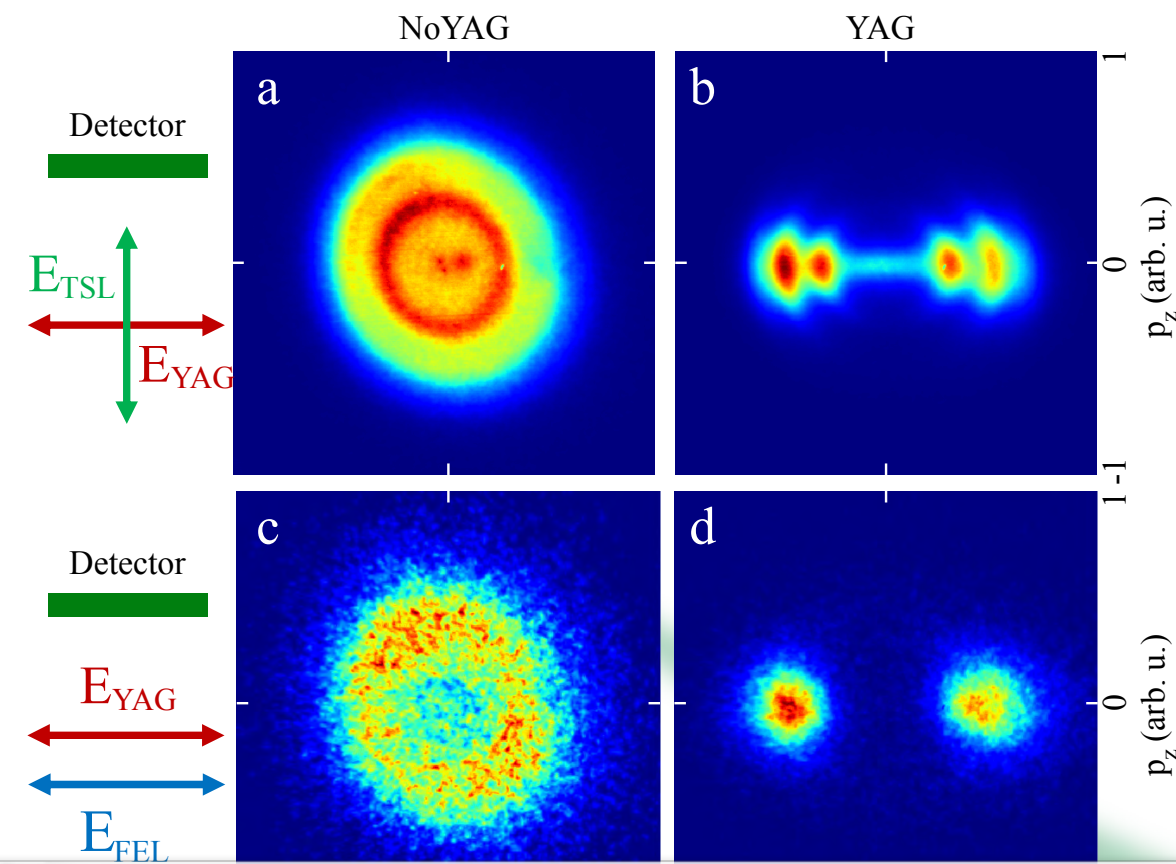


Imaging molecular dynamics with controlled molecules



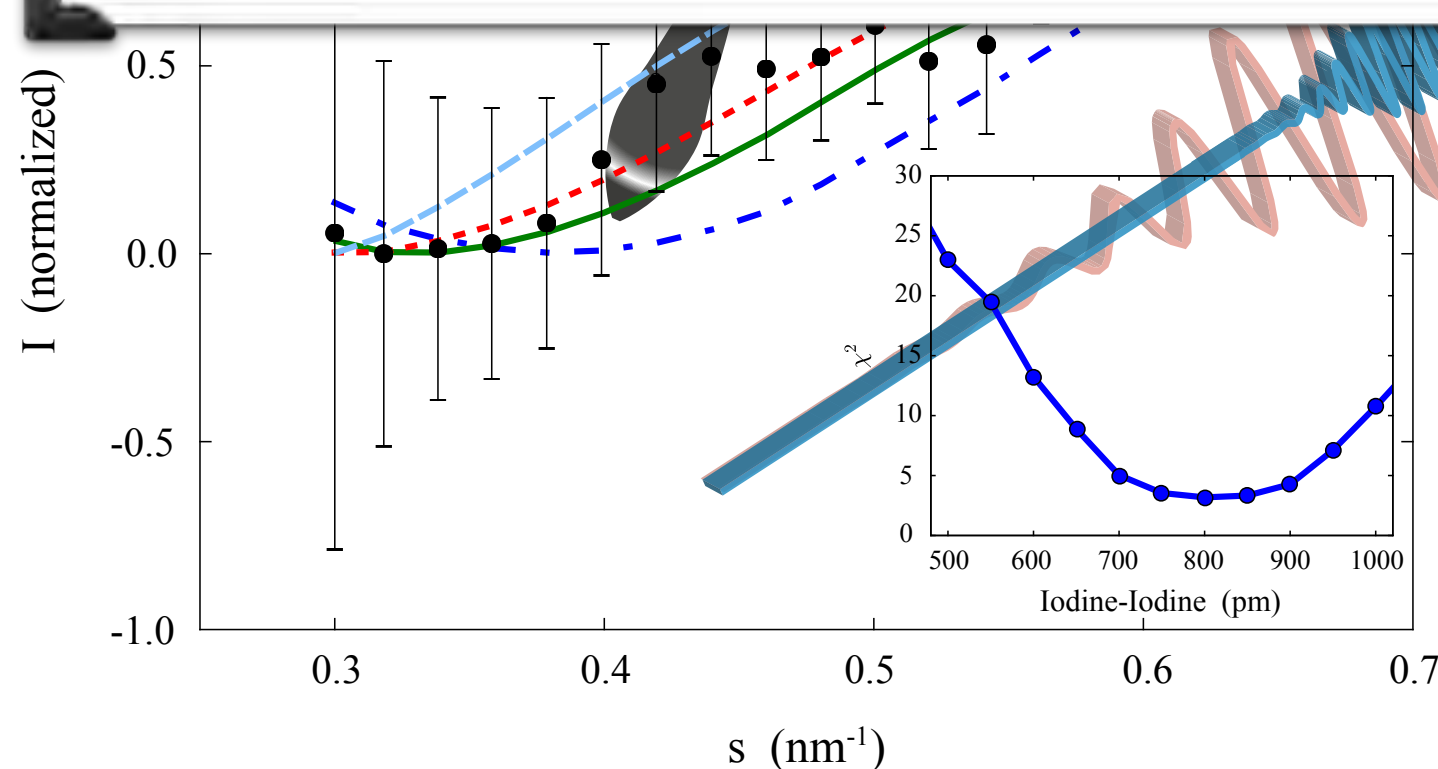
Coherent diffractive imaging of isolated molecules

100 pm precision from 0.5 nm^{-1} diffraction ($\lambda = 620 \text{ pm}$)



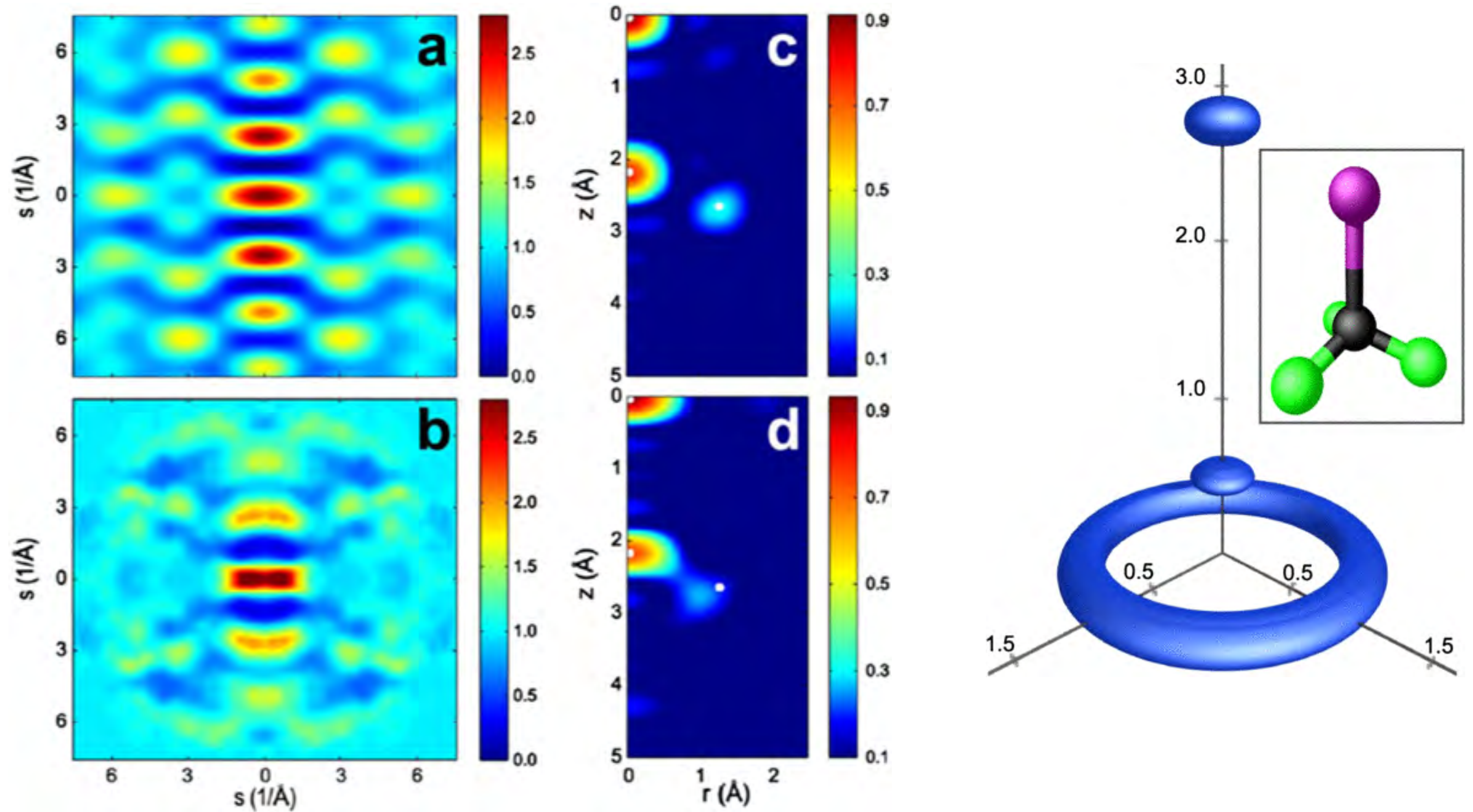
simultaneous diffractive-imaging determination of “atomic resolution” distance and angle

electron diffraction: Hensley, Yang, Centurion, *Phys. Rev. Lett.* **109**, 133202 (2012)

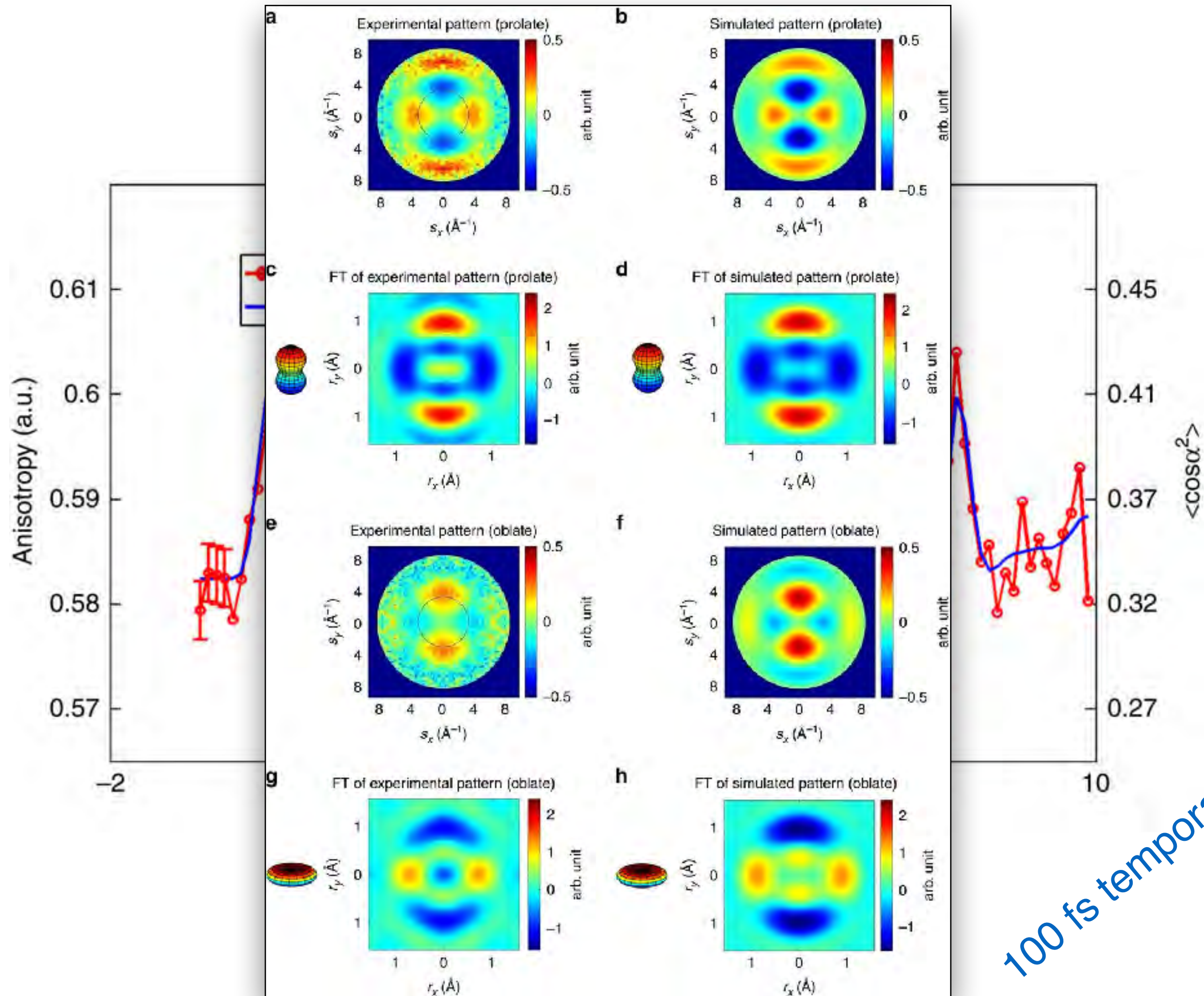


diffraction data yields
 $\langle \cos^2 \theta \rangle_{2D} = 0.8$ (vs. 0.84)
 $r(I-I) \approx 800 \text{ pm}$ (vs. 700 pm)

Imaging of Isolated Molecules with Ultrafast Electron Pulses



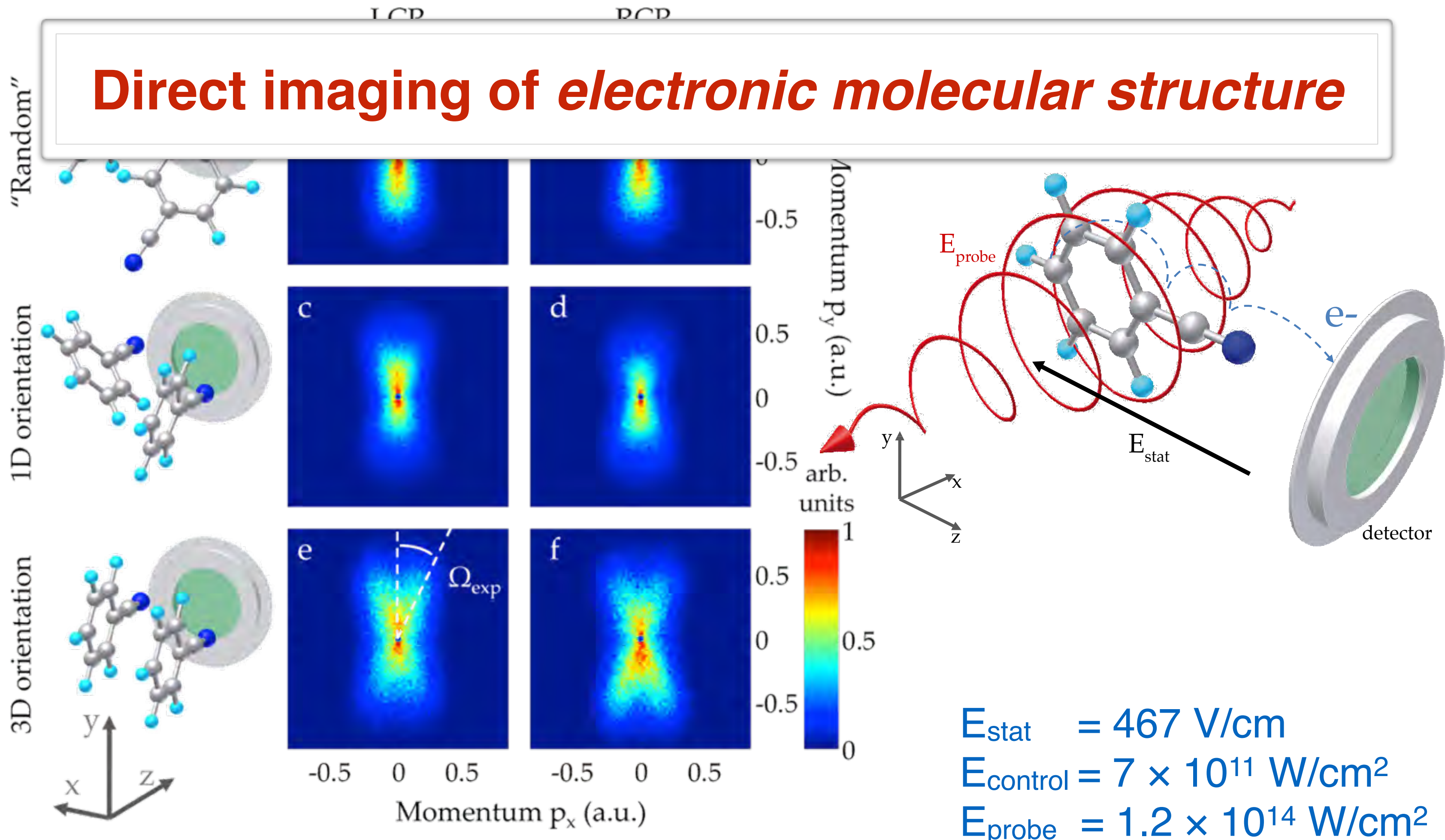
Diffractive imaging of (rotational) dynamics



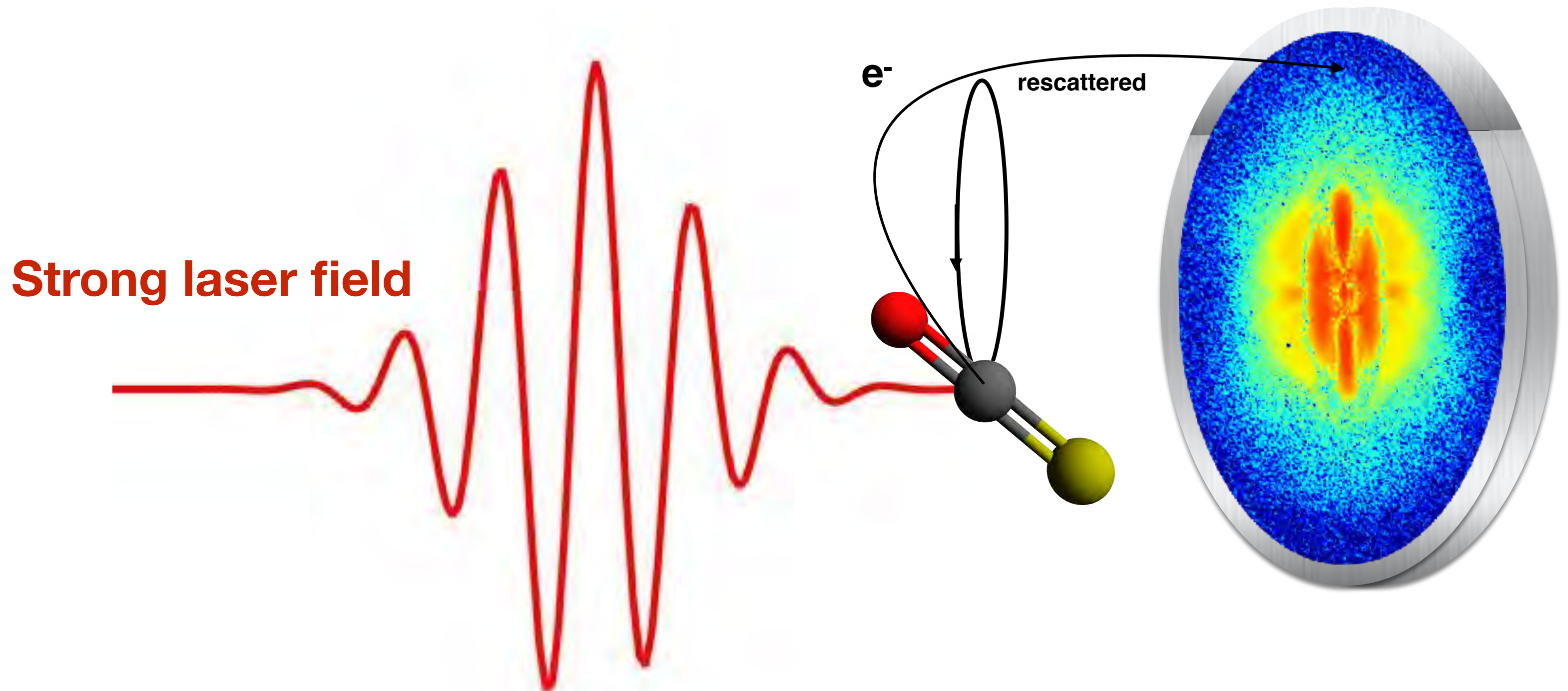
100 fs temporal resolution!

Molecular frame photoelectron angular distributions of 3D-oriented benzonitrile molecules

Direct imaging of *electronic molecular structure*



Laser induced electron diffraction



Direct imaging of *nuclear molecular structure*

“game”: electron trajectories in an external field

